

Magnetic properties of [NiFe/Au/Co/Au] multilayers

magnetostatic coupling and giant magnetoresistance

Zmniejszona jakość ilustracji
Low quality version
Zakopane
2009

Maciej Urbaniak, Feliks Stobiecki, Bogdan Szymański
Institute of Molecular Physics, Polish Academy of Sciences

Magnetic properties of [NiFe/Au/Co/Au] multilayers

magnetostatic coupling and giant magnetoresistance

Cooperation between:

Institute of Molecular Physics, Polish Academy of Sciences, Poznań, Poland

Prof. A. Ehresmann

Department of Physics , University of Kassel, Kassel, Germany

Prof. M. Kopcewicz

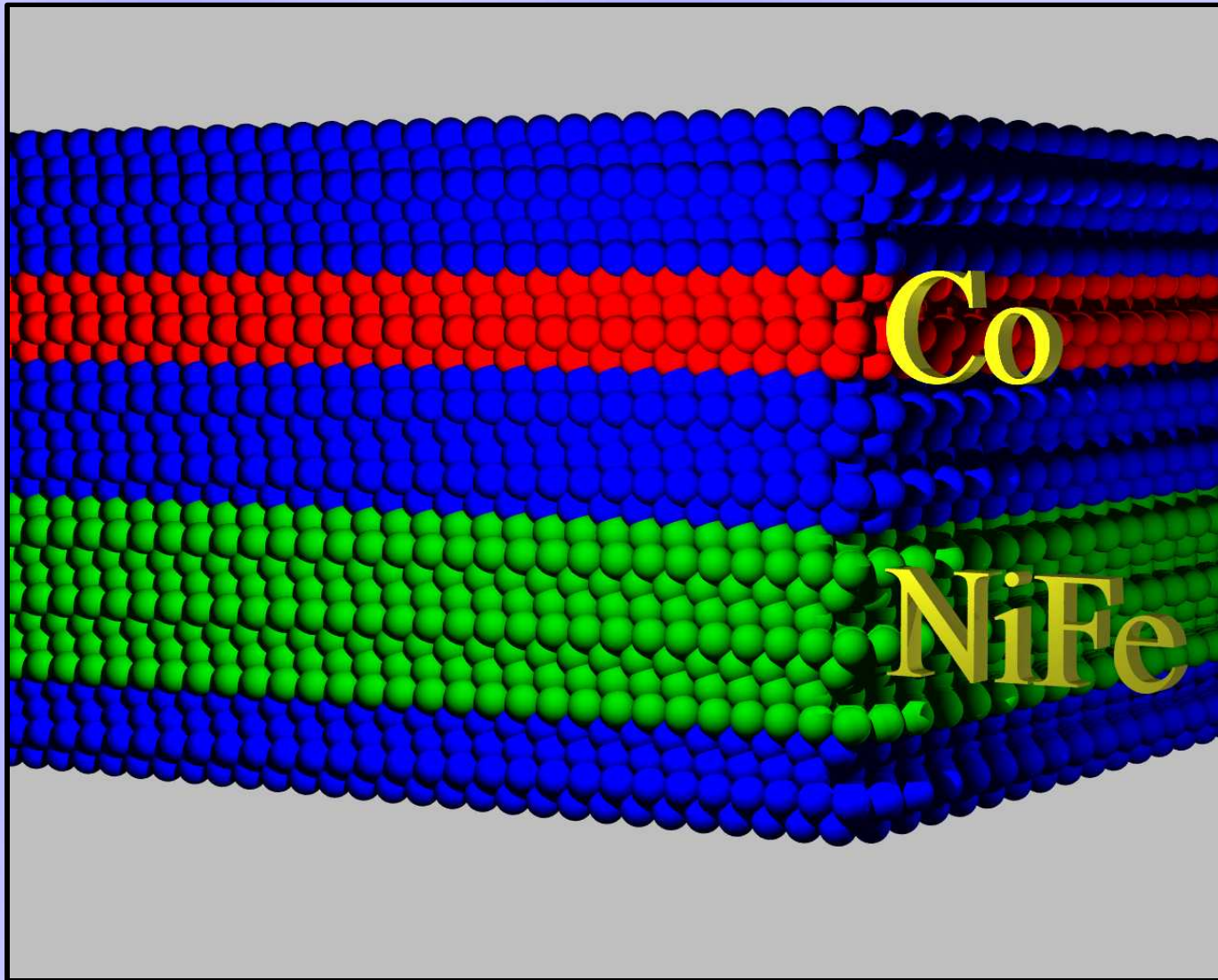
Institute of Electronic Materials Technology, Warszawa, Poland

Prof. A. Maziewski

Laboratory of Magnetism, Faculty of Physics, University of Białystok, Poland

Magnetic properties of [NiFe/Au/Co/Au] multilayers

- Introduction
- Magnetic properties
- Giant magnetoresistance
- Magnetostatic coupling
- Magnetic patterning
- Conclusions



Substrate:
naturally oxidized
Si(100)

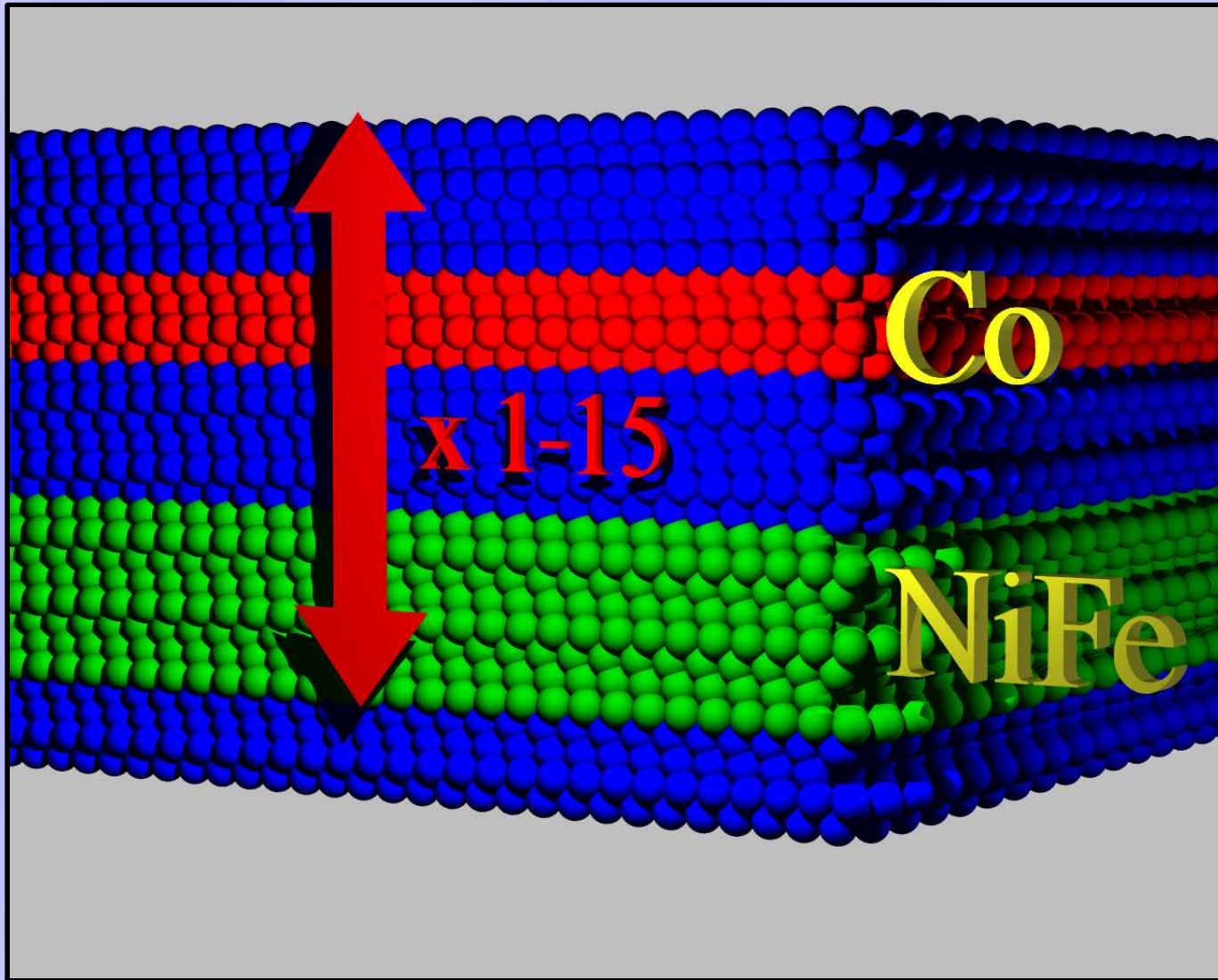
$$t_{\text{Co}(\text{CoFe})} = 0.2-1.5 \text{ nm}$$

$$t_{\text{NiFe}} = 0.5-4 \text{ nm}$$

$$t_{\text{Au}} = 1.5-3 \text{ nm}$$

Magnetron sputtering

Introduction



Substrate:
naturally oxidized
Si(100)

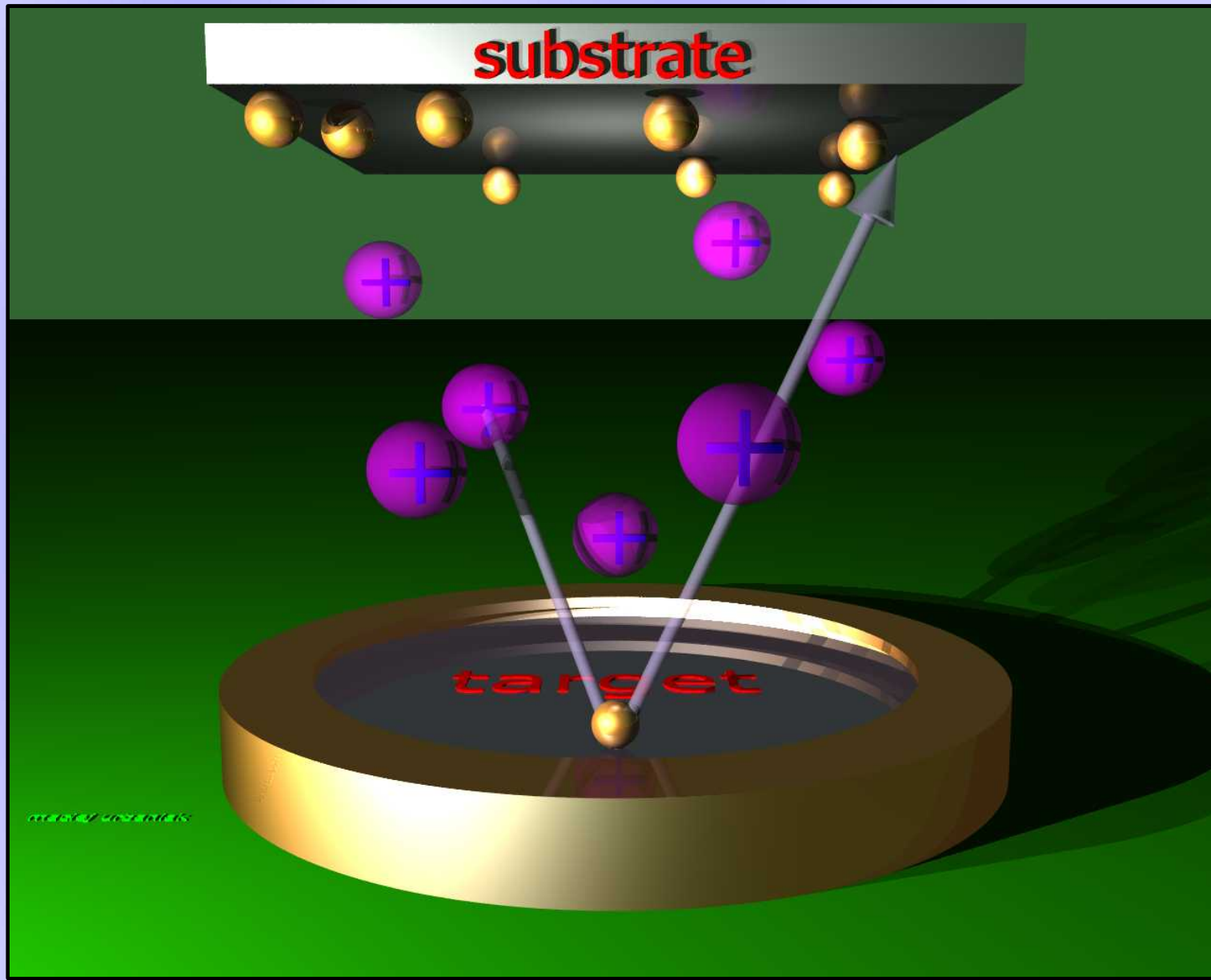
$$t_{\text{Co}(\text{CoFe})} = 0.2-1.5 \text{ nm}$$

$$t_{\text{NiFe}} = 0.5-4 \text{ nm}$$

$$t_{\text{Au}} = 1.5-3 \text{ nm}$$

Magnetron sputtering

Introduction



Target:

NiFe, Au, Co

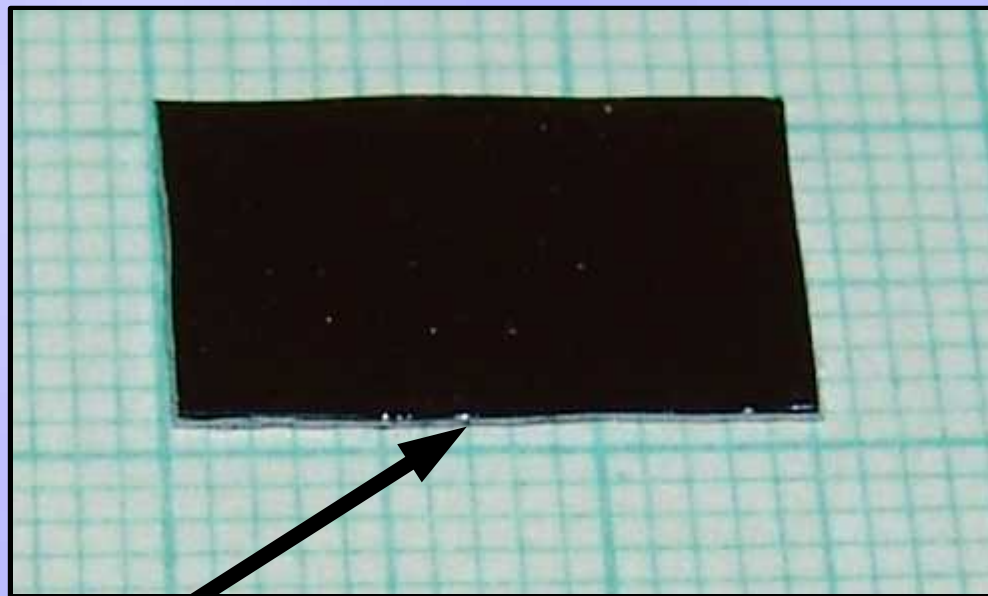
Substrate:

Si(100), glass,
adhesive tape

Magnetron sputtering

Target-
negative potential

Introduction

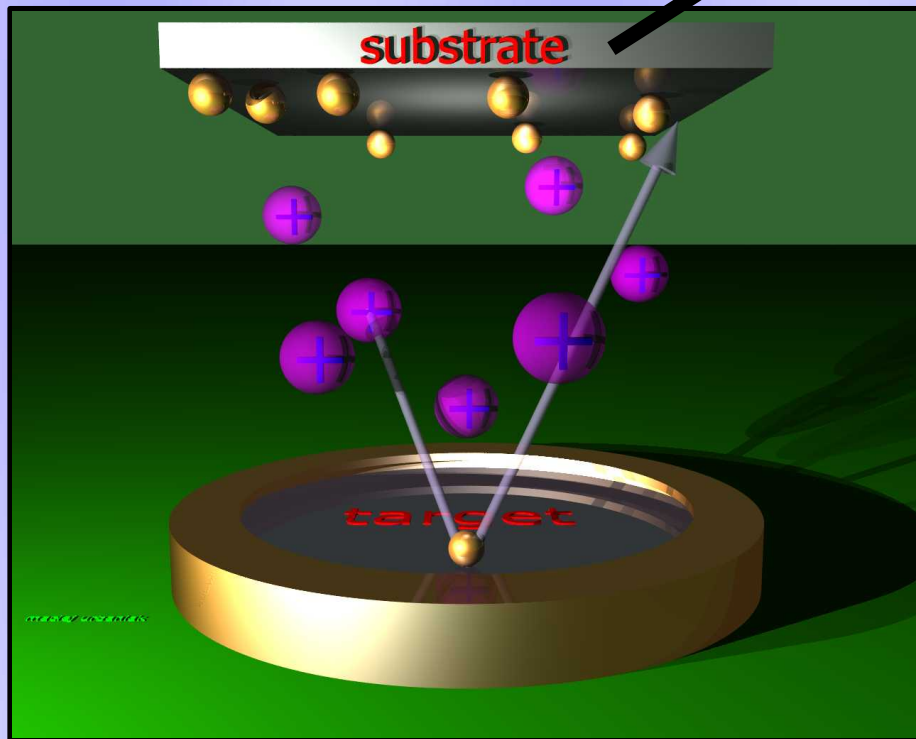


Target:

NiFe, Au, Co

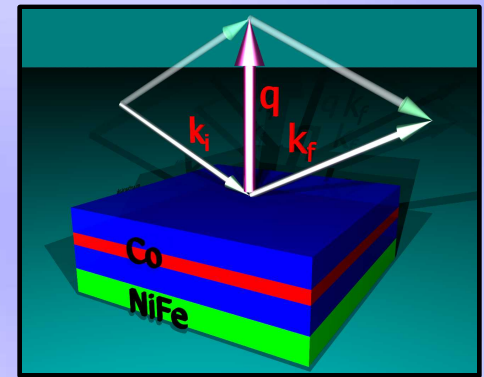
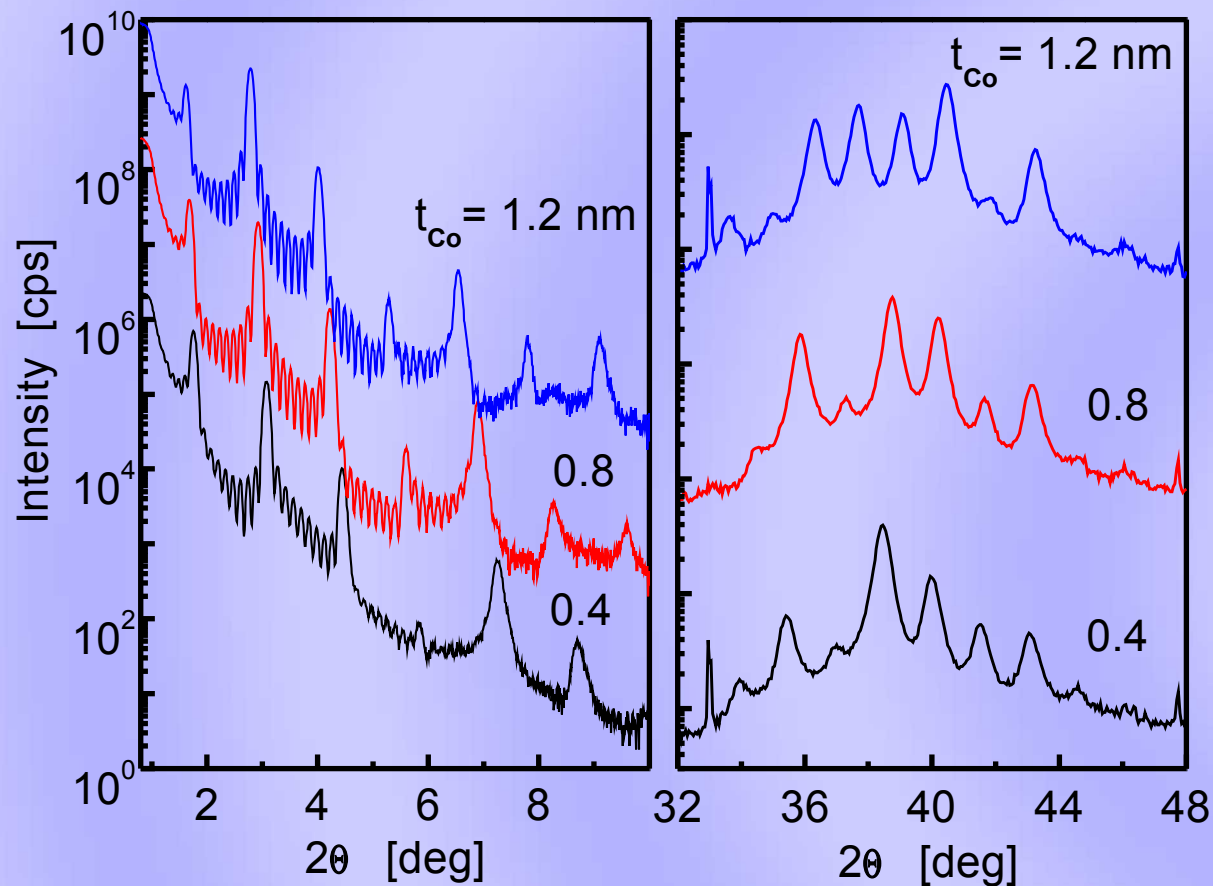
Substrate:

Si(100), glass,
adhesive tape



Magnetron sputtering

Target-
negative potential

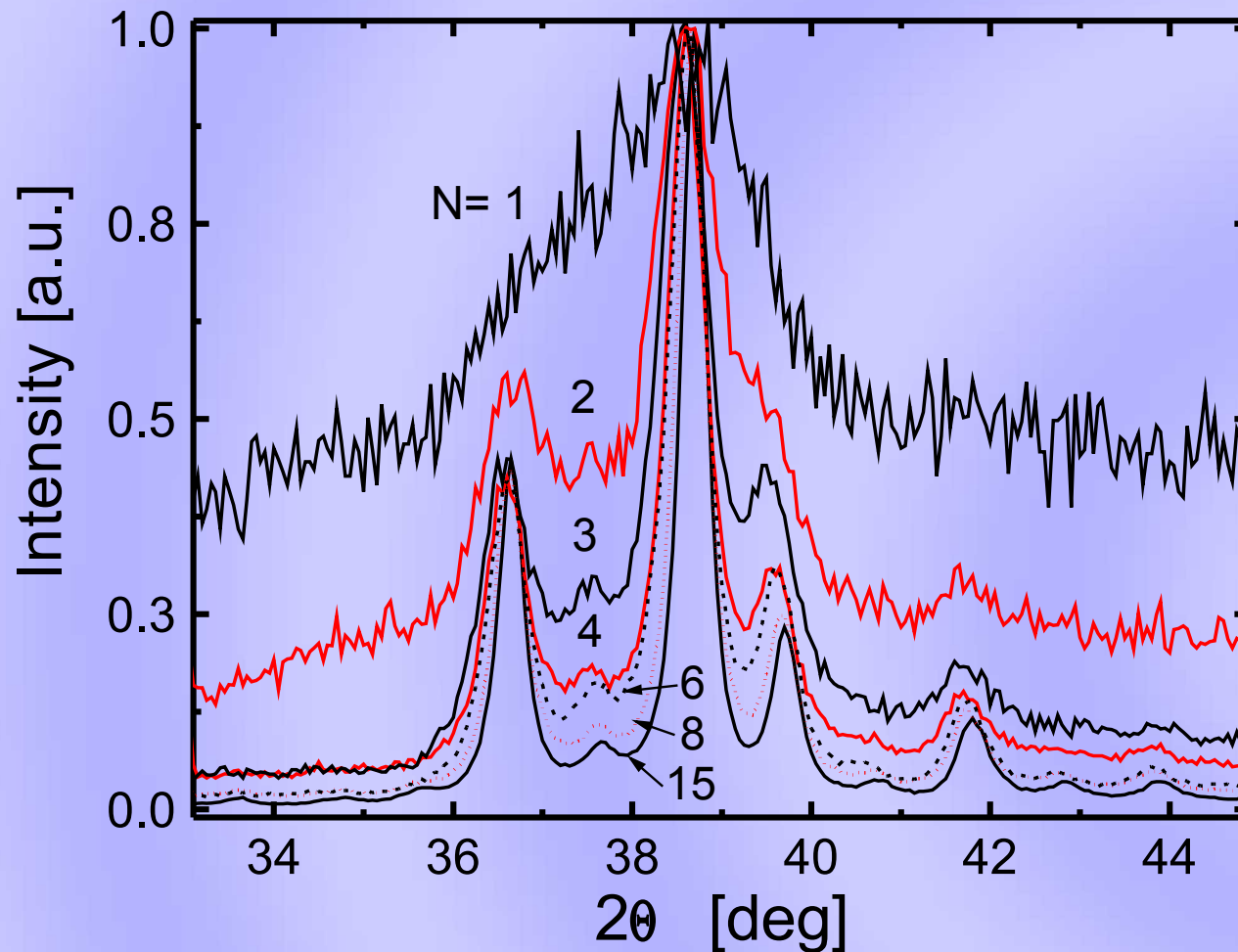


good periodicity

defined fcc(111) texture

$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(1.9 \text{ nm})/\text{Co}(t_{Co})/\text{Au}(1.9 \text{ nm})]_{10}$

Cu $K\alpha$ 0.154nm

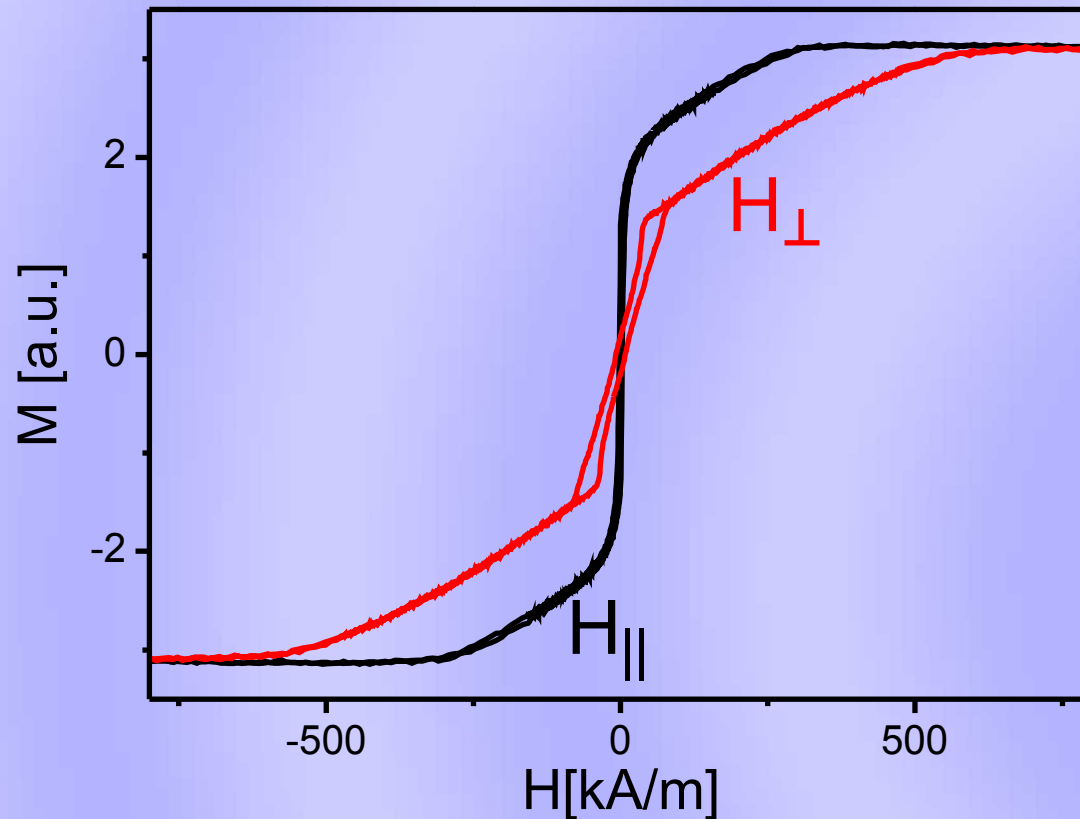


“the profile for MLs with $N = 3$ shows all features typical of profiles for large N .”

$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(3 \text{ nm})/\text{Co}(0.8)/\text{Au}(3 \text{ nm})]_N$

Cu $K\alpha$ 0.154nm

Magnetic properties

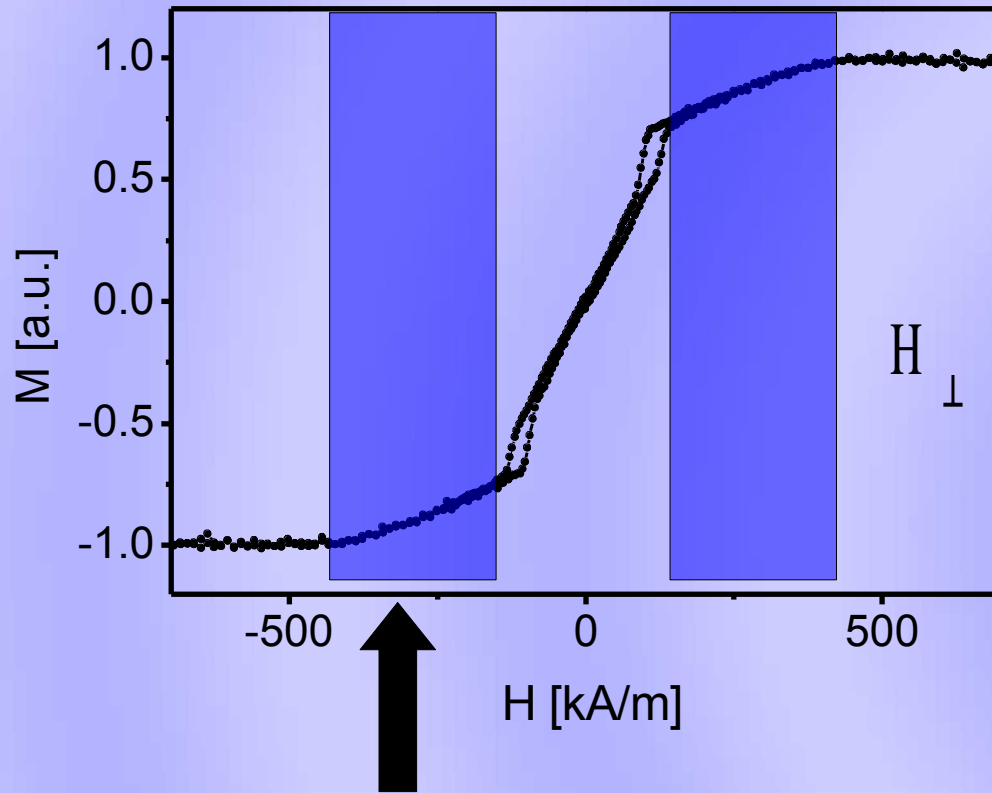


NiFe – magnetic easy axis
in-plane

Co – magnetic easy axis
perpendicular to plane of
the sample

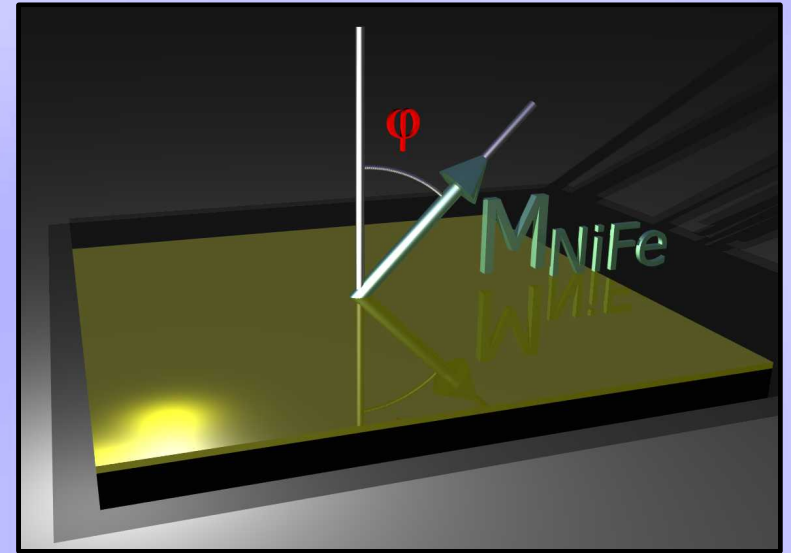
- Small field range of hysteresis with field applied perpendicularly is characteristic of systems with **stripe domains**.
- In both field configurations NiFe and Co layers reverse quasi independently.

Magnetic properties



Reversal of NiFe only

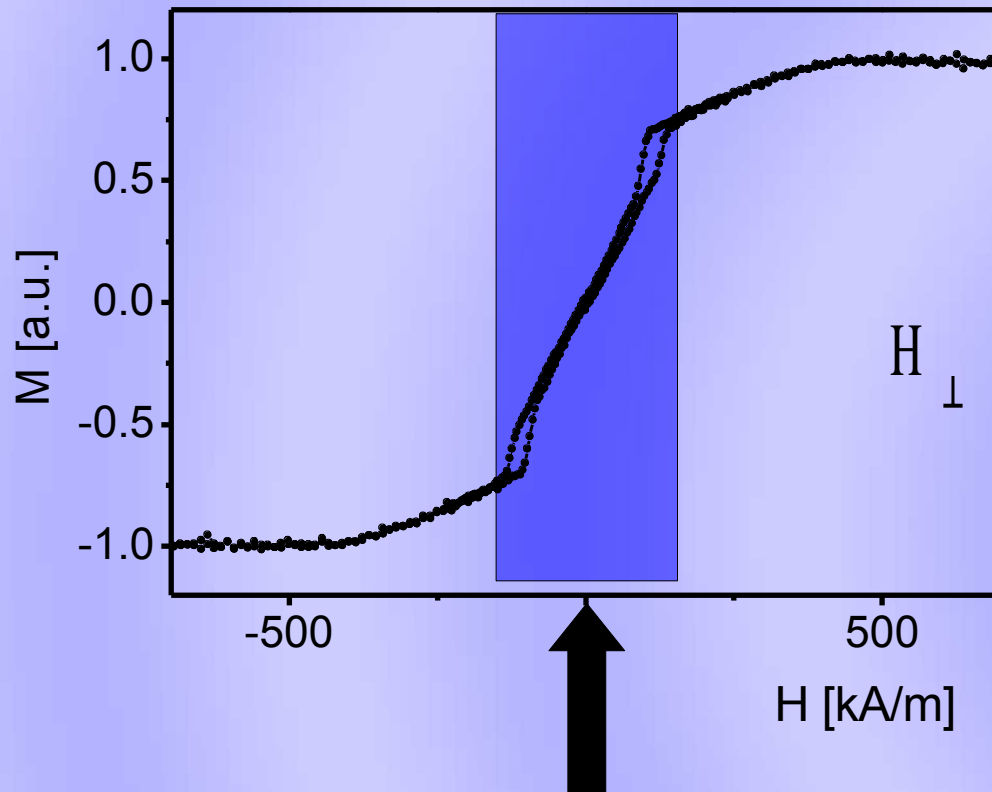
$$K_u = \frac{1}{2} \mu_0 (M_S^{NiFe})^2$$



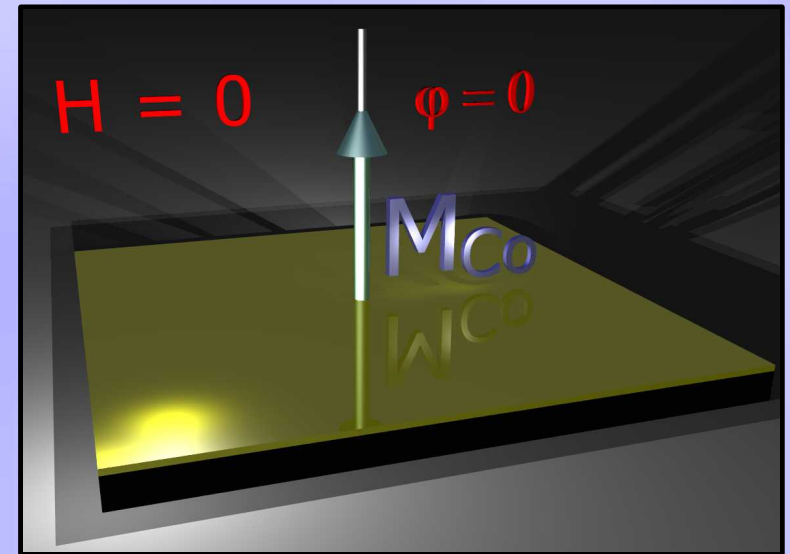
Shape anisotropy:

$$\cos(\varphi) = \frac{H}{M_S}$$

Magnetic properties

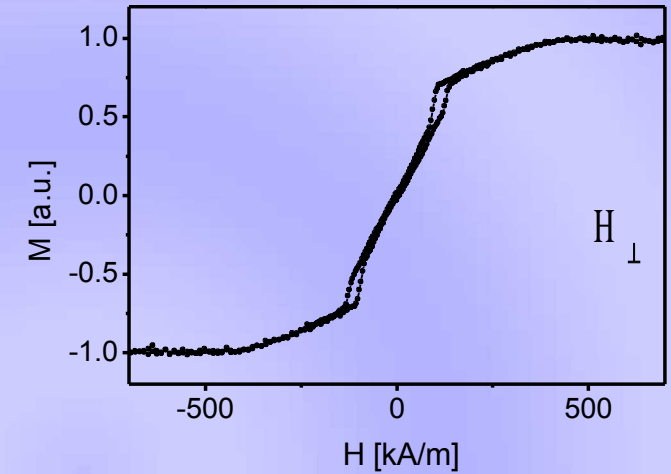
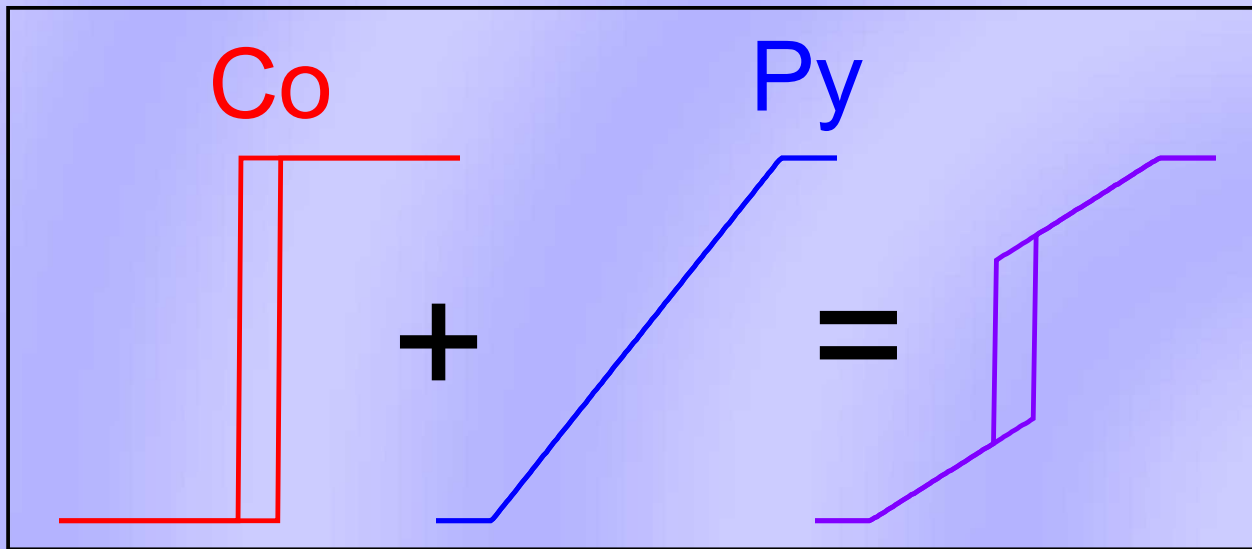


Simultaneous reversal
of
NiFe and Co



An easy axis of the Co layers
is perpendicular to surface of
multilayer.

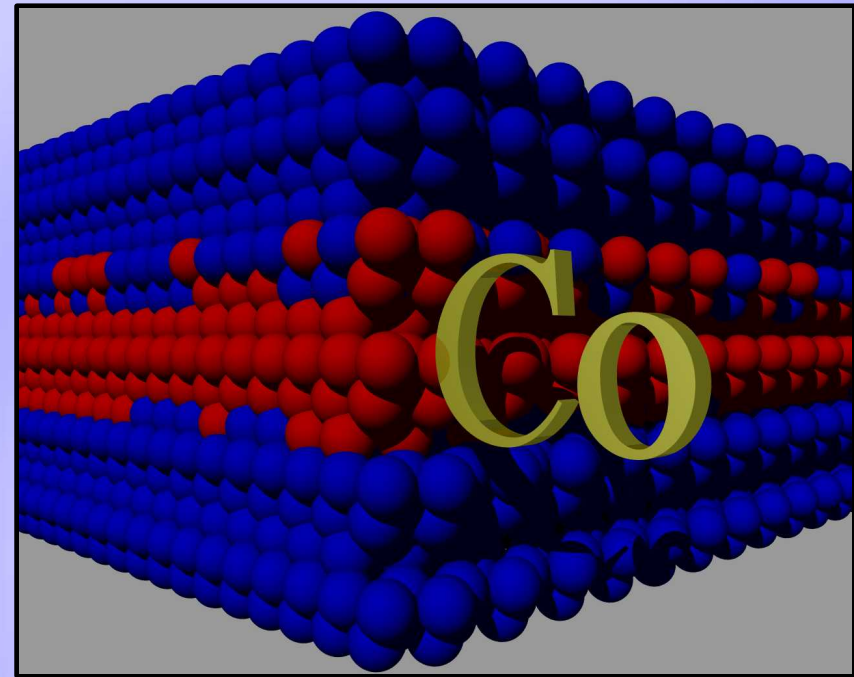
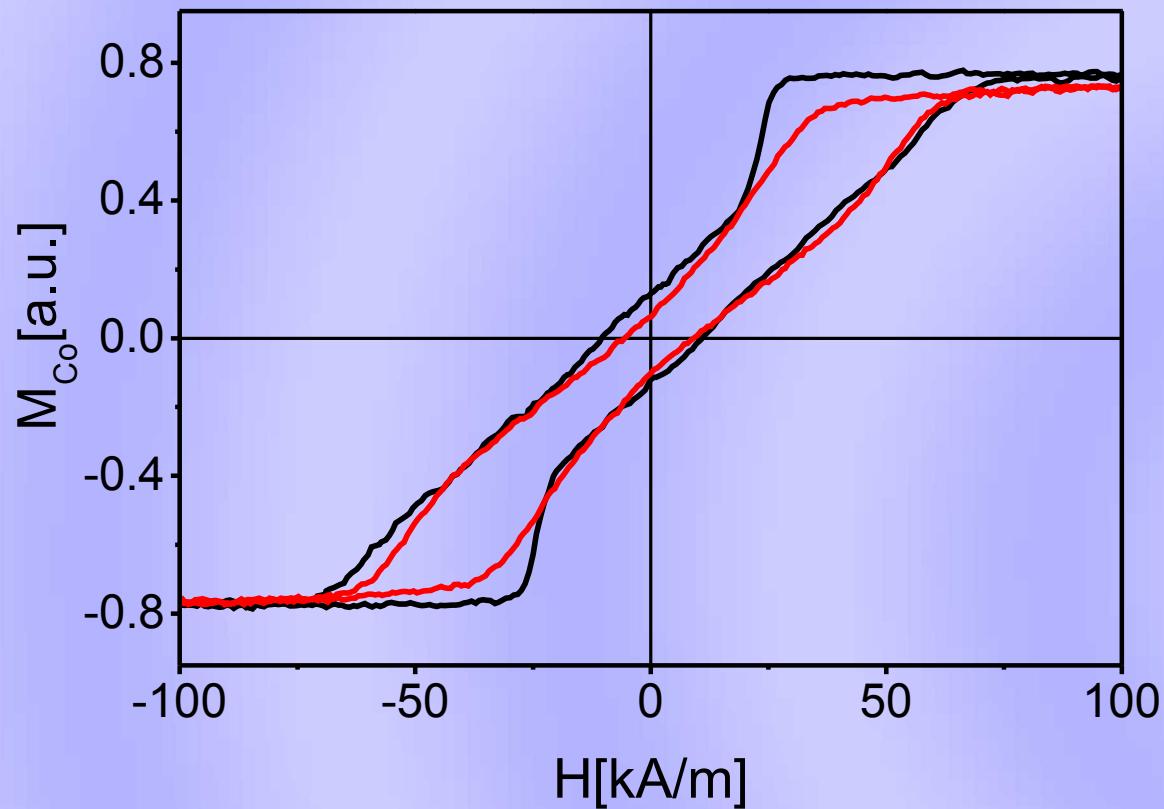
Magnetic properties



In the first approximation Co and NiFe layers can be thought of as uncoupled.

$M(H)$ dependence of the NiFe/Au/Co structure is then an arithmetic sum of the $M(H)$ dependencies of Co and NiFe layers.

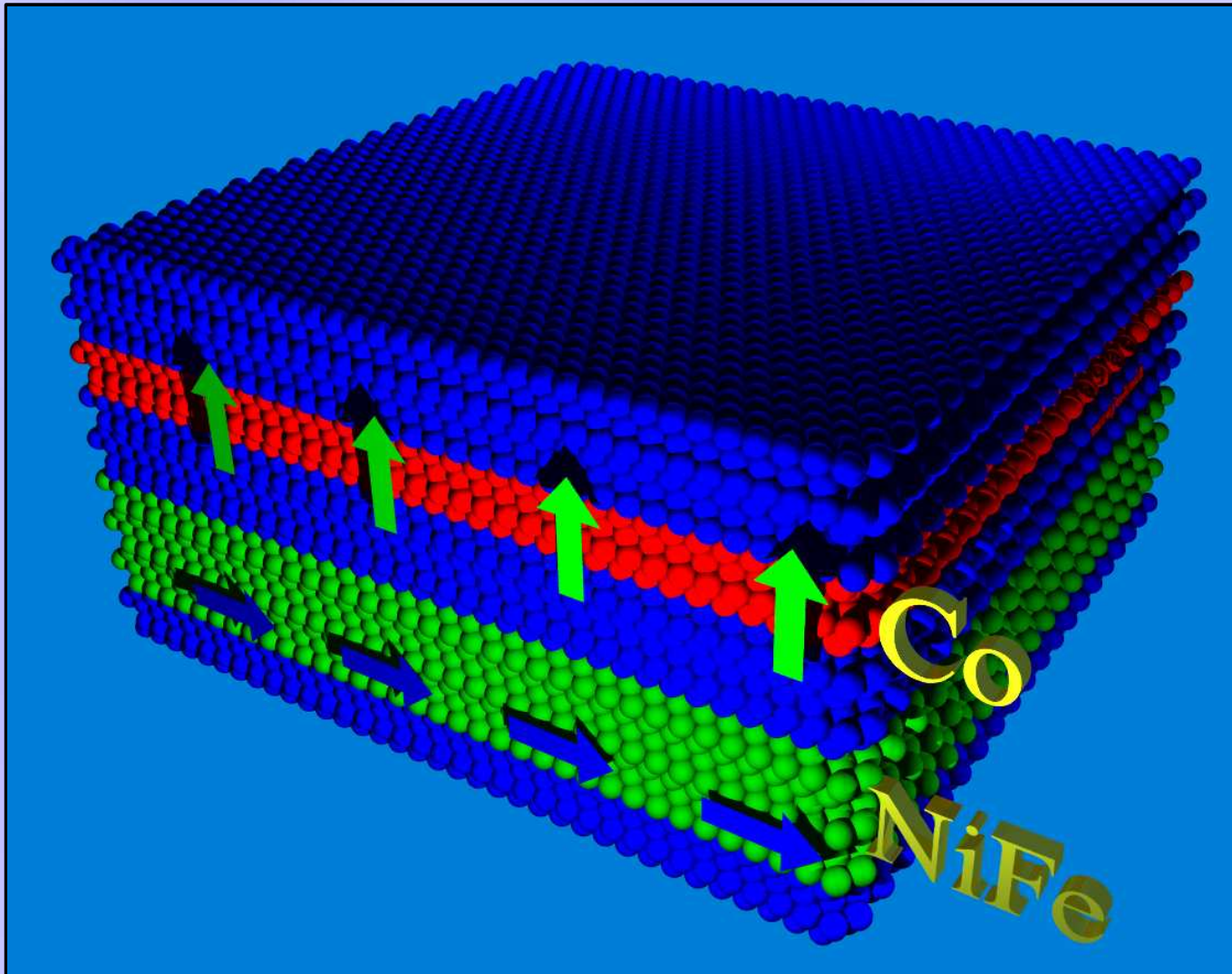
Magnetic properties



$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(1.9 \text{ nm})/\mathbf{\text{Co}(0.6 \text{ nm})}/\text{Au}(1.9 \text{ nm})]_{10}$

$[\mathbf{\text{Co}(0.6 \text{ nm})}/\text{Au}(4.4 \text{ nm})]_{15}$

NiFe sublayers do not considerably influence the reversal of Co sublayers.



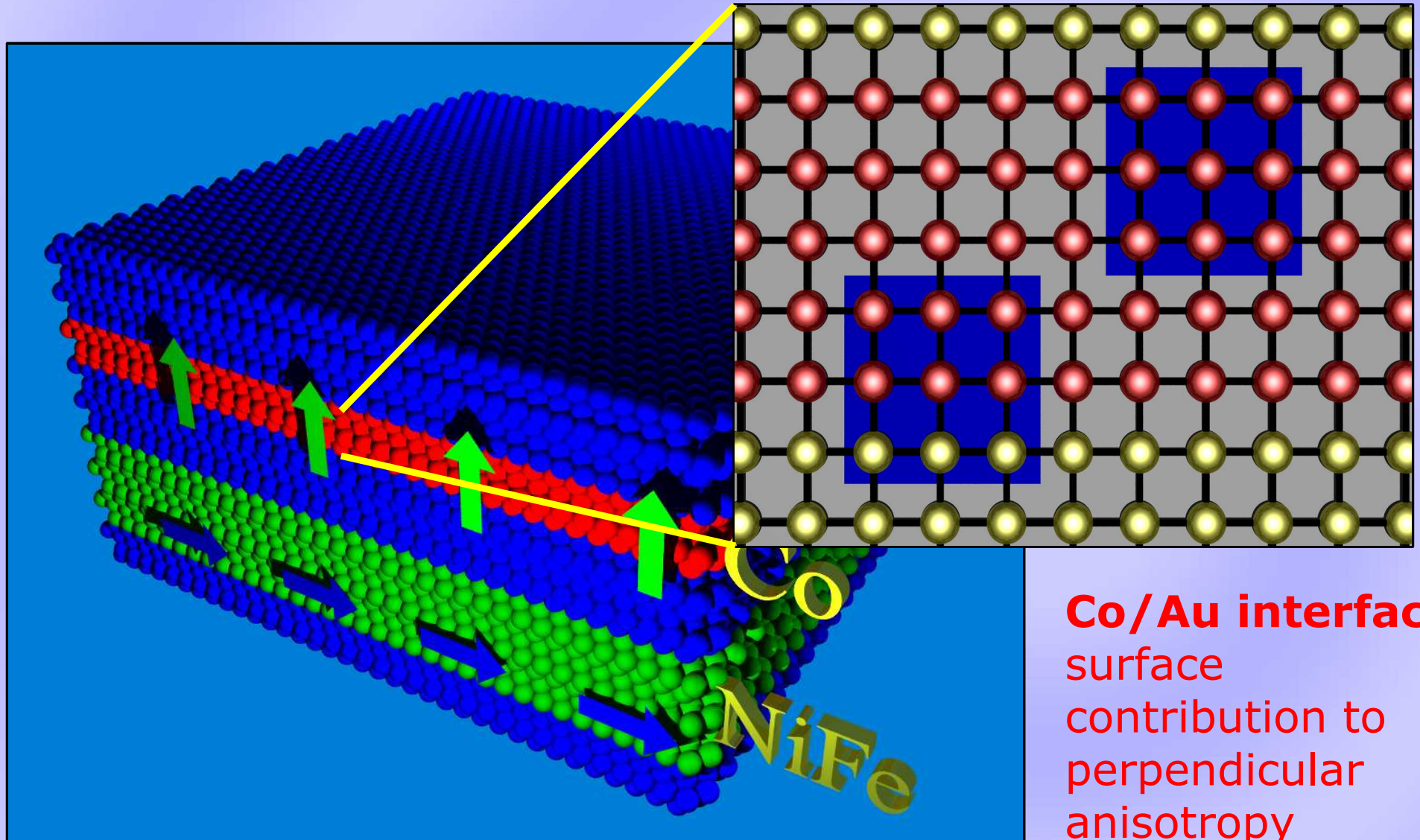
K_{eff} strongly
depends on t_{Co}

Co sublayers:
perpendicular
effective magnetic
anisotropy for
 $t_{\text{Co}} = 0.5 \div 1.2 \text{ nm}$

$$K_{\text{eff}} = \frac{2K_{1s}}{t_{\text{Co}}} + K_{1v} - \frac{1}{2} \mu_0 (M_S^{\text{Co}})^2$$

Magnetic properties

Stripe domains

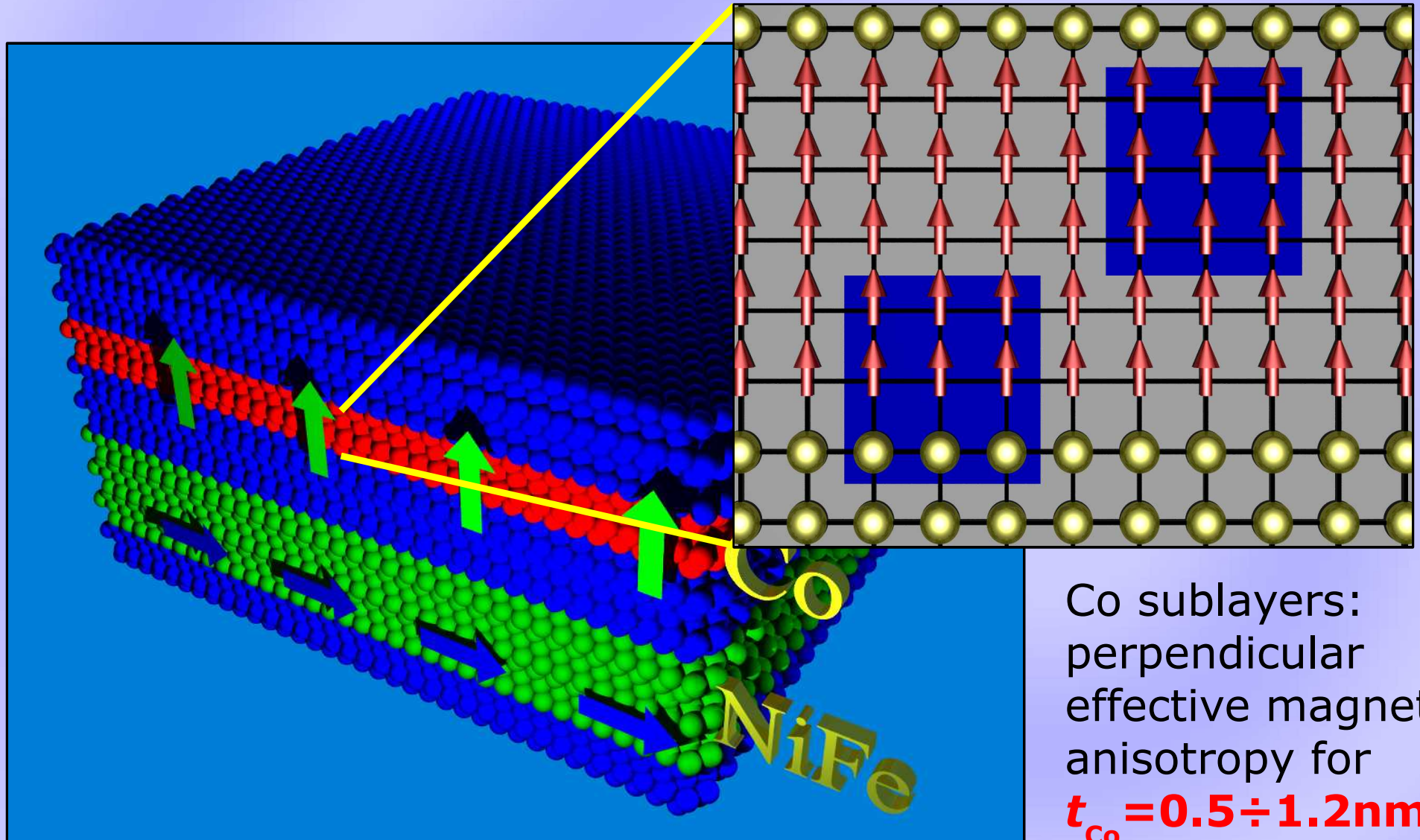


Co/Au interface-
surface
contribution to
perpendicular
anisotropy

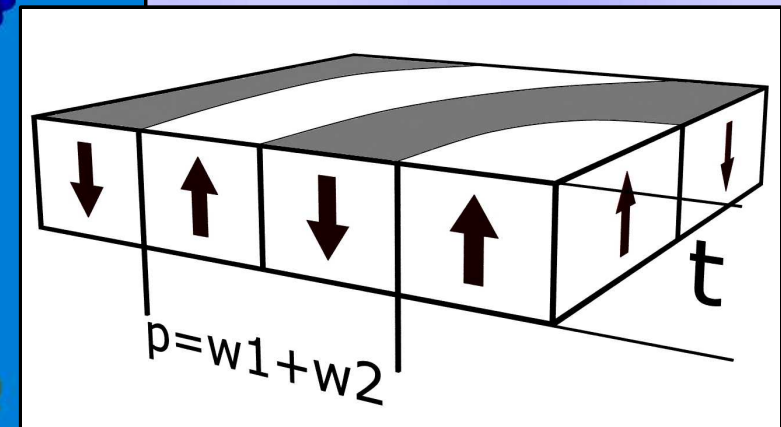
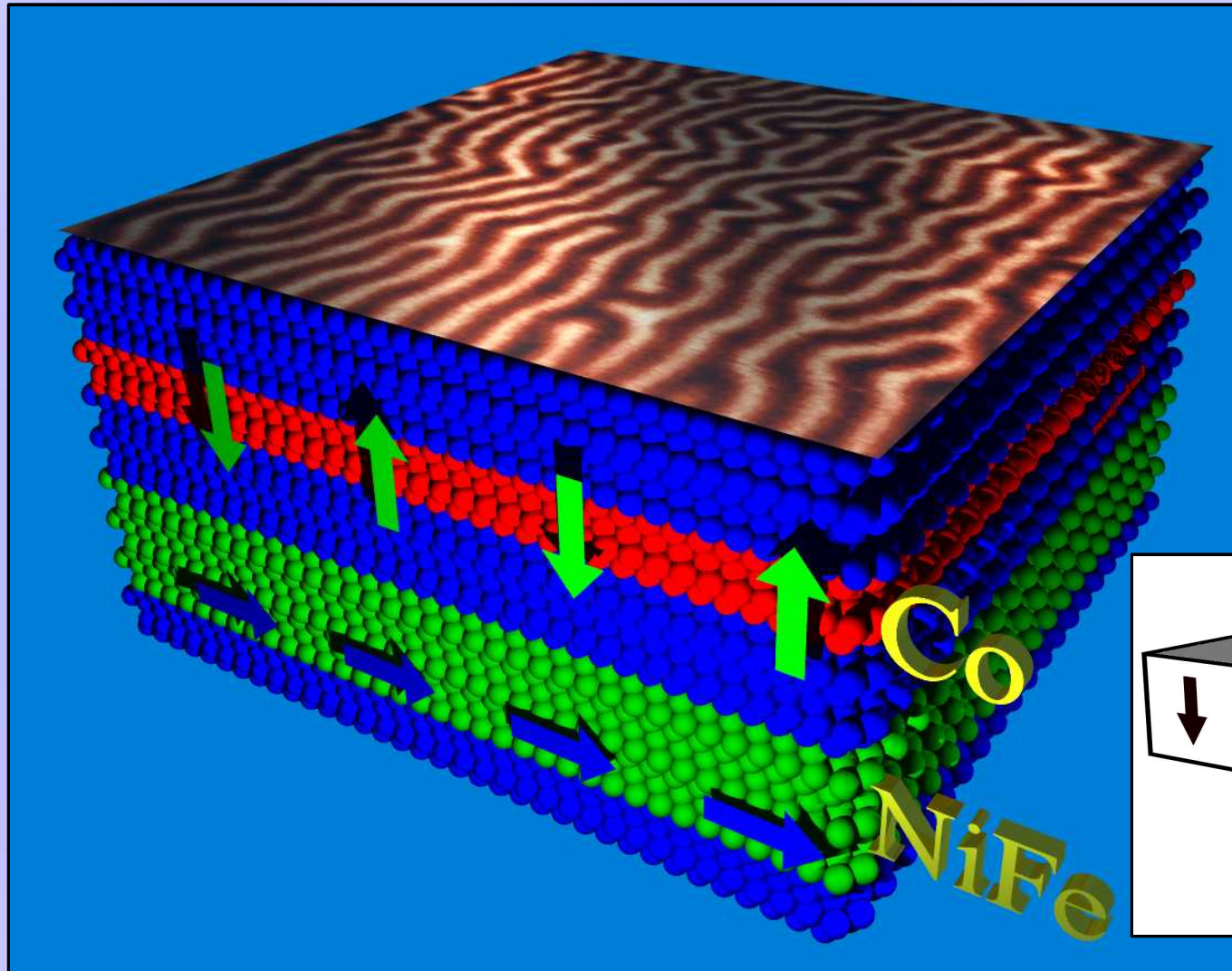
$$K_{eff} = \frac{2K_{1s}}{t_{Co}} + K_{1v} - \frac{1}{2}\mu_0(M_S^{Co})^2$$

Magnetic properties

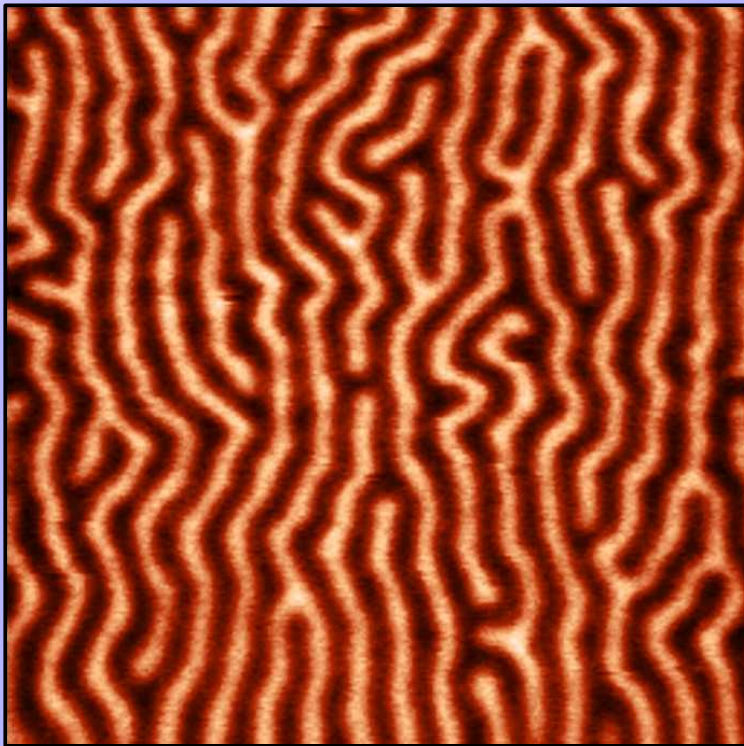
Stripe domains



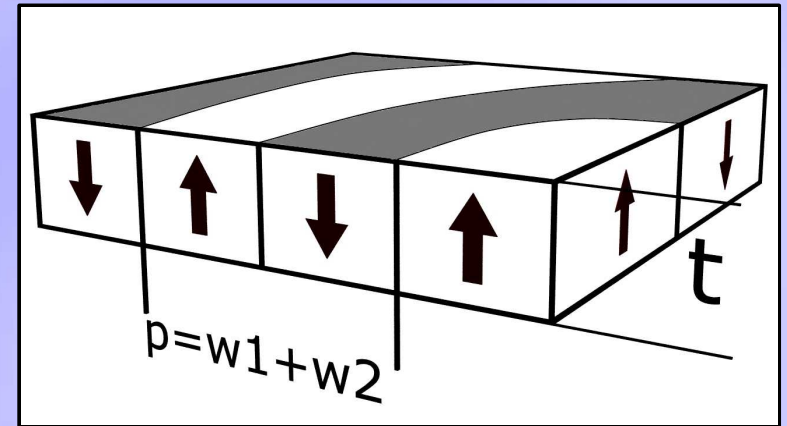
$$K_{eff} = \frac{2K_{1s}}{t_{Co}} + K_{1v} - \frac{1}{2} \mu_0 (M_S^{Co})^2$$



Magnetic Force Microscopy confirms the presence of **the stripe domain structure** characteristic for systems with perpendicular anisotropy.



AC demagnetization
 $5 \times 5 \mu\text{m}^2$

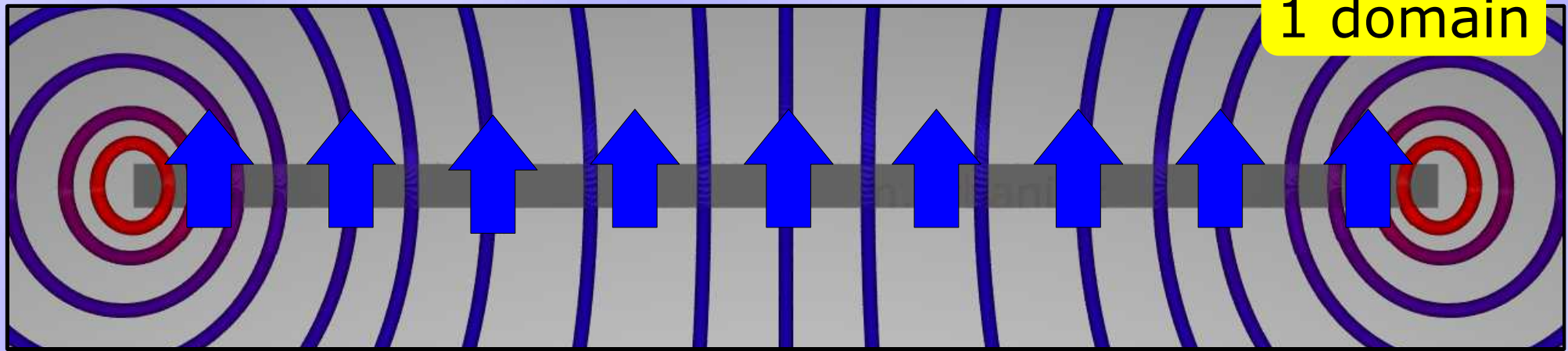


spatial period 400-1000 nm

Spatial period of the stripe domain structure depends strongly on the thicknesses of Co and Au sublayers.

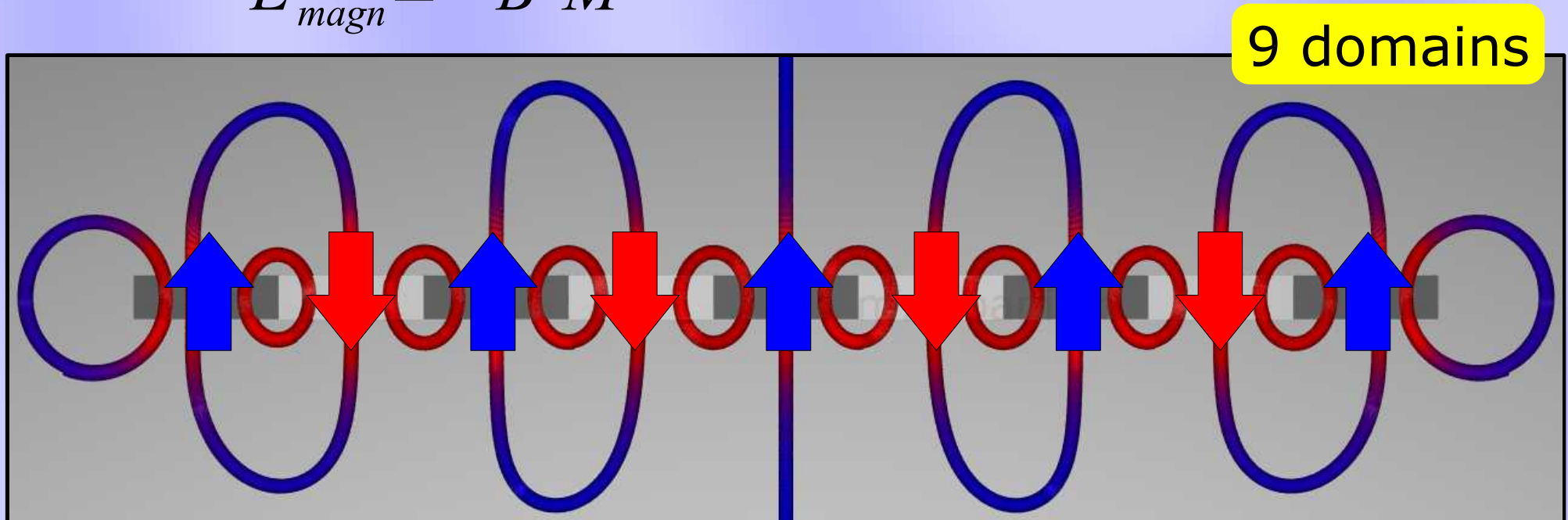
$[\text{Ni}_{80}\text{Fe}_{20}^*(2 \text{ nm})/\text{Au}(2.4 \text{ nm})/\text{Co}(1.2 \text{ nm})/\text{Au}(2.4 \text{ nm})]_{10}$

* with ^{57}Fe

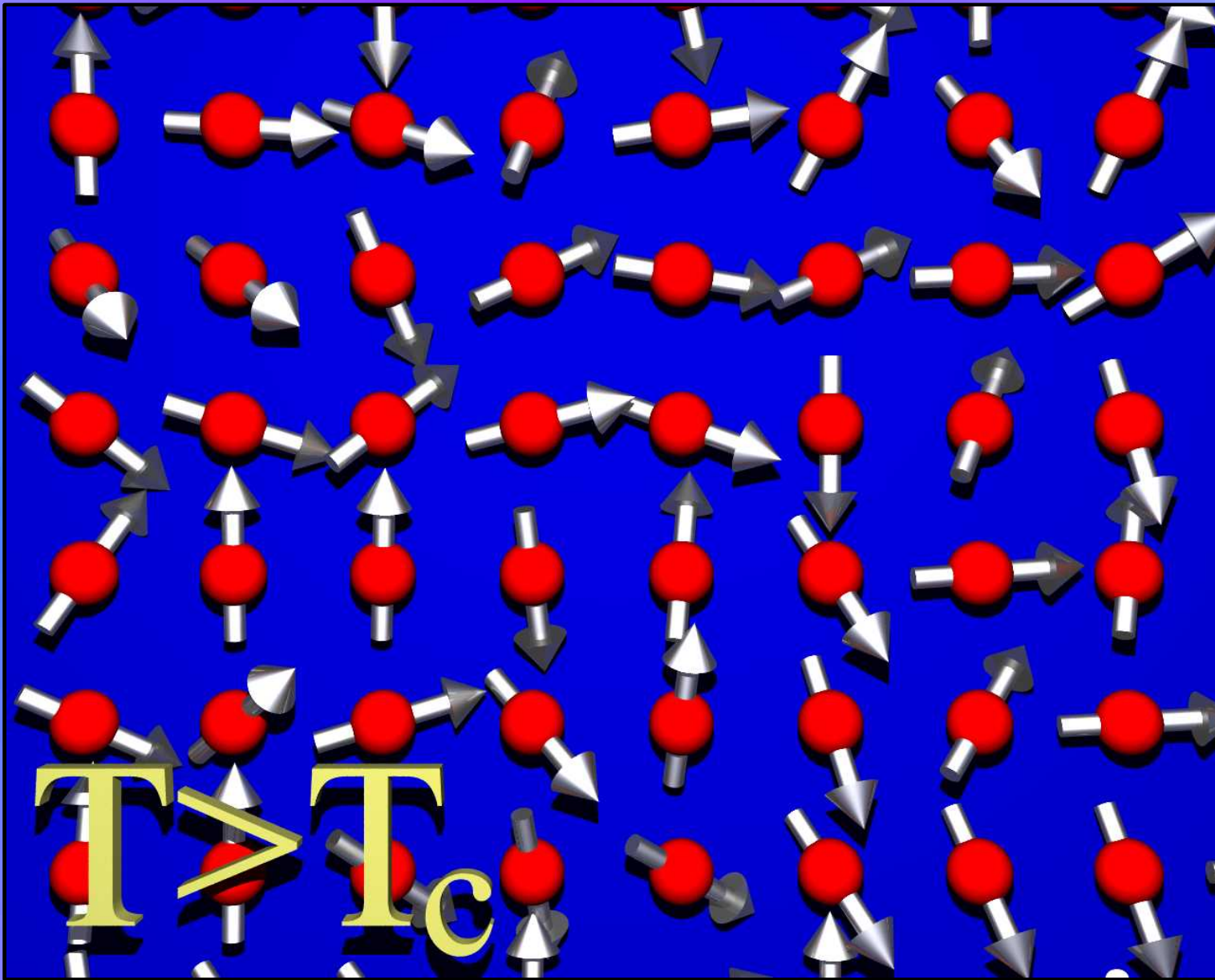


Division into **the magnetic stripe domains** increases the magnetic induction within the layer and leads to the decrease of magnetostatic energy:

$$E_{magn} = -\vec{B} \cdot \vec{M}$$



Giant magnetoresistance



T_{Curie} :

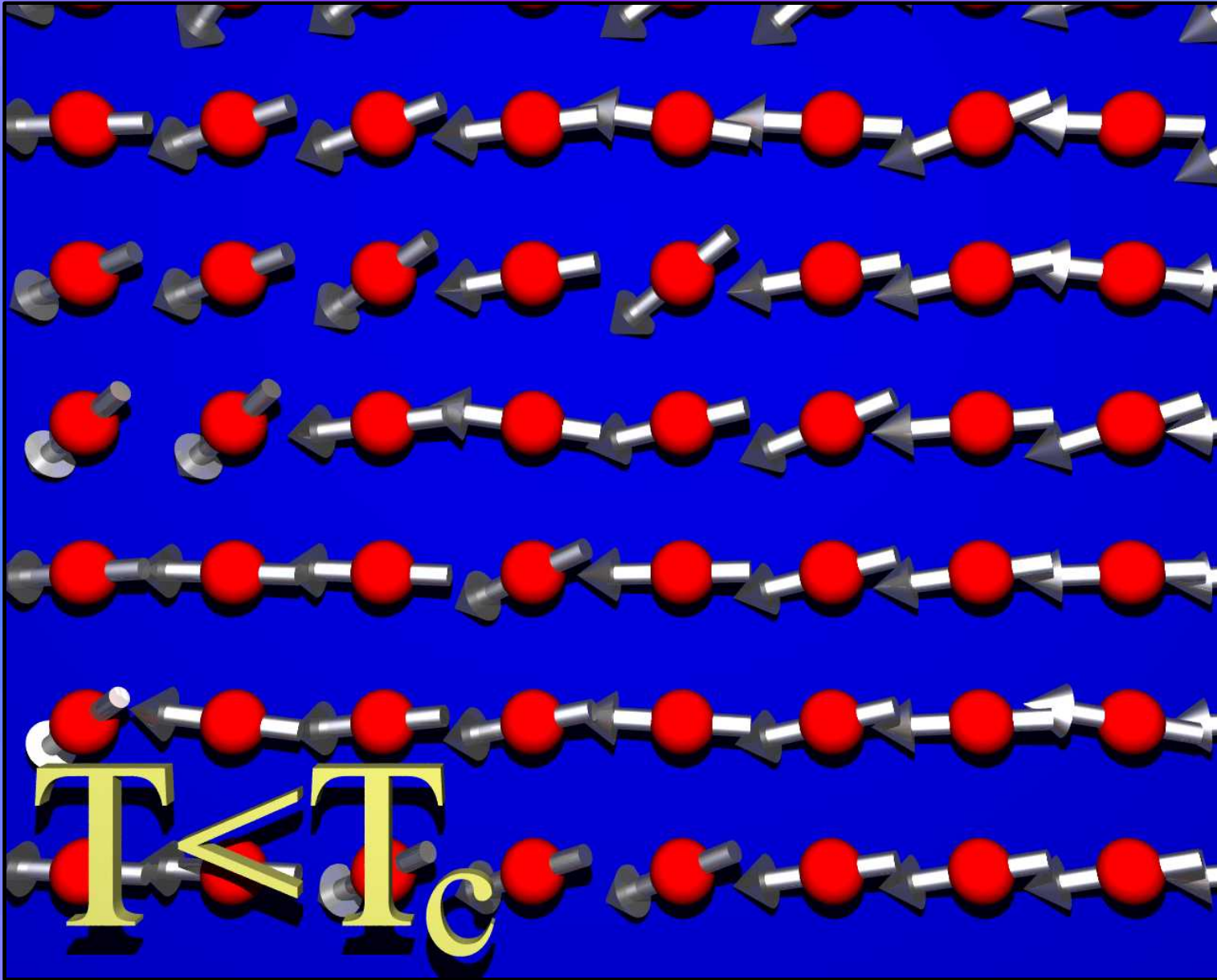
Fe 1044 K

Co 1388 K

Ni 627 K

NiFe 660 K

Giant magnetoresistance



$T_{\text{Curie}}:$

Fe 1044 K

Co 1388 K

Ni 627 K

NiFe 660 K

Giant magnetoresistance

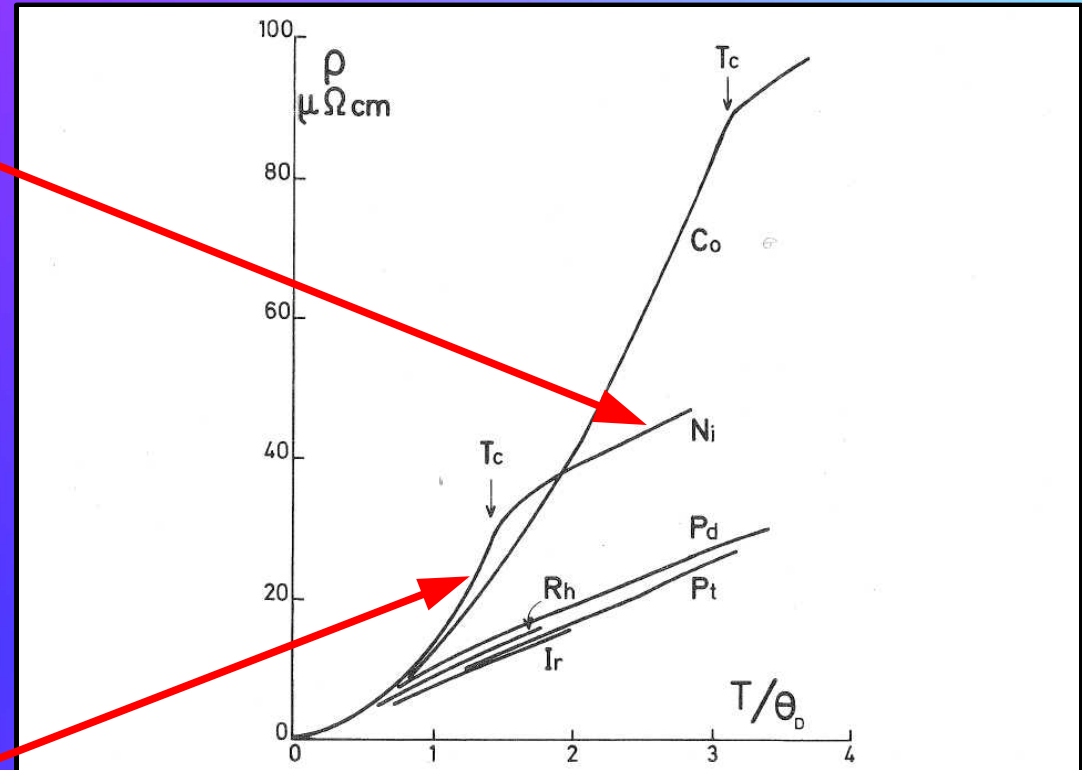
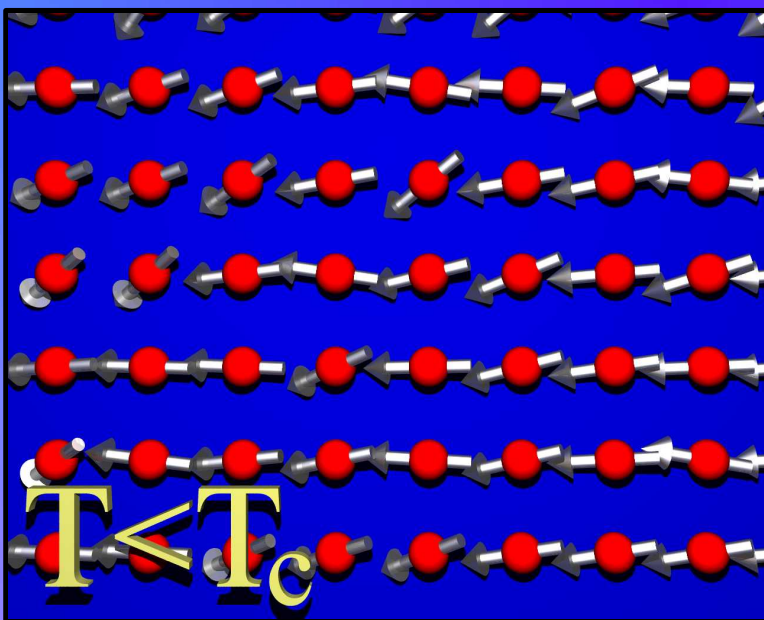
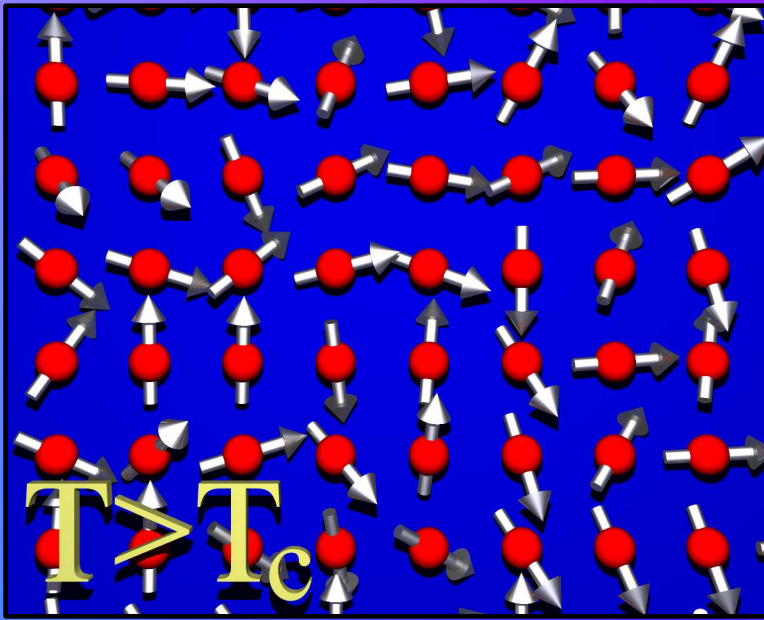
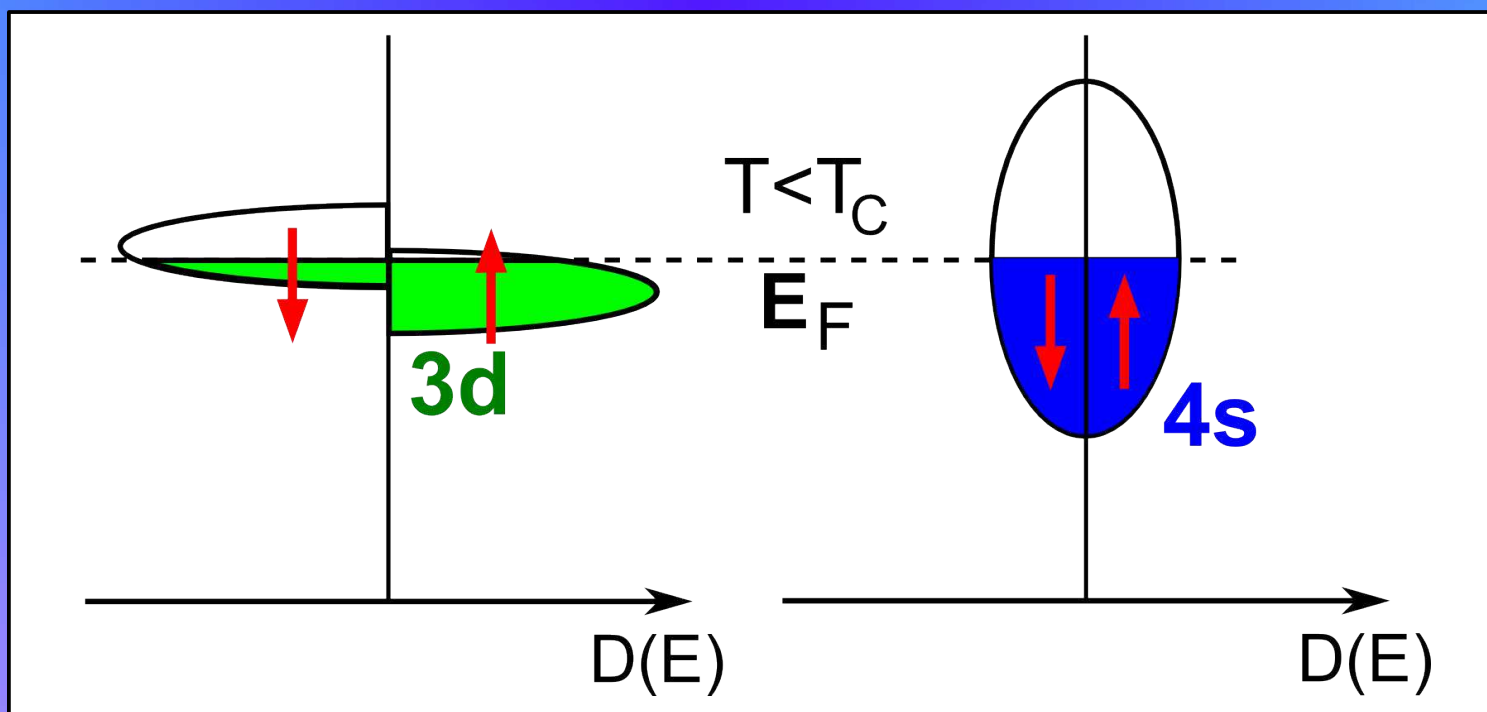
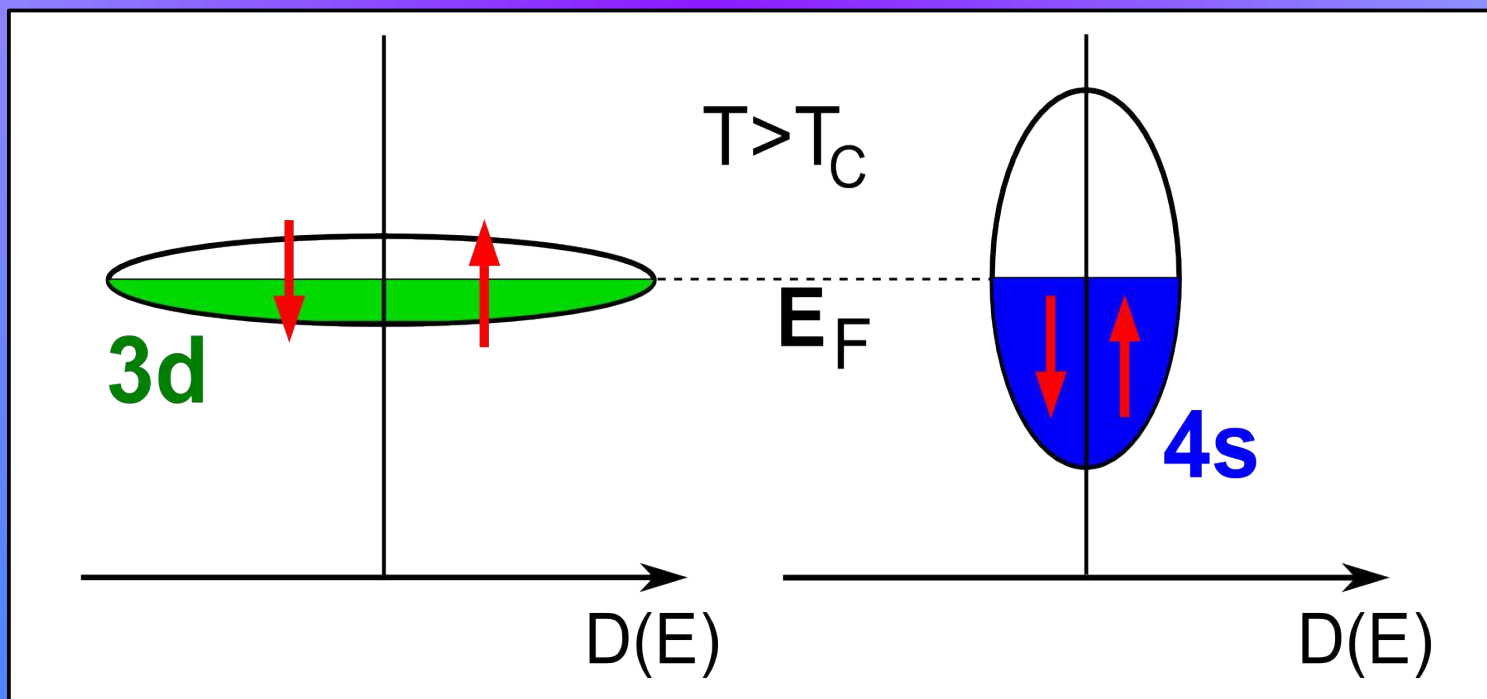


Fig. 10. Resistivity of several transition metals as a function of T/θ_D . θ_D is the Debye temperature.

Below Curie temperature the resistance of ferromagnetic materials decreases below that of non-ferromagnetic metals.

Giant magnetoresistance



Scattering:

$$s\uparrow \rightarrow d\uparrow$$

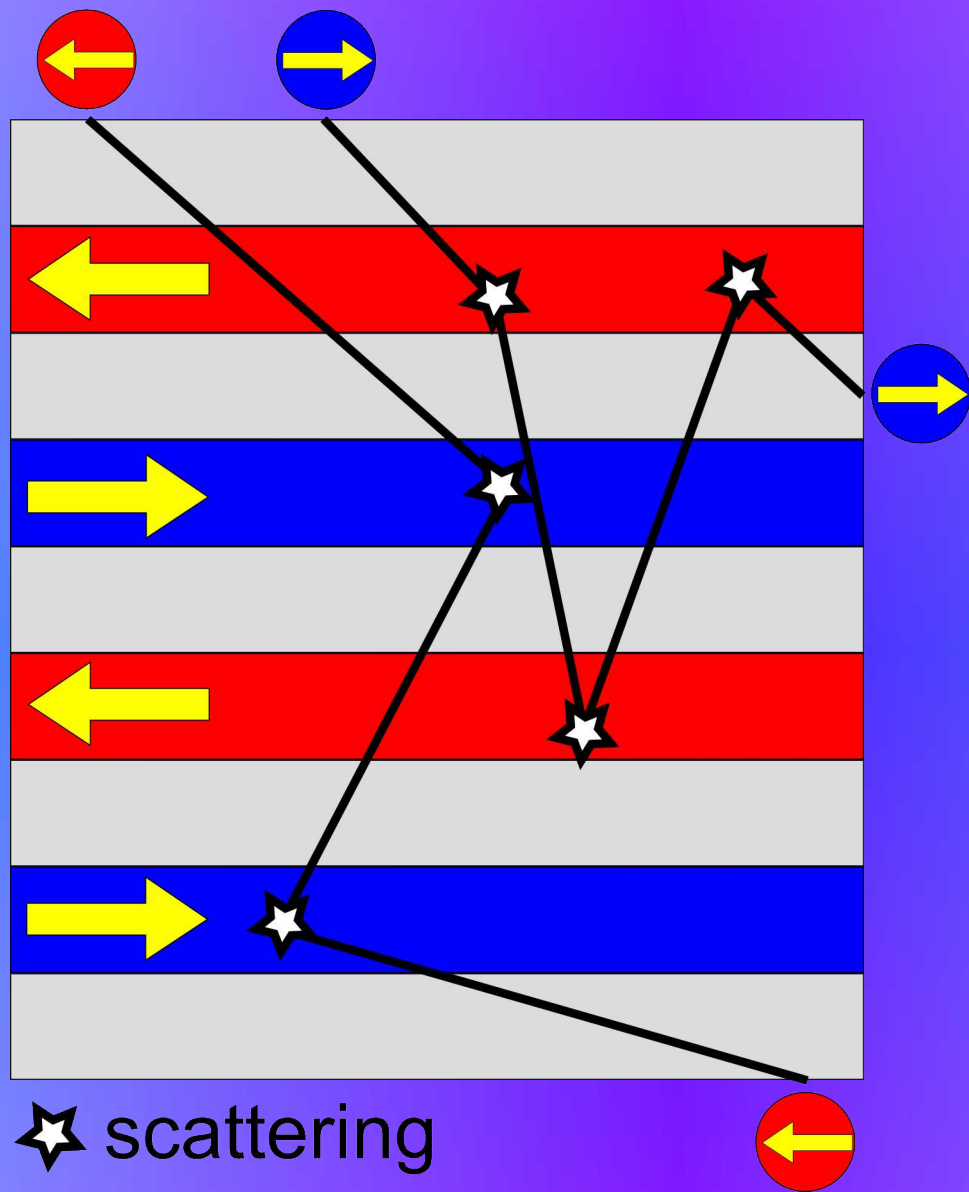
$$s\downarrow \rightarrow d\downarrow$$

$$\rho \sim D(E_F)$$

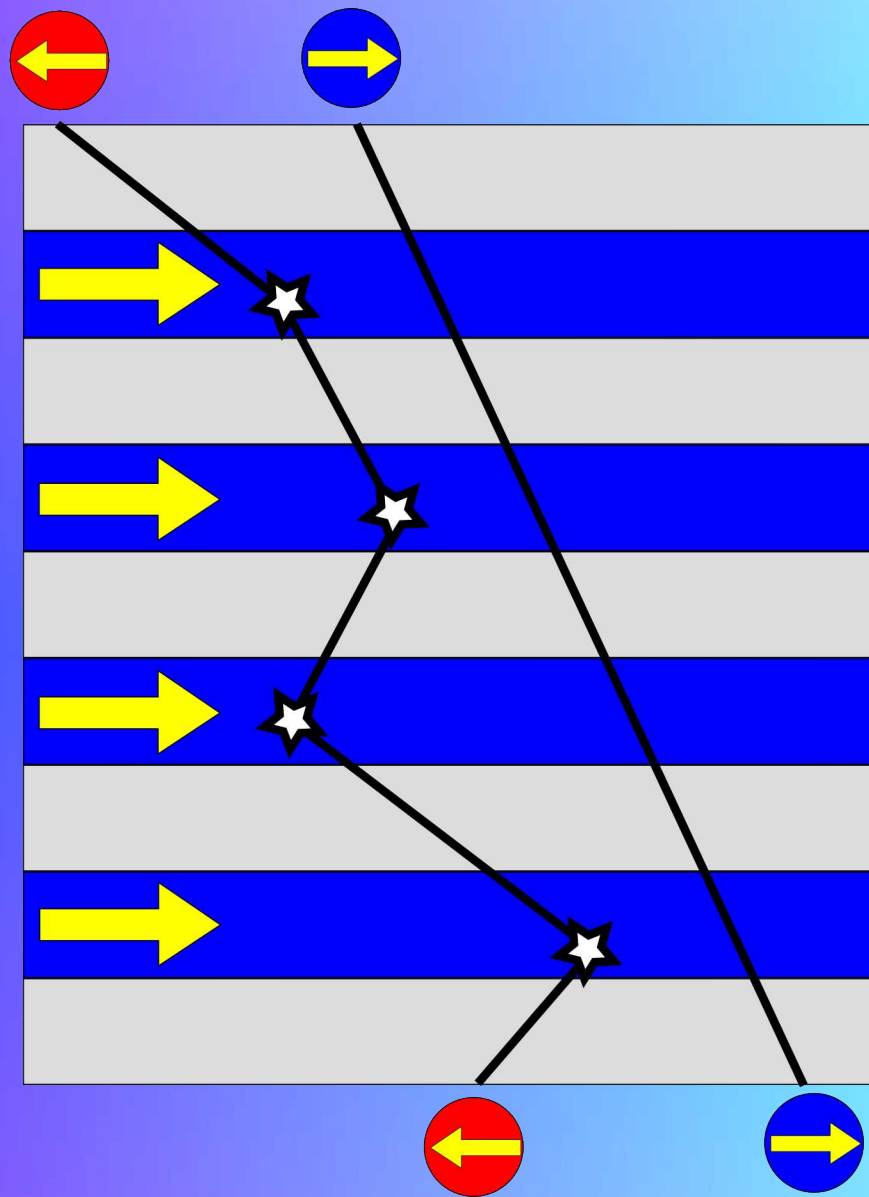
$$s\uparrow \rightarrow d\uparrow$$

$$s\downarrow \rightarrow d\downarrow$$

Giant magnetoresistance



Antiparallel – high resistance



Parallel – low resistance

Giant magnetoresistance

Nobel 2007 (Fert, Grünberg)

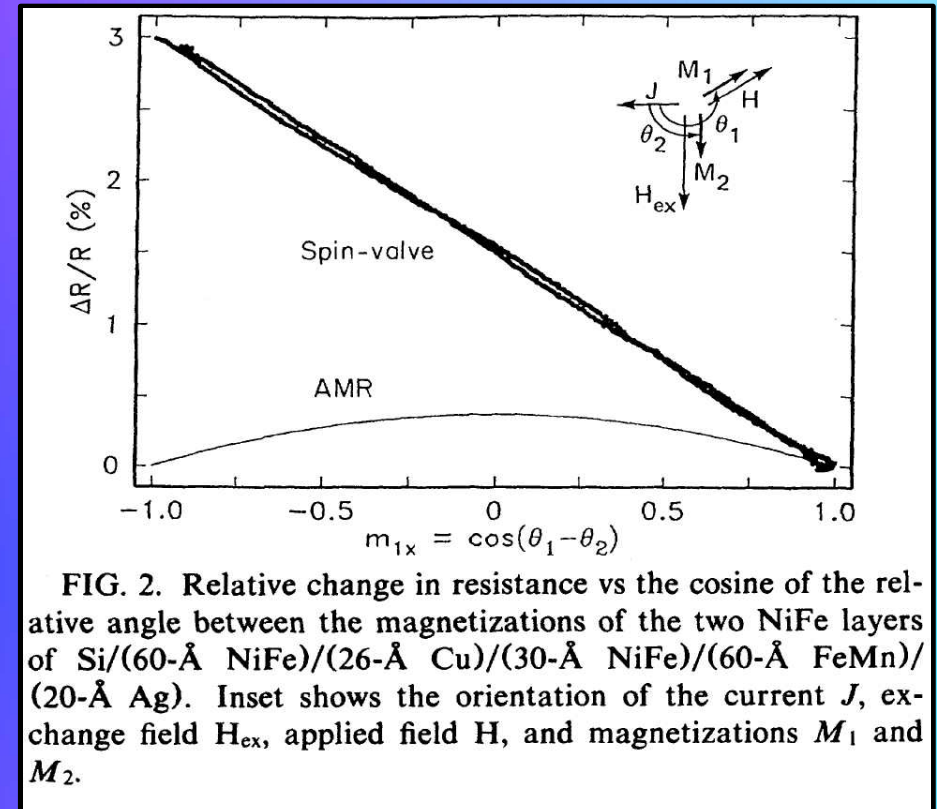
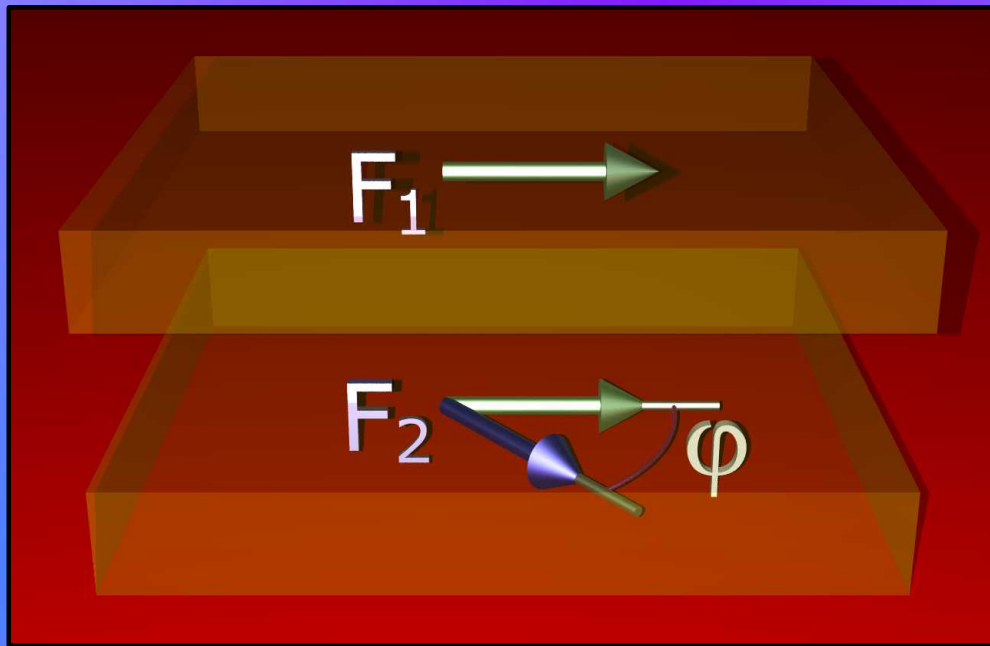
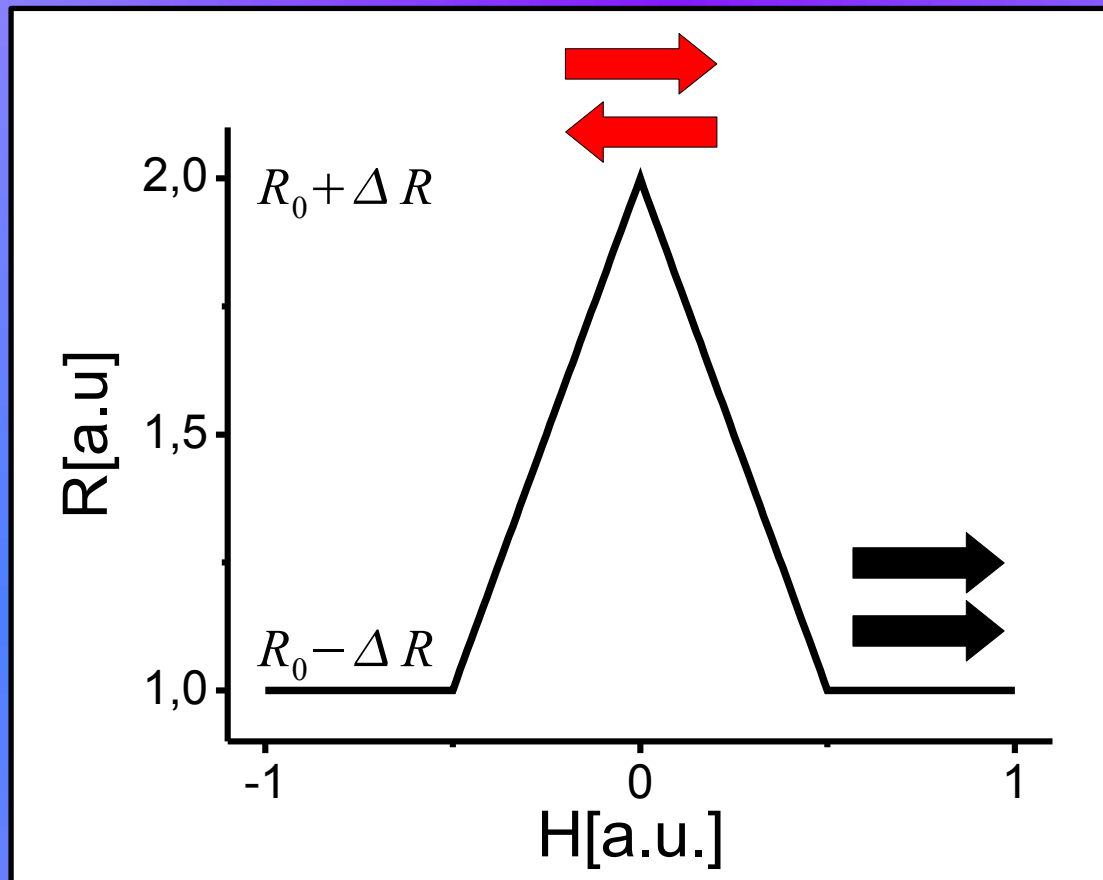


FIG. 2. Relative change in resistance vs the cosine of the relative angle between the magnetizations of the two NiFe layers of Si/(60-Å NiFe)/(26-Å Cu)/(30-Å NiFe)/(60-Å FeMn)/(20-Å Ag). Inset shows the orientation of the current J , exchange field H_{ex} , applied field H , and magnetizations M_1 and M_2 .

$$\Delta R \propto \cos(\varphi)$$

Giant magnetoresistance

Nobel 2007 (Fert, Grünberg)

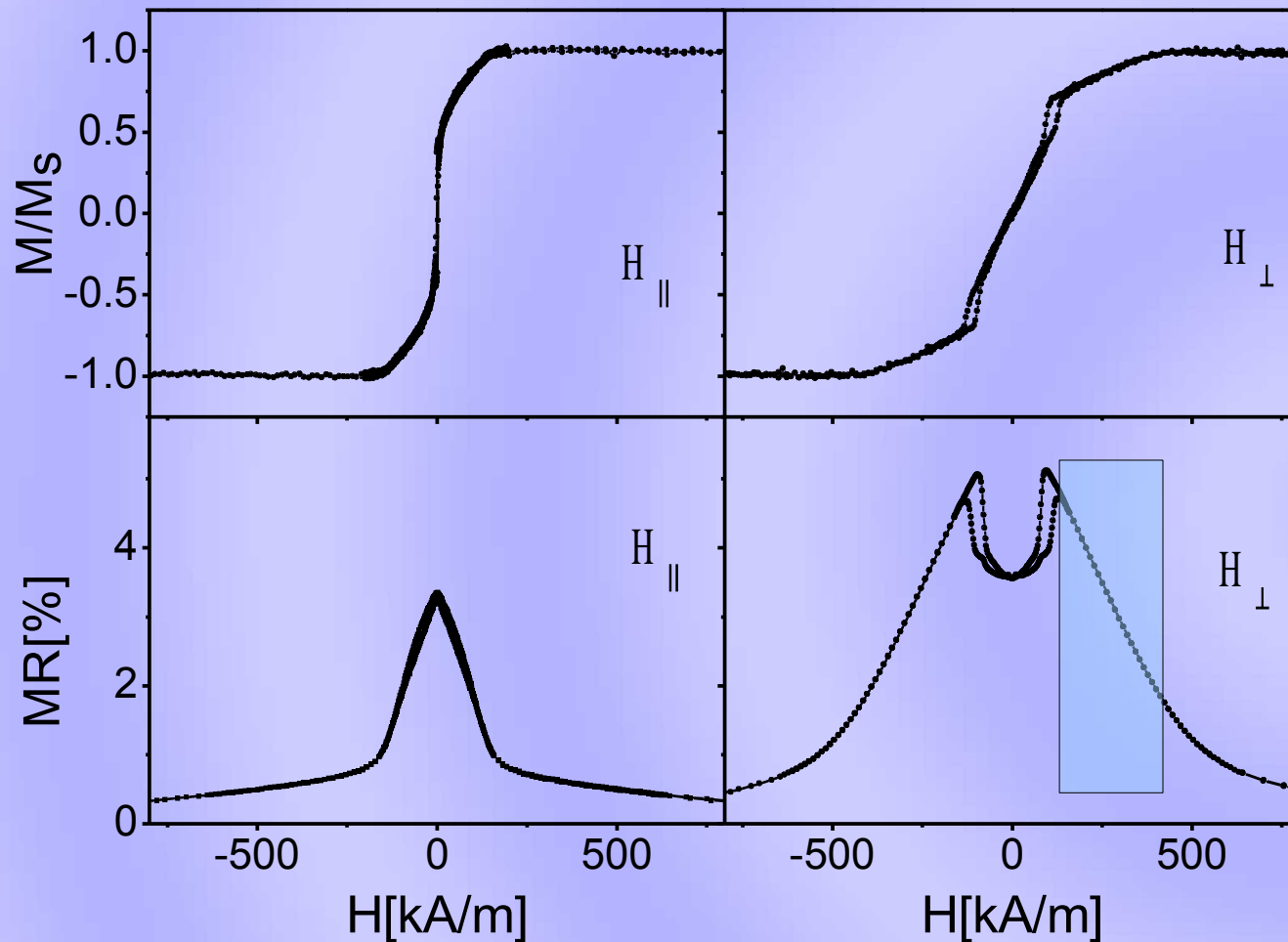


$$2\Delta R / (R_0 - \Delta R) = 1 \div 100 \%$$

Co/Au,
NiFe/Au,
NiFe/Cu,
Fe/Au,.....

$$R = R_0 - \Delta R \cos(\varphi)$$

Giant magnetoresistance



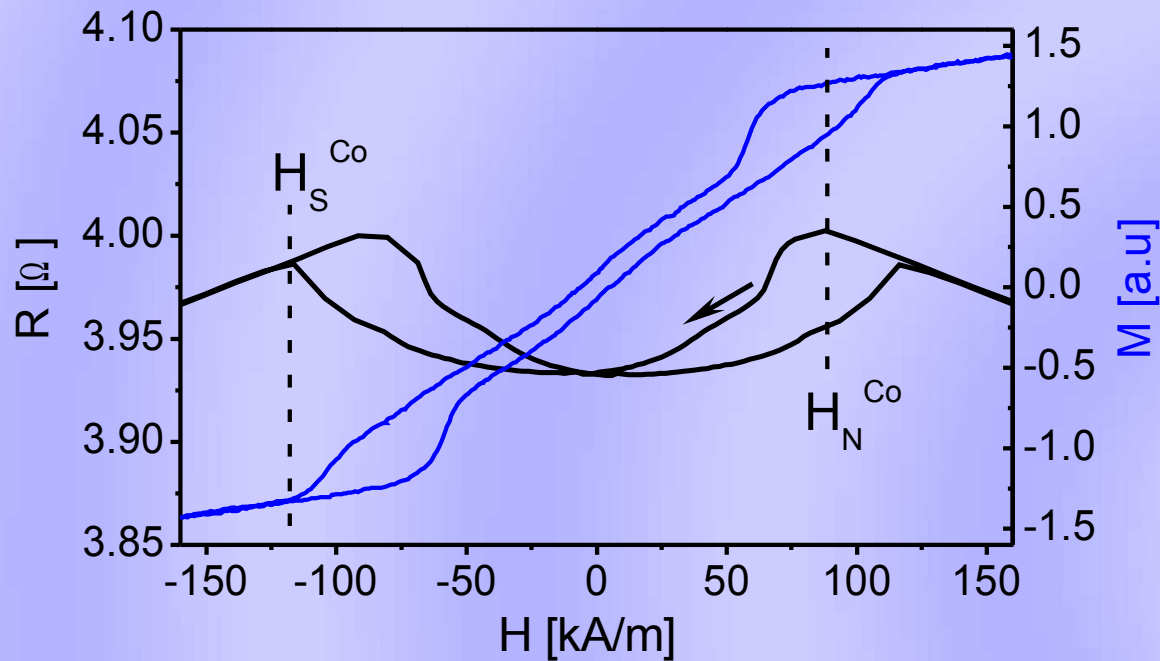
Broad **linearity** range in $R(H)$ dependence:

-magnetic layer magnetized along hard axis

-no hysteresis in linear range

$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(1.9 \text{ nm})/\text{Co}(1 \text{ nm})/\text{Au}(1.9 \text{ nm})]_{10}$

Giant magnetoresistance

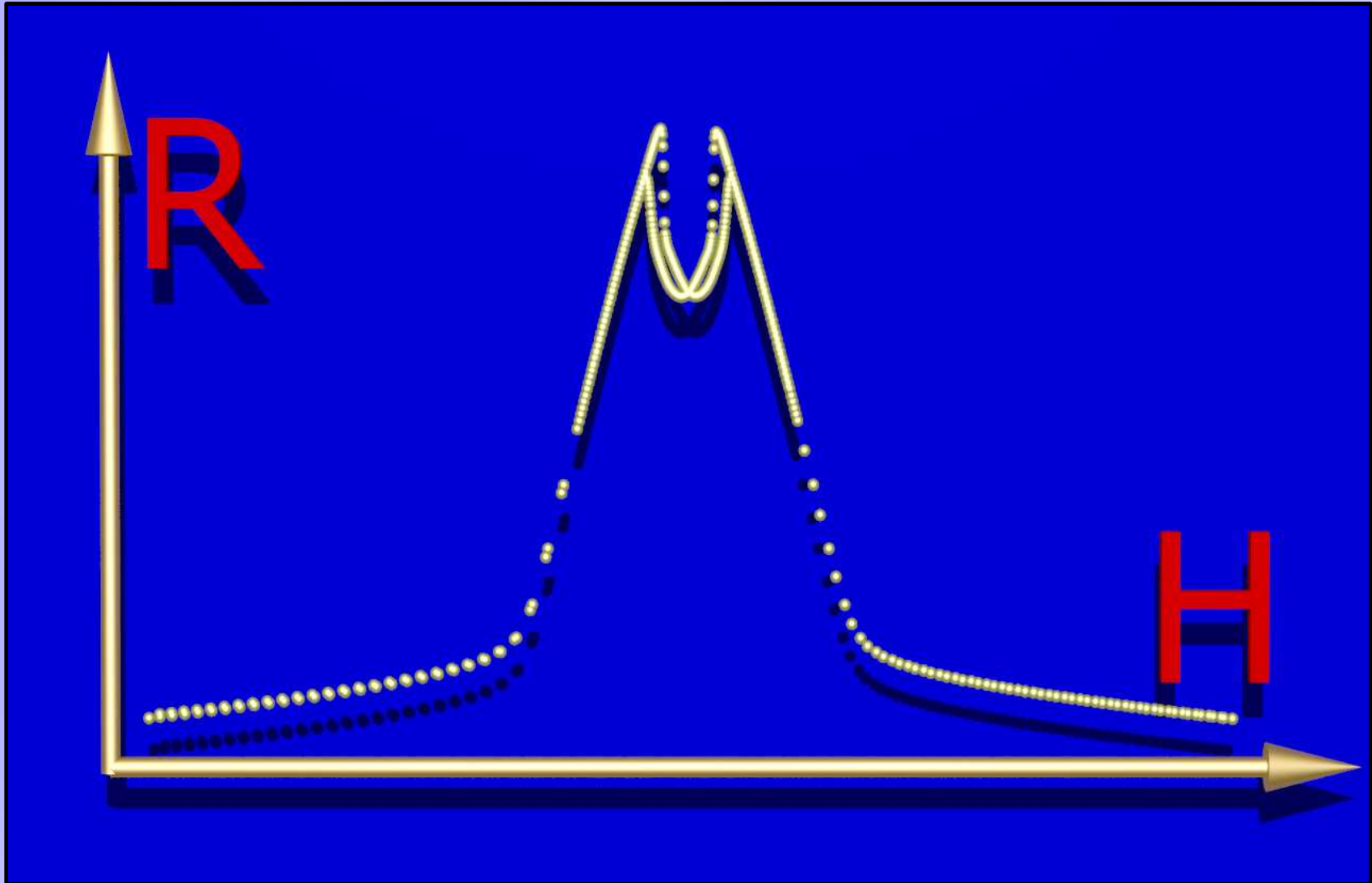


There is a local minimum of resistance in the $R(H)$ dependence.

The nucleation field (creation of the domain structure) and the annihilation field (saturation of Co layers) are visible both in $R(H)$ and $M(H)$ dependencies.

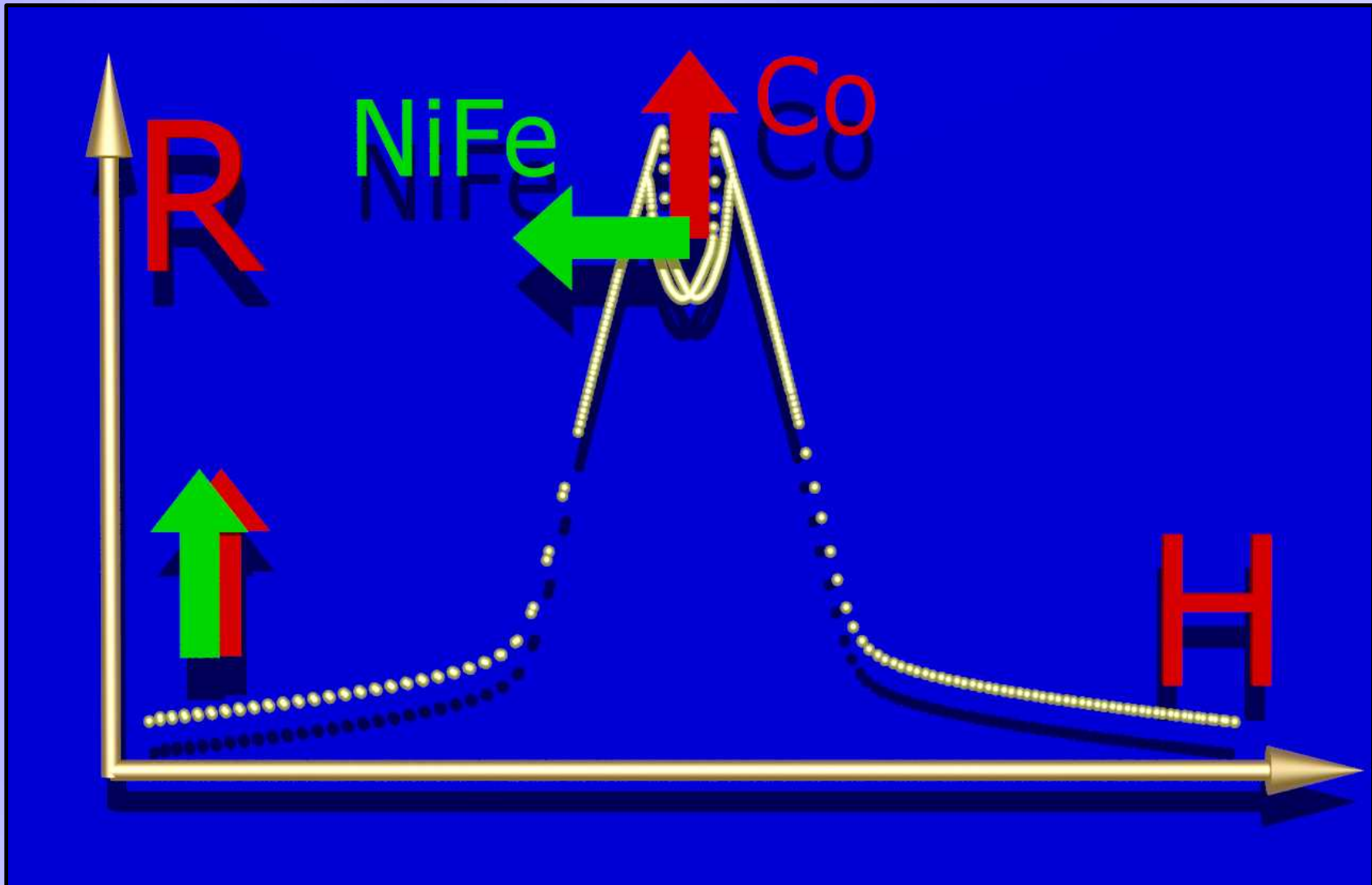
Giant magnetoresistance

Explaining the $R(H)$ dependence



Giant magnetoresistance

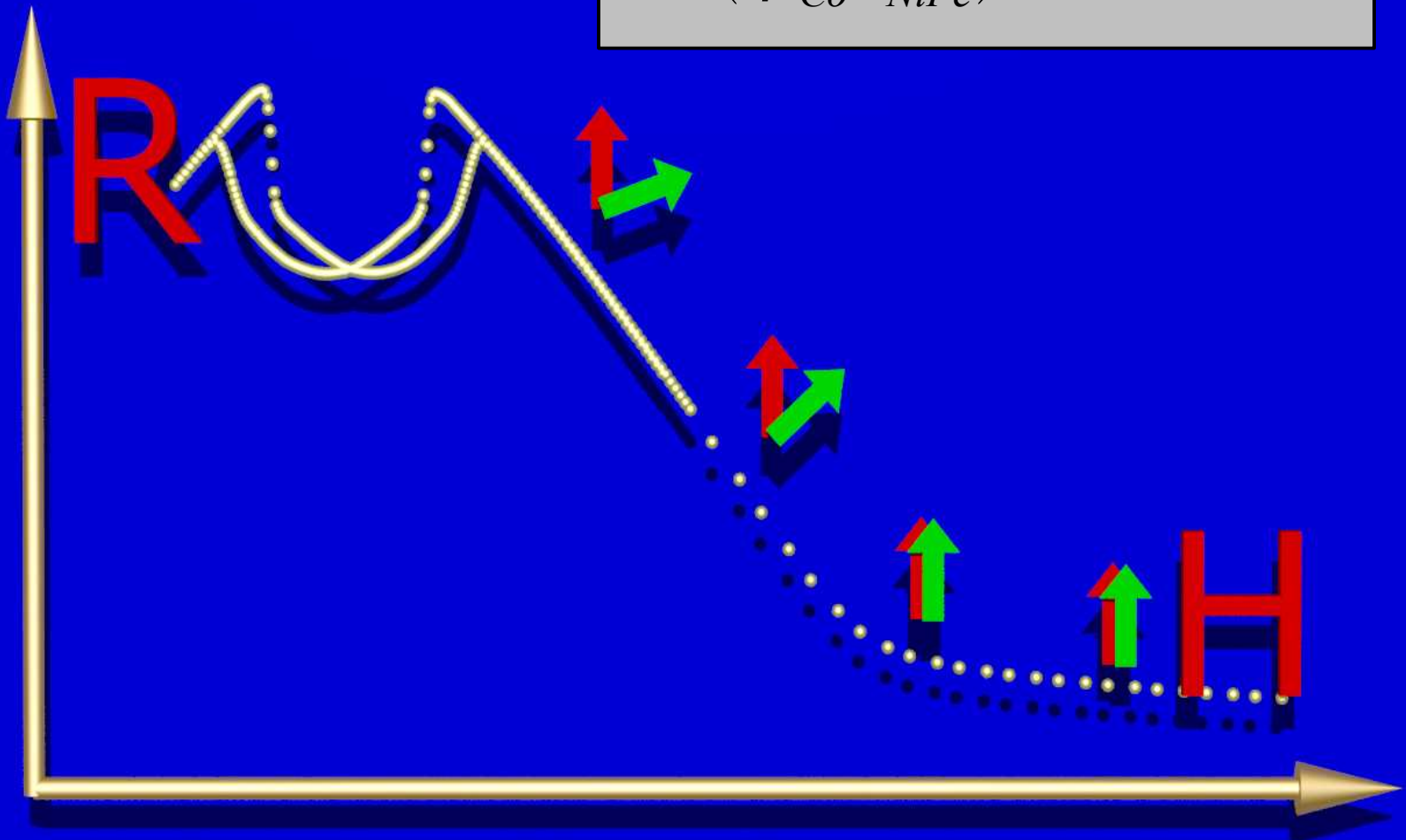
Explaining the $R(H)$ dependence



Giant magnetoresistance

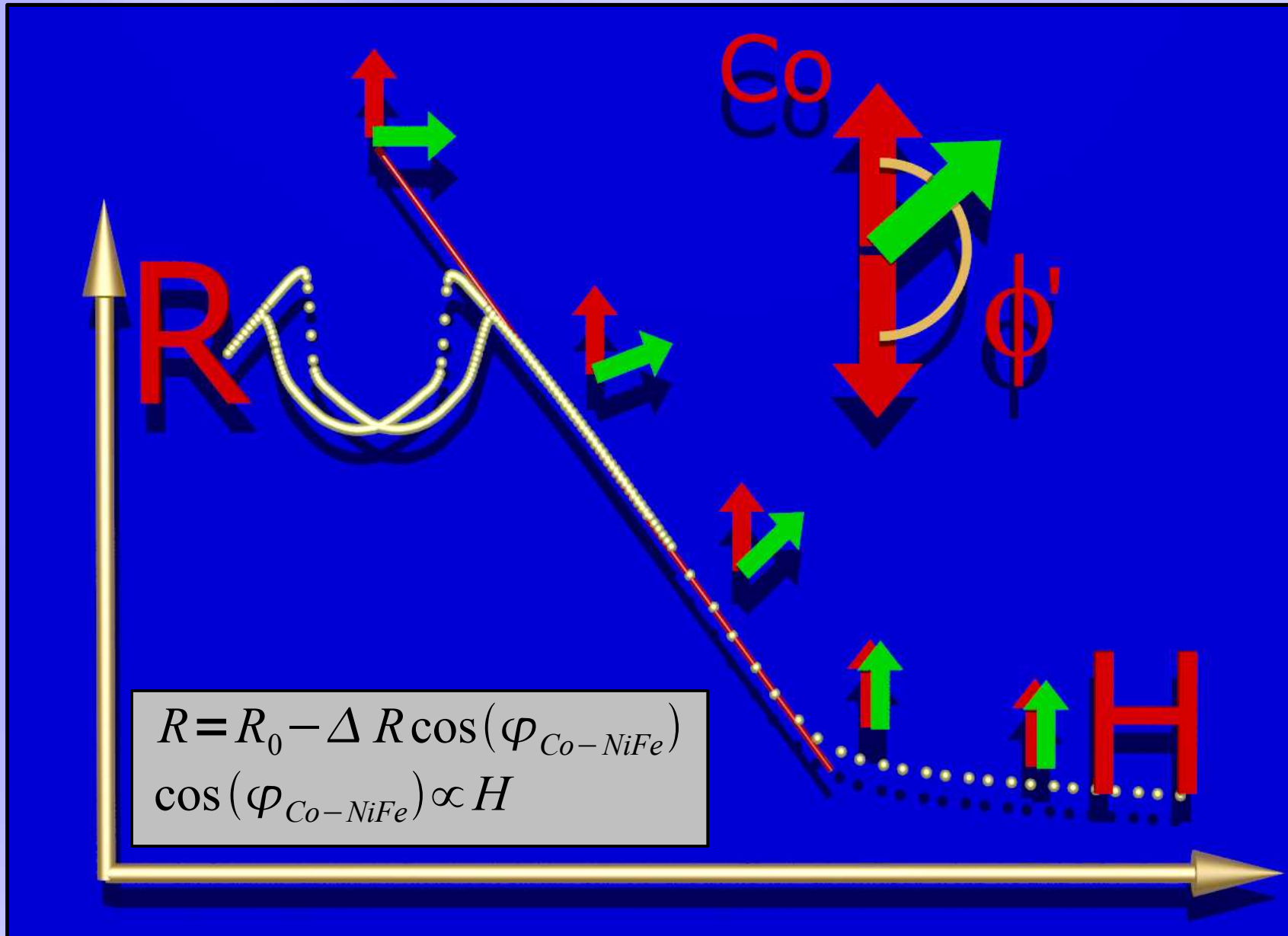
Explaining the $R(H)$ dependence

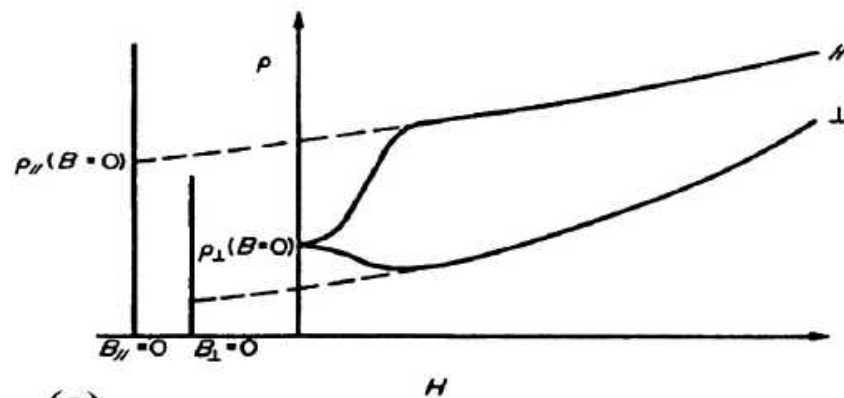
$$R = R_0 - \Delta R \cos(\varphi_{Co-NiFe})$$
$$\cos(\varphi_{Co-NiFe}) \propto H$$



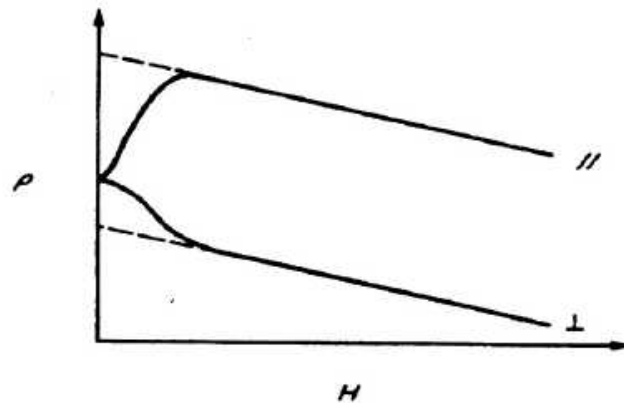
Giant magnetoresistance

Explaining the $R(H)$ dependence





(a)

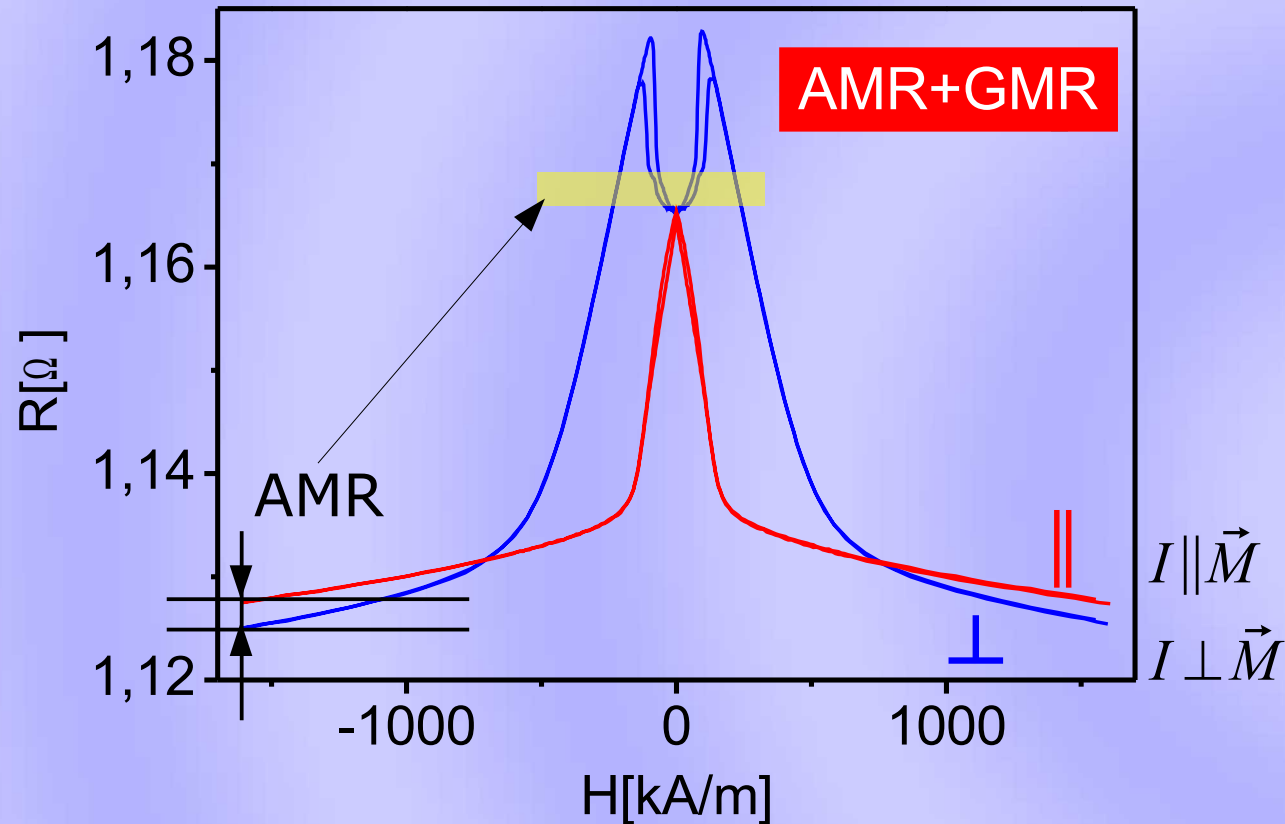


(b)

Abbildung 2.3: Schematische Darstellung der Magnetfeldabhängigkeit des spezifischen Widerstands in einem ferromagnetischen Metall. Im Niederfeldbereich dominiert der AMR, wogegen oberhalb des Sättigungsmagnetfeldes der positive Magnetwiderstand (a) oder der negative Magnetwiderstand durch Reduzierung der Spinunordnung (b) dominiert. Die gestrichelten Linien deuten an, wie man durch Rückextrapolation den positiven oder negativen Magnetwiderstandsbeitrag bei kleinen Feldern eliminieren kann, um den reinen AMR-Effekt zu erhalten.

AMR in NiFe alloys reaches 5%.

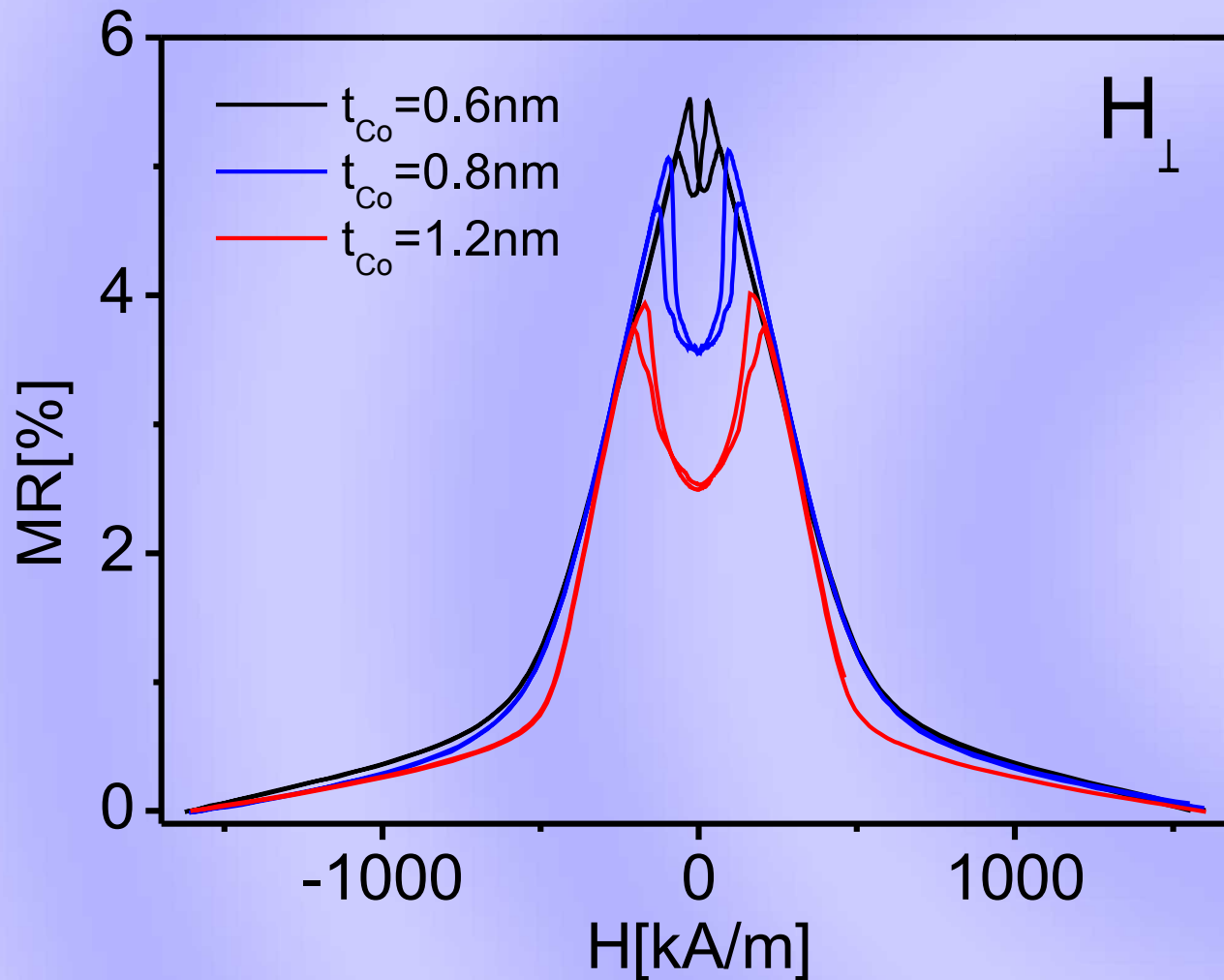
resistance of the system depends on **the angle** between the measuring current and the local magnetic moment



Anisotropic magnetoresistance is too small to account for the observed local minima of resistance.



Giant magnetoresistance



A relative "depth" of resistance minimum is a strong function of Co layers thickness.

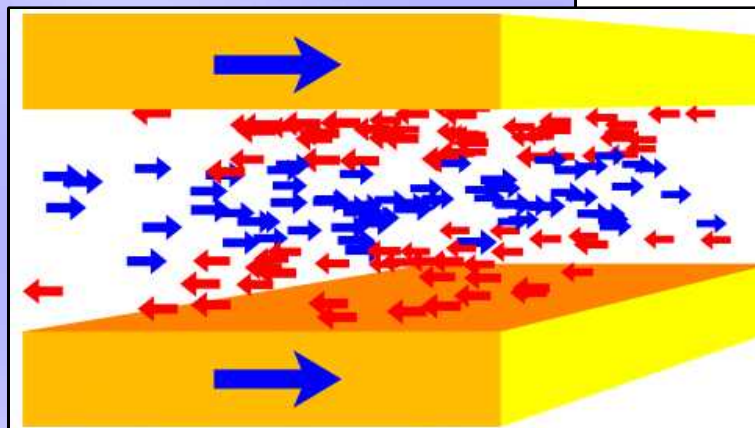
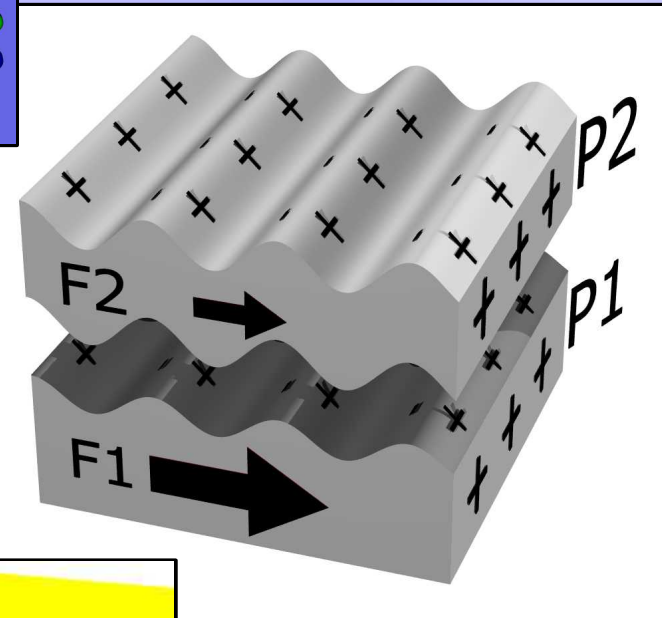
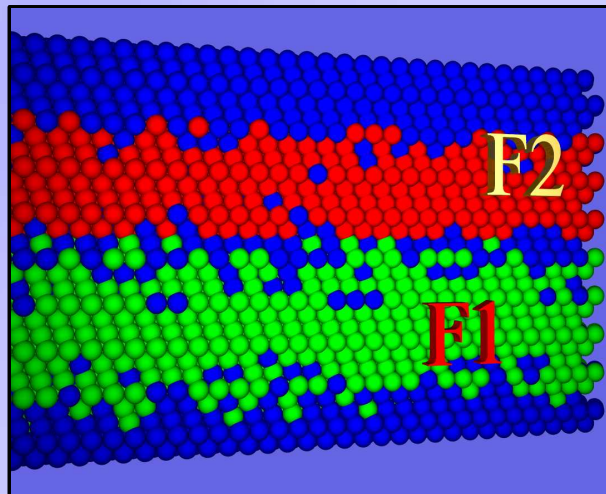
?

Magnetostatic coupling

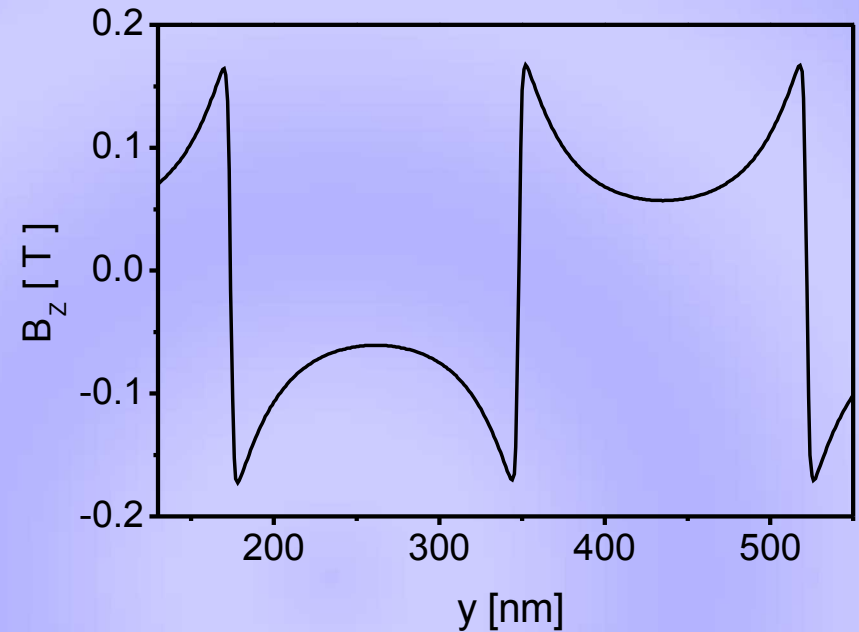
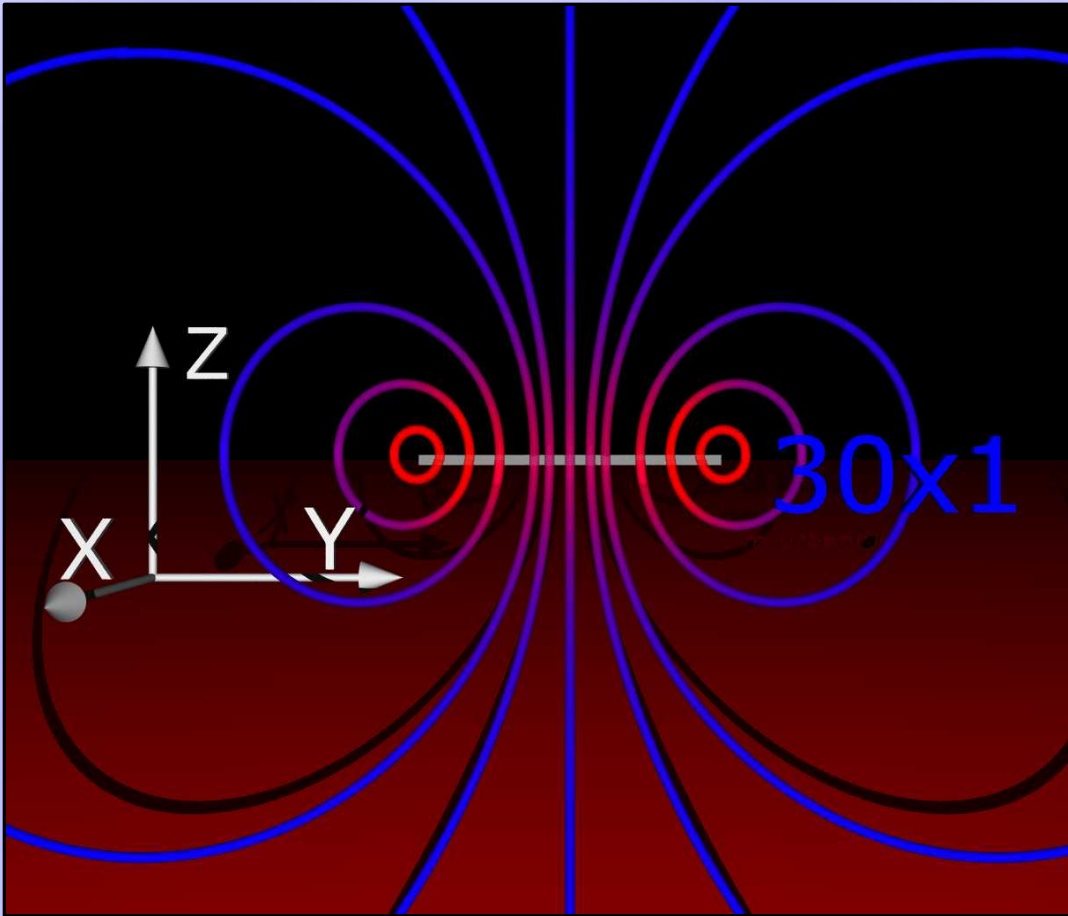
$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(1.9 \text{ nm})/\text{Co}(t_{Co})/\text{Au}(1.9 \text{ nm})]_{10}$

Interlayer coupling in magnetic multilayers

- coupling through magnetic bridging
- magnetostatic coupling
- Ruderman-Kittel-Kasuya-Yosida like coupling



Magnetostatic coupling



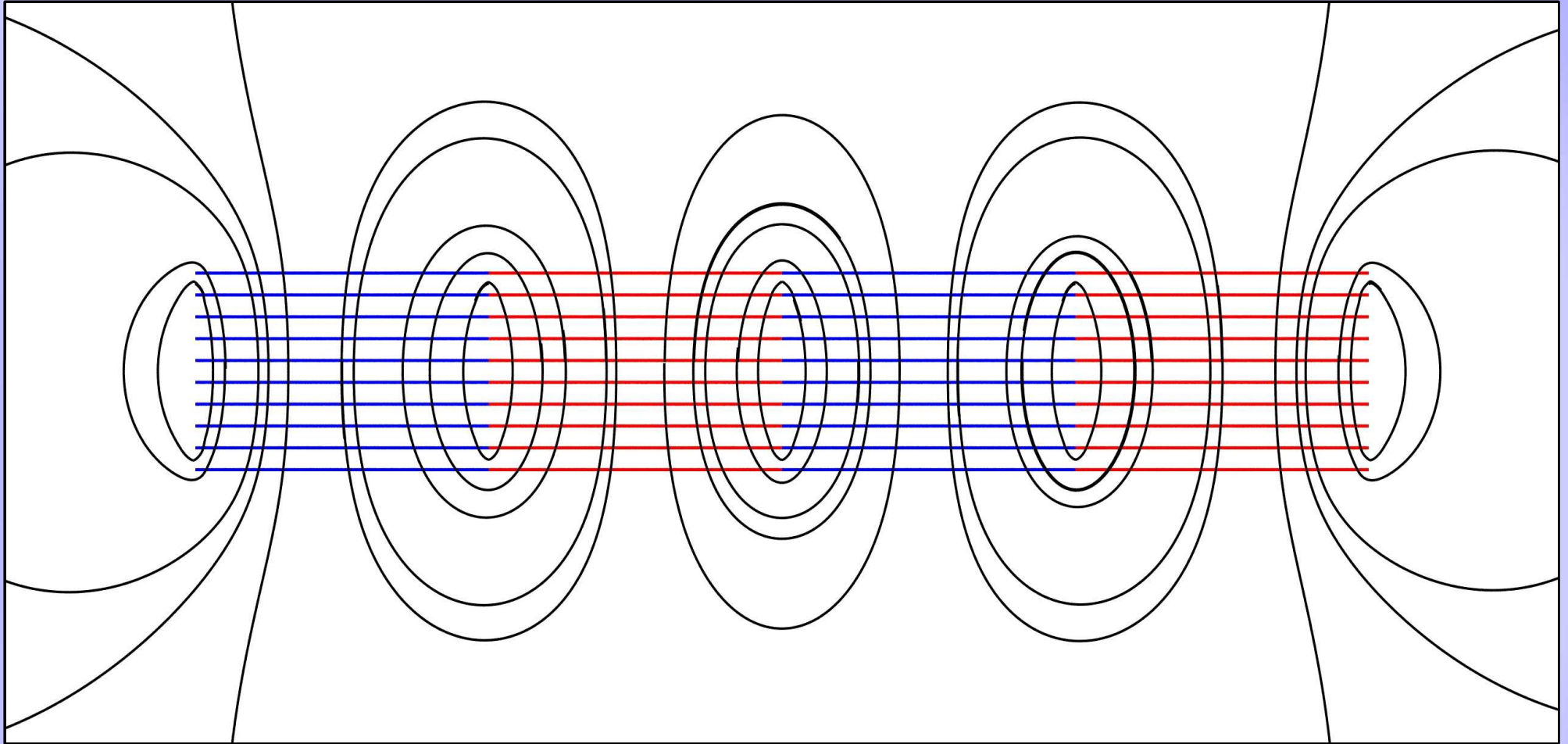
Model:

Domains width: 174 nm

$t_{\text{Co}}=0.6$ nm **10 Co layers**

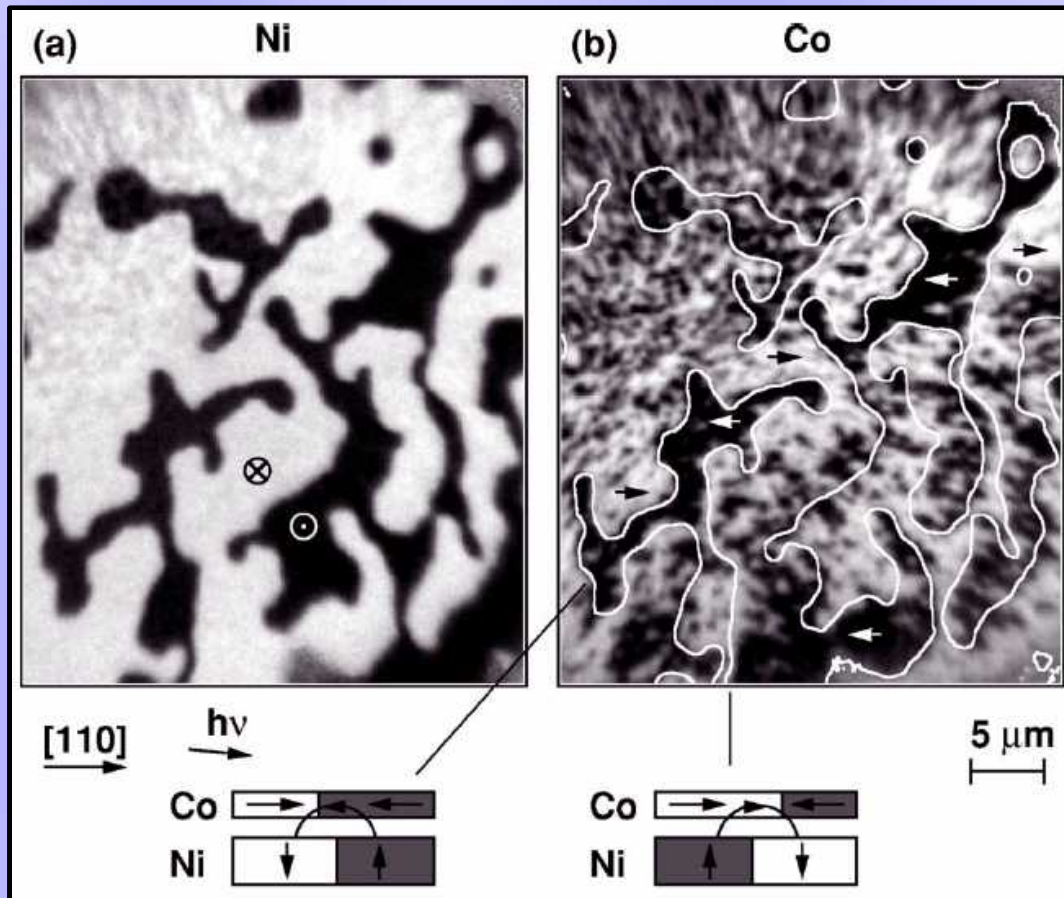
Magnetic fields that originate from the stripe domain structure in $[\text{NiFe}/\text{Au}/\text{Co}/\text{Au}]_N$ multilayers are of the order of 0.1 T.

Magnetostatic coupling



Magnetic field of the stack of the infinite "stripe domains" (from Biot-Savart law). Domain width=100, thickness=1, multilayer period=7.

Magnetostatic coupling

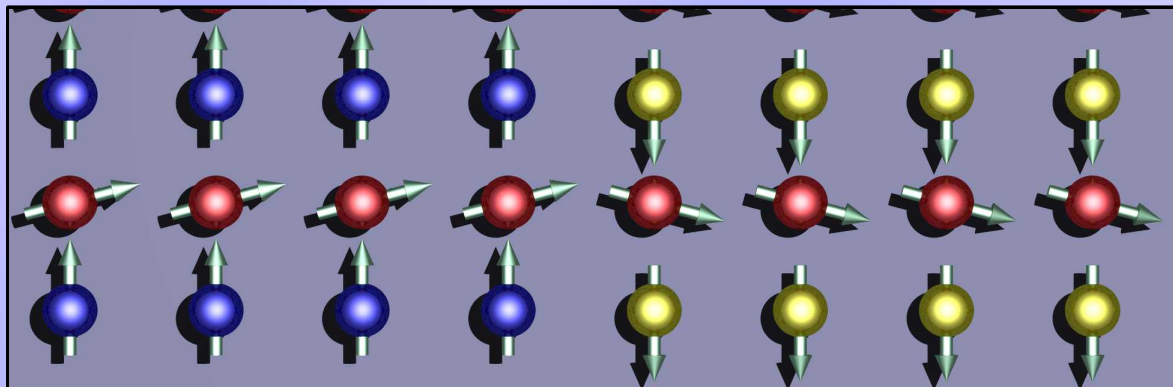


photoemission electron microscopy (PEEM) +
X ray magnetic circular dichroism (XMCD)



- Cu(001)/Ni/Cu/Co
- Cu – wedge (1ML/10m)
- electron beam evaporation
- Ni – **perpendicular anisotropy**
- field of Ni DW in Co: 250Oe

Magnetostatic coupling



Co

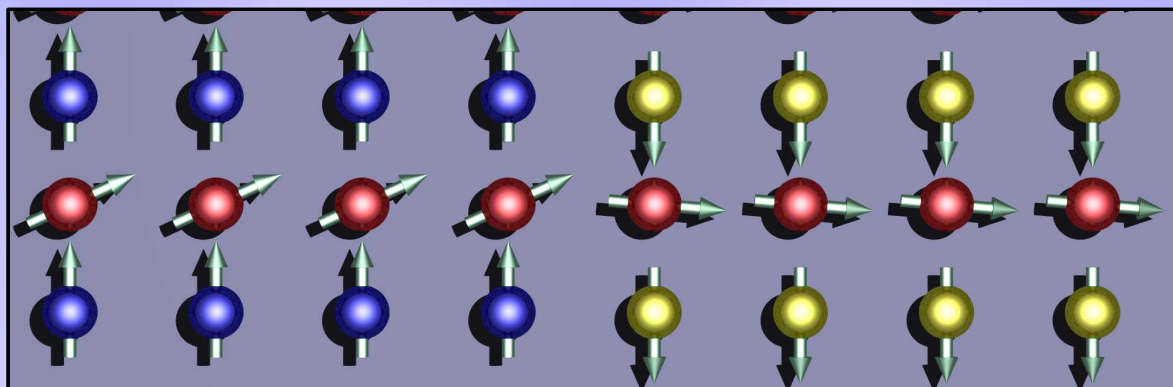
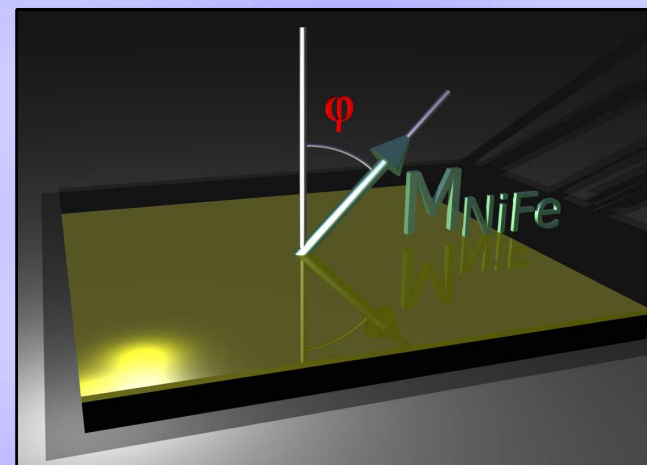
$$H = H_{\text{domains}}$$

NiFe

Co

$$\cos(\varphi_{\uparrow}) = \frac{H_{\text{appl}} + H_d}{M_S^{\text{Co}}}$$

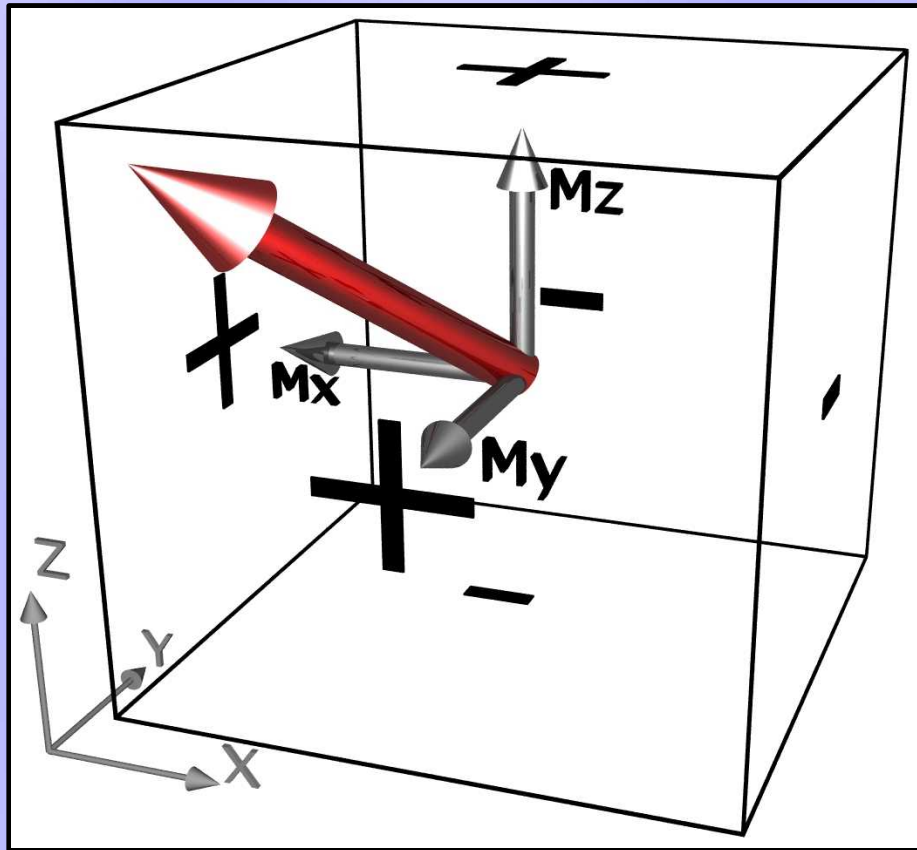
$$\cos(\varphi_{\downarrow}) = \frac{H_{\text{appl}} - H_d}{M_S^{\text{Co}}}$$



$$H = H_{\text{domains}} + H_{\text{appl}}$$

\uparrow H_{appl}

Magnetostatic coupling



С. В. Вонсовский, МАГНЕТИЗМ
«Наука», 1971

“Magnetic charges” present within the volume of the magnetized body
And on its outer boundaries are the sources of magnetic field.

Micromagnetic simulation

$$r = [(x_n - x_q)^2 + (z_n - z_q)^2 + (z_n - z_q)^2]^{1/2}$$

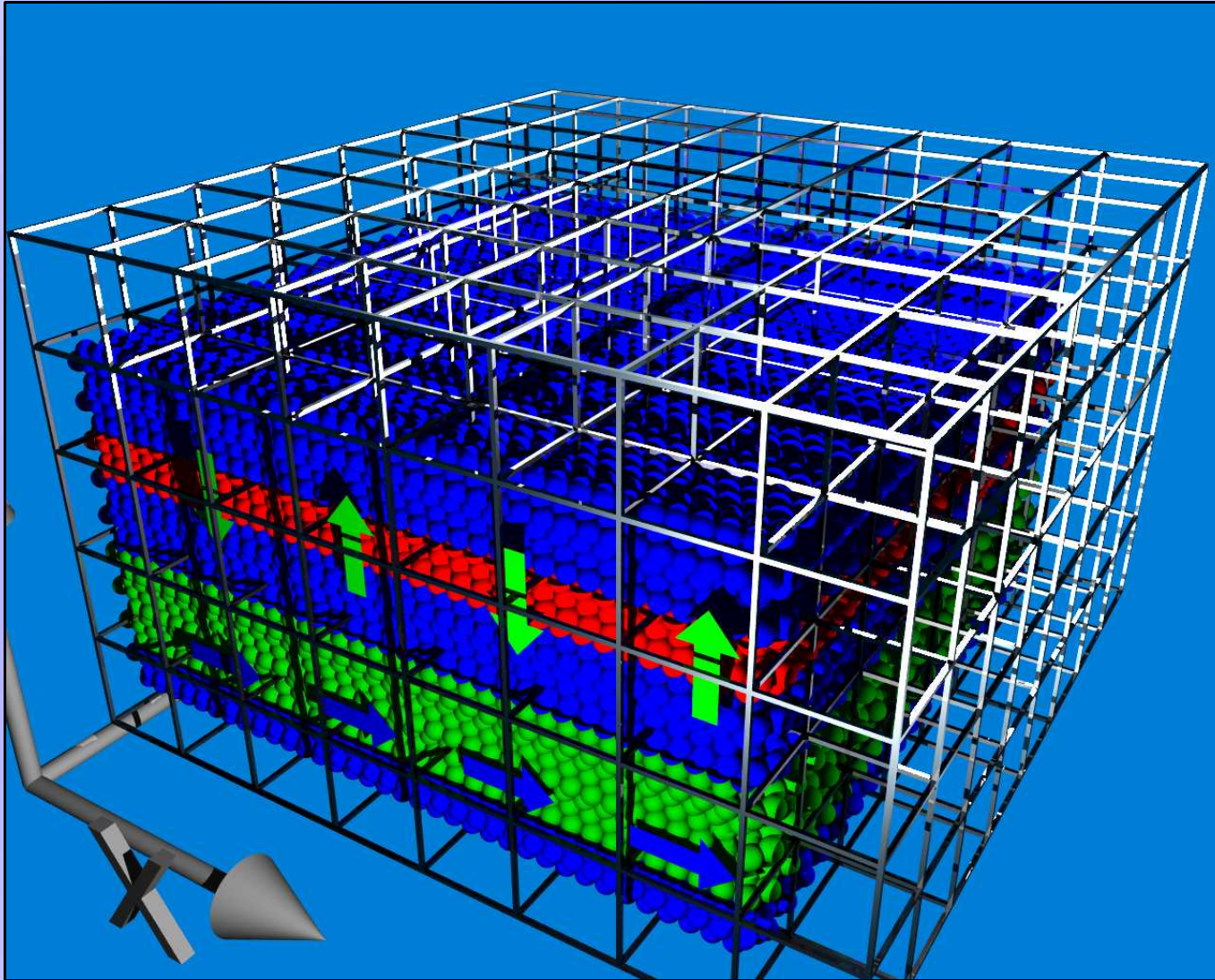
$$\phi_m^{(i)} = \frac{1}{4\pi} \frac{(\vec{\mu} \vec{r})}{r^3}$$

$$\vec{H} = -\vec{\nabla} \phi$$

$$\phi_m = \frac{1}{4\pi} \int d\tau (\vec{M} \nabla_q r^{-1})$$

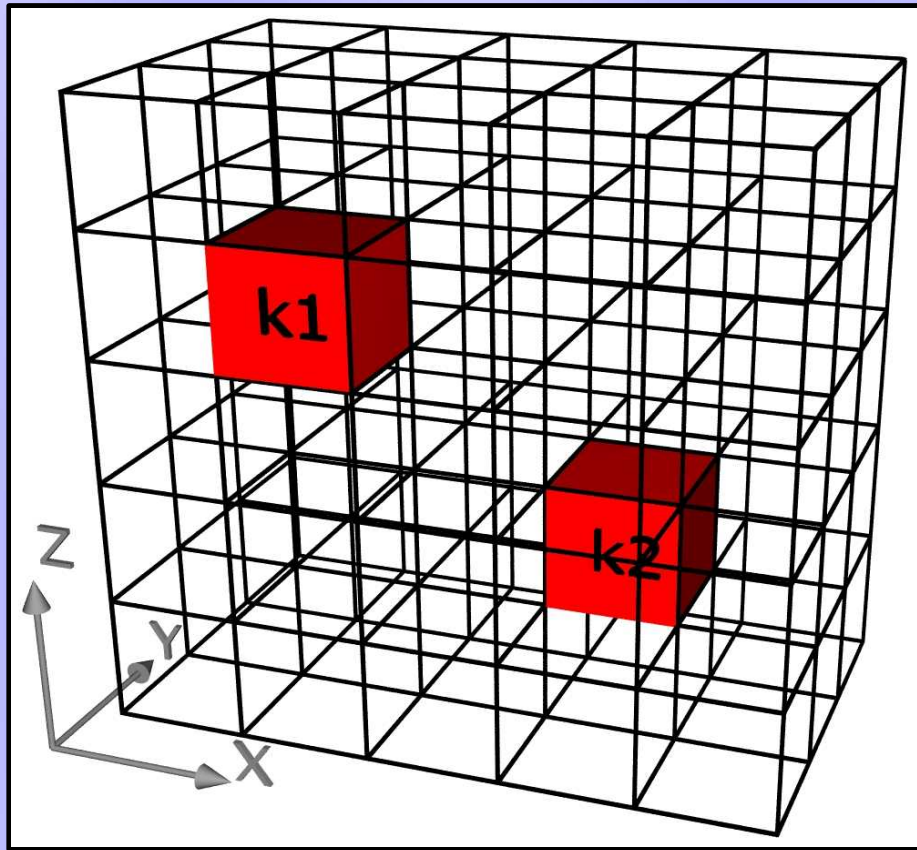
$$(\vec{M} \nabla_q r^{-1}) = \nabla_q (r^{-1} \vec{M}) - \frac{1}{r} \nabla_q \vec{M}$$

$$\phi_m = \frac{1}{4\pi} \left(-\int d\tau \frac{\nabla_q \vec{M}}{r} + \oint dS \frac{\vec{n} \vec{M}}{r} \right)$$



A single cell contains MANY atoms.

Continuous approximation - the magnetization is a continuous function of the position.



LLG equation: $\vec{m} = \frac{\vec{M}}{M_S}$

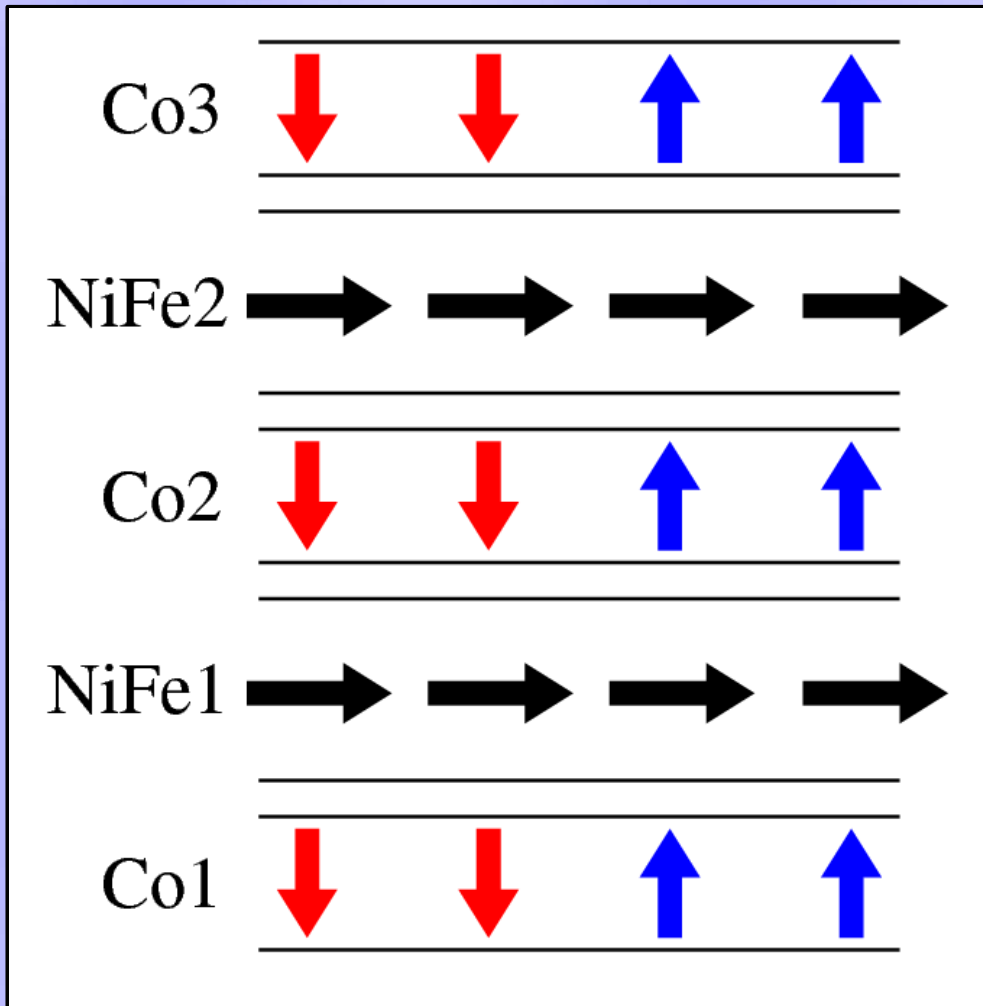
$$\frac{d\vec{m}}{dt} = \gamma_0 (\vec{m} \times \vec{H}_{eff}) + \alpha \left[\vec{m} \times \frac{d\vec{m}}{dt} \right]$$

$$\vec{H}_{eff} = \frac{-1}{\mu_0 \vec{M}_S} \frac{\delta \epsilon}{\delta \vec{m}}$$

$$\vec{H}_{eff} = \frac{A}{\mu_0 \vec{M}_S} \frac{\delta (\nabla \vec{m})^2}{\delta \vec{m}} - \frac{1}{\mu_0 \vec{M}_S} \frac{\delta \epsilon_K}{\delta \vec{m}} + \vec{H}_{appl} + \vec{H}_d$$

\vec{H}_{eff} = “exchange energy” + “anisotropy energy” + “external field” + “own field”

Magnetostatic interactions between cells have of **a long-range character**.

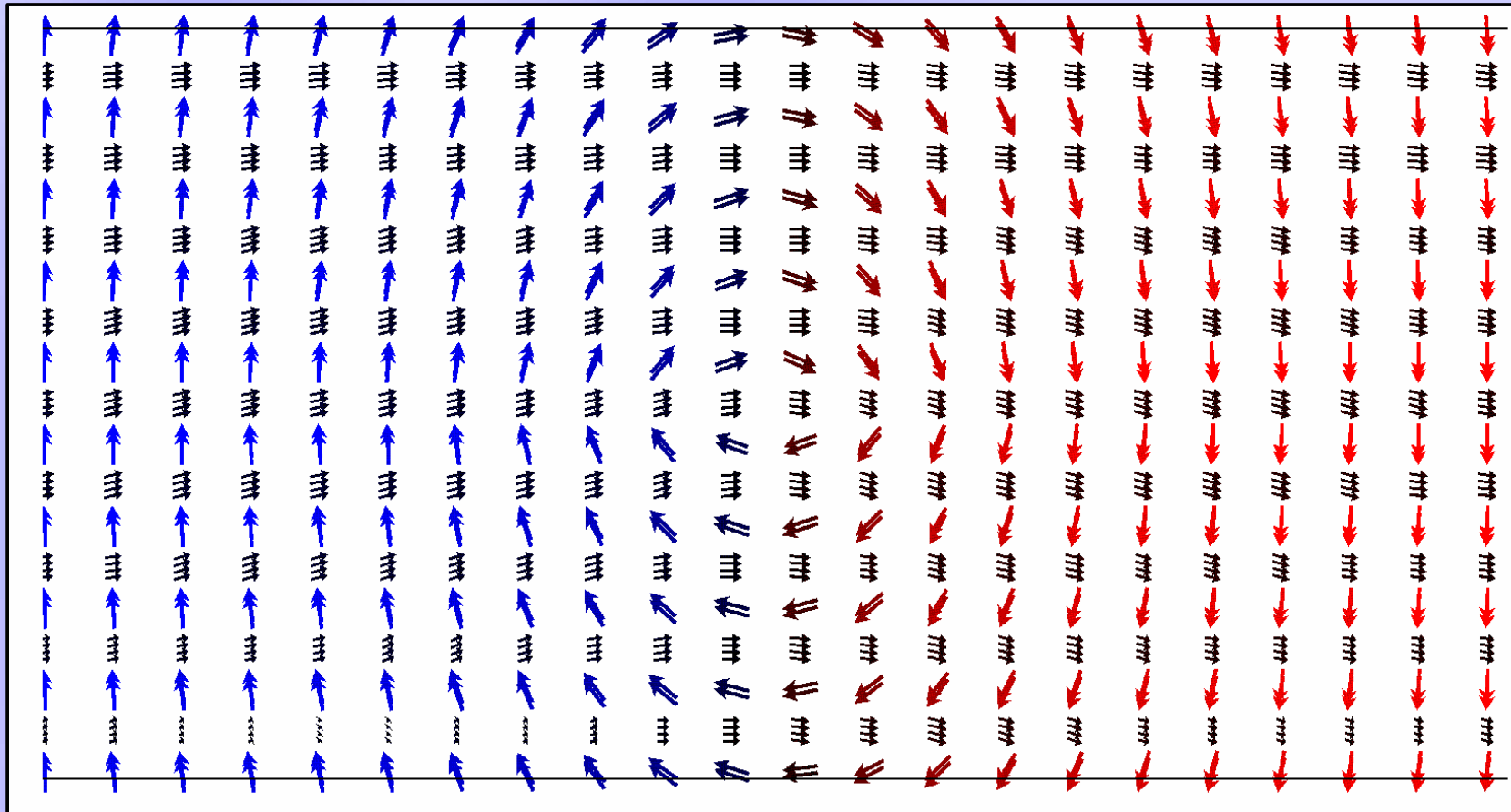


- simulation of remanent state
- starting configuration-stripe domains in Co sublayers
- starting configuration-monodomain state in NiFe sublayers

*Simulation with free oommf package from NIST; $(1 \times 1 \mu\text{m}^2) \times 55\text{nm}$; Co domains 200 nm wide; $\alpha=0.5$; regular mesh with cell size of $(5 \times 0.5 \times 50\text{nm}^3)$; stiffness: Co: $30\text{e-}12 \text{ J/m}$, NiFe: $13\text{e-}12 \text{ J/m}$

Magnetostatic coupling

Micromagnetic simulation*



↑
thickness
↓

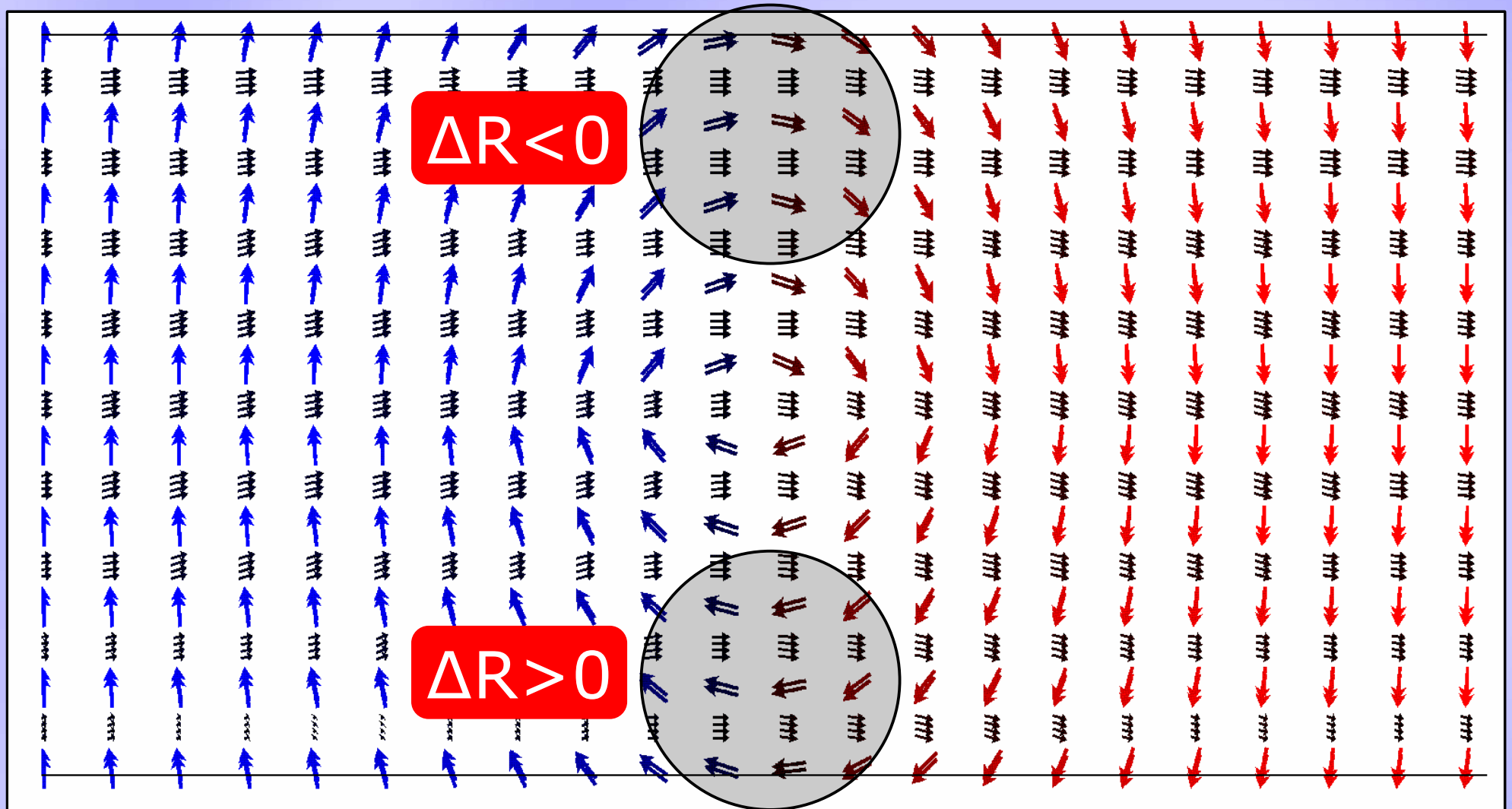
[Co(1nm)/Au(1.5nm)/Ni₈₀Fe₂₀(2 nm)/Au(1.5 nm)]₉/Co(1nm)

$H=0$

*Simulation with free oommf package from NIST; $(1 \times 1 \mu\text{m}^2) \times 55\text{nm}$;
Co domains 200 nm wide; $\alpha=0.5$; regular mesh with cell size of
 $(5 \times 0.5 \times 50\text{nm}^3)$; stiffness: Co: $30\text{e-}12 \text{ J/m}$, NiFe: $13\text{e-}12 \text{ J/m}$

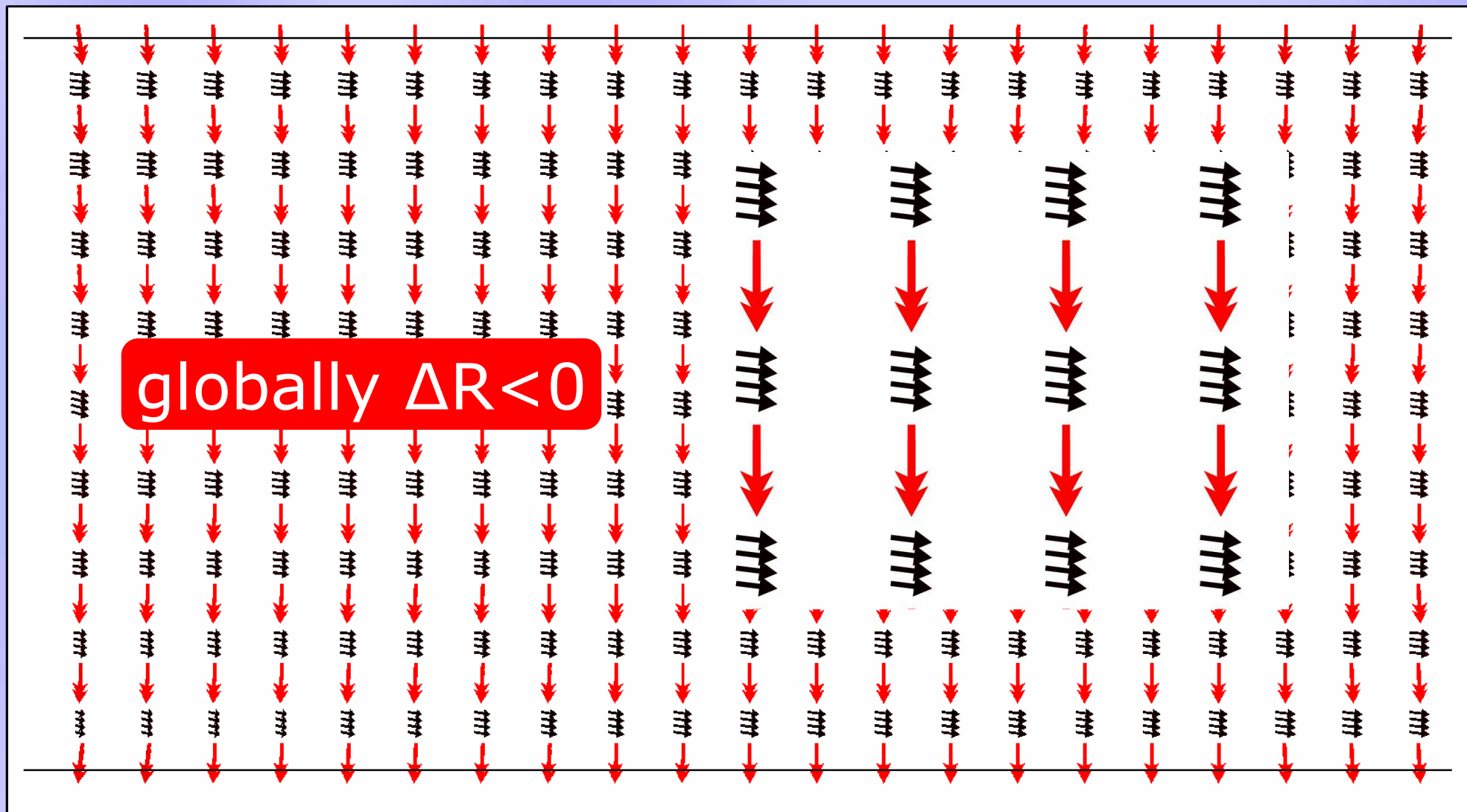
Magnetostatic coupling

Micromagnetic simulation



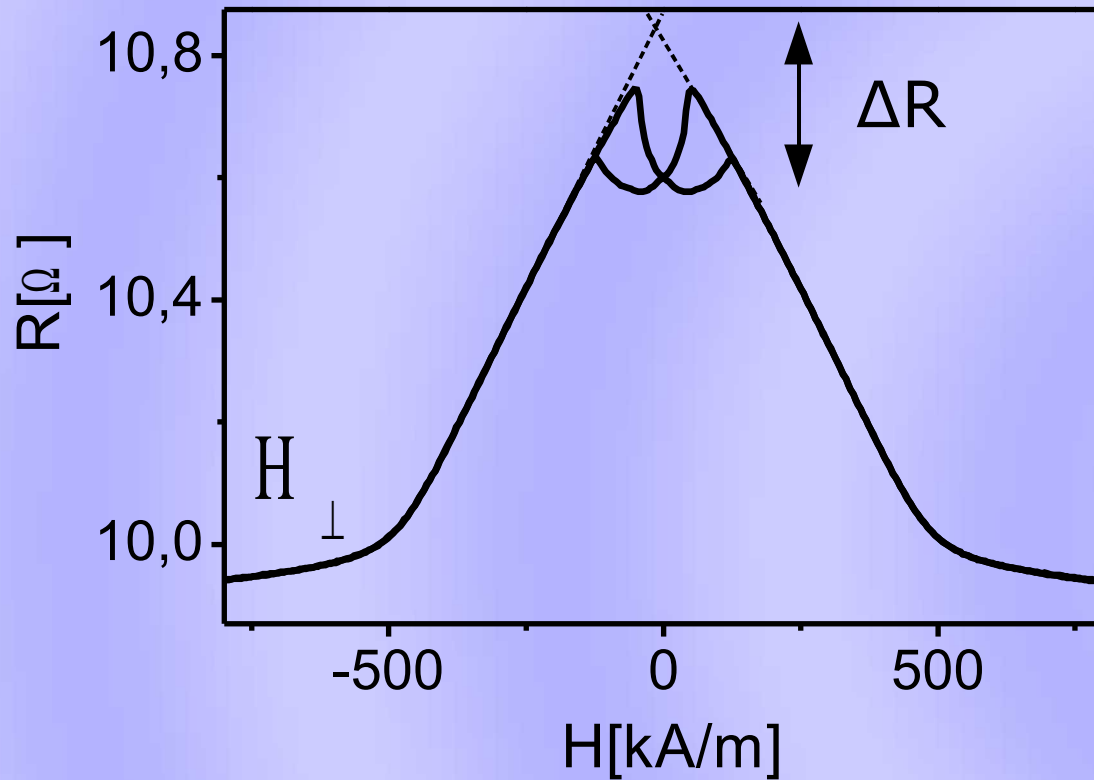
$$\Delta R \propto \cos(\varphi)$$

$$H=0$$



$$\Delta R \propto \cos(\varphi)$$

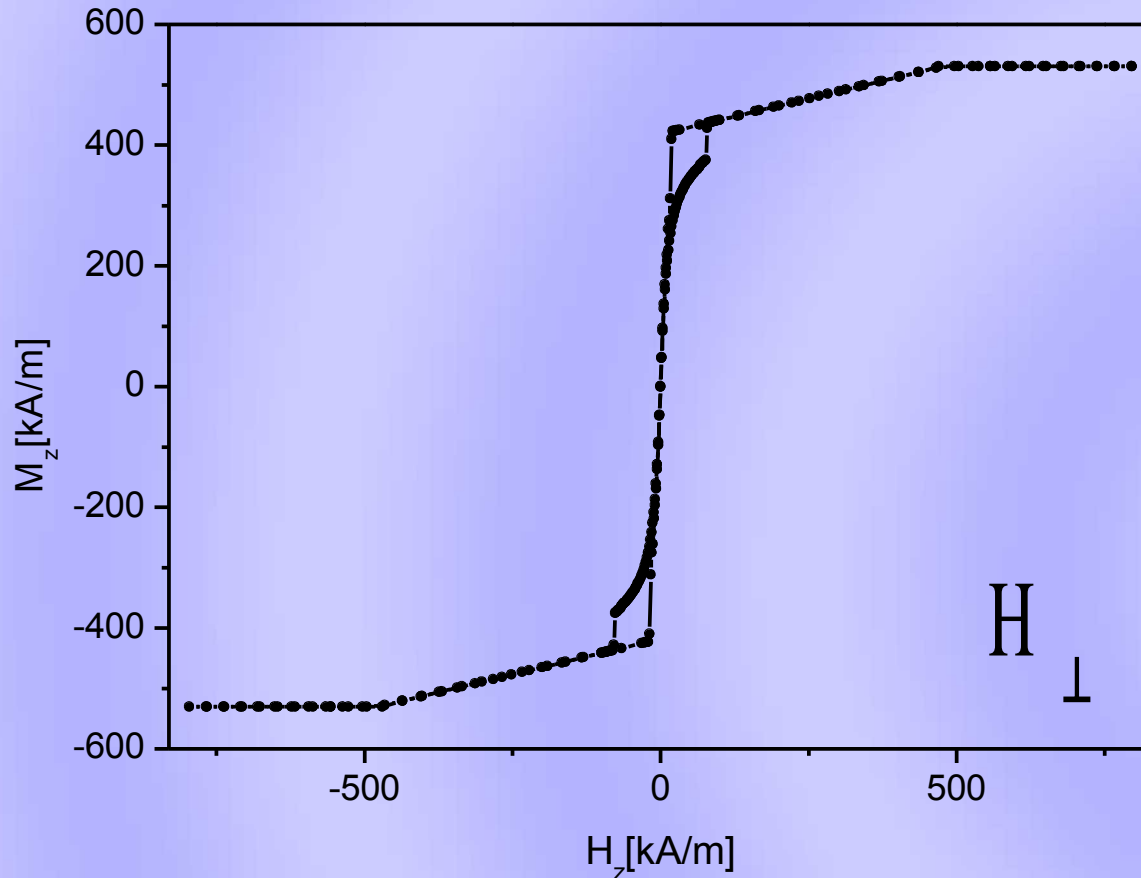
$$H=0$$



$$\Delta \cos(\varphi) = 0.265$$

Magnetostatic interactions between Co and NiFe layers lead to the decrease of the average cosine of the angle between magnetic moments of neighboring sublayers \Rightarrow **resistance decrease**.

Magnetostatic coupling

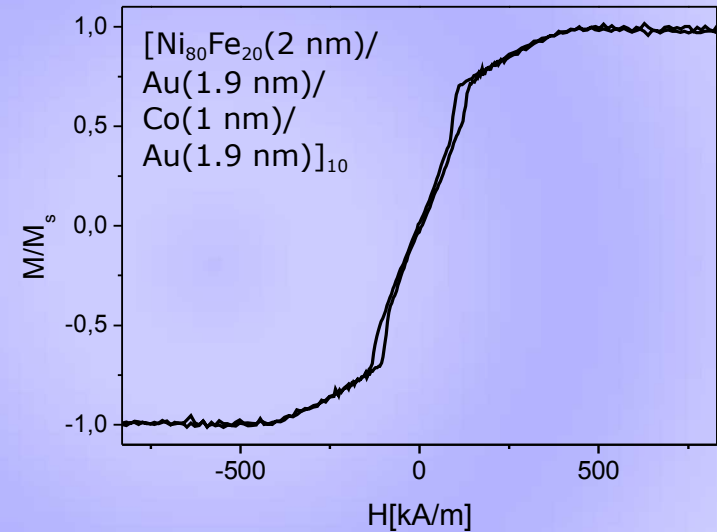


Micromagnetic simulation*

Co-perpendicular anisotropy

NiFe-shape anisotropy

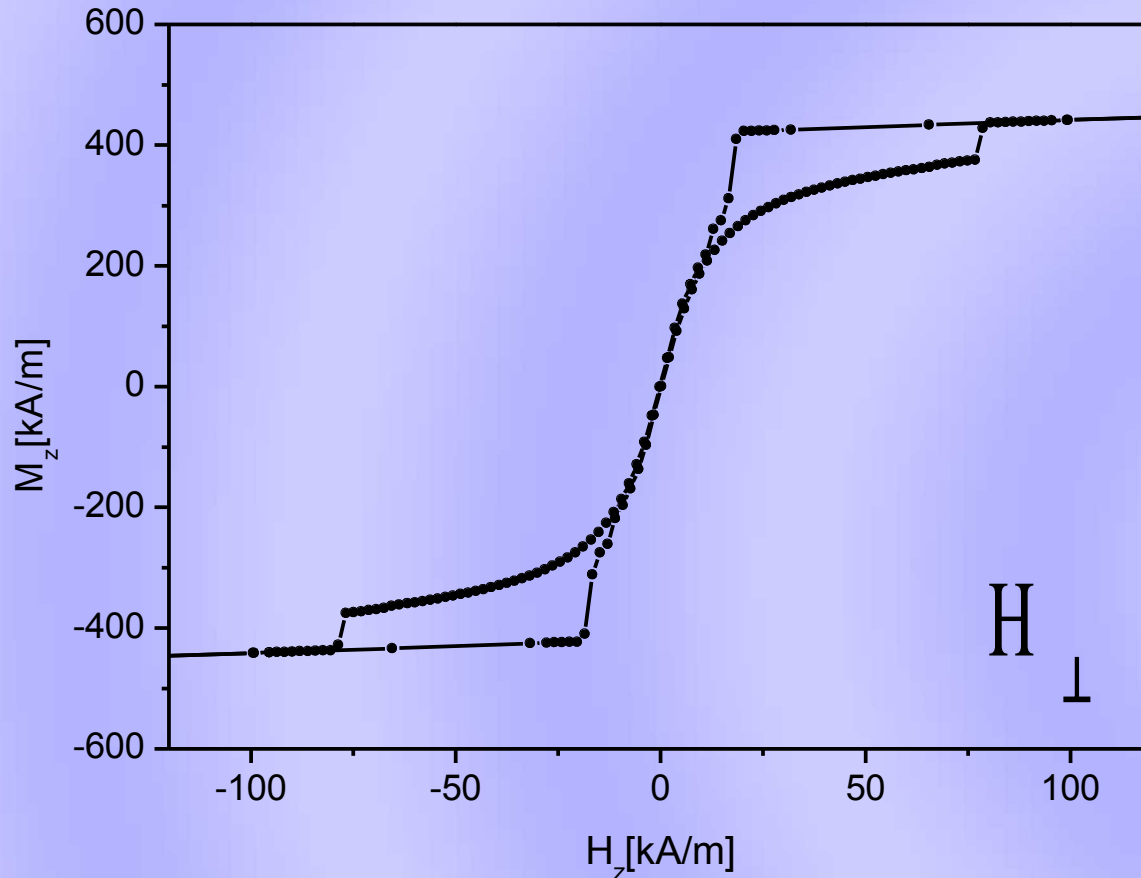
measurement:



$[\text{Co}(1\text{nm})/\text{spacer}(1\text{nm})/\text{NiFe}(1\text{nm})/\text{spacer}(1\text{nm})]_4/\text{Co}(1\text{nm})$

*Simulation with free oommf package from NIST (M.J. Donahue and D.G. Porter);
 $\alpha=0.5$; regular mesh with cell size of $(5 \times 20000 \times 1 \text{ nm}^3)$;
stiffness: Co: $30 \text{ e-}12 \text{ J/m}$, NiFe: $13 \text{ e-}12 \text{ J/m}$

Magnetostatic coupling

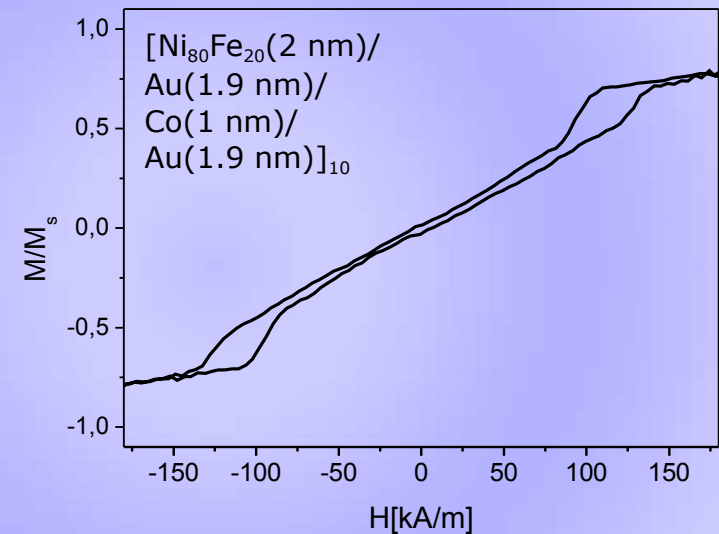


Micromagnetic simulation

Co-perpendicular anisotropy

NiFe-shape anisotropy

measurement:

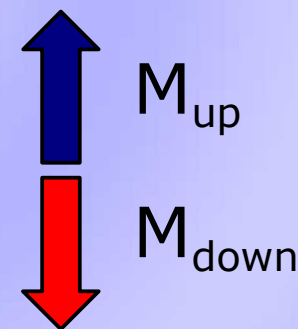
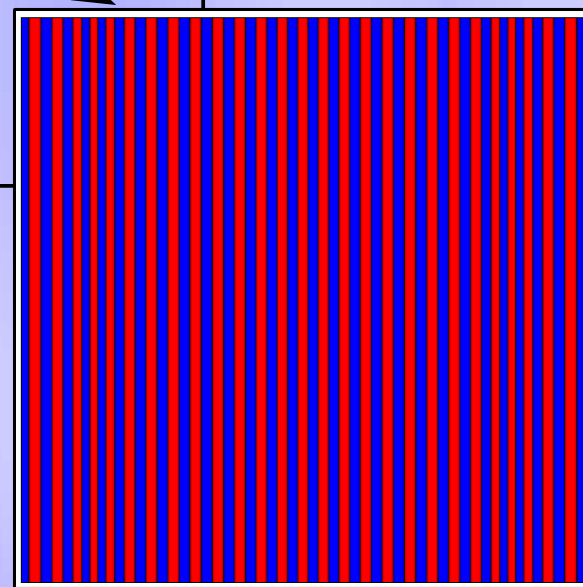
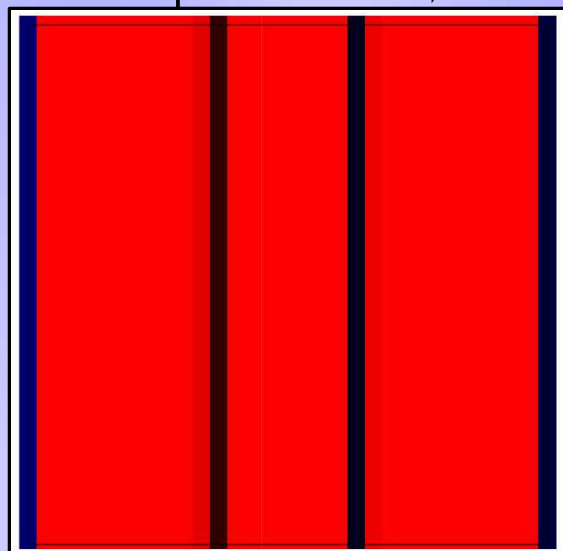
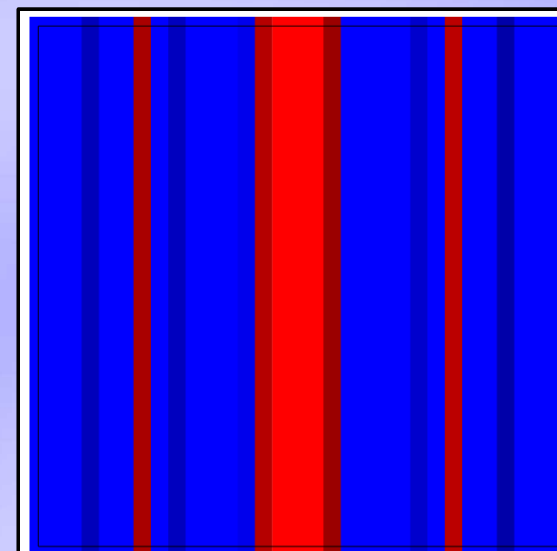
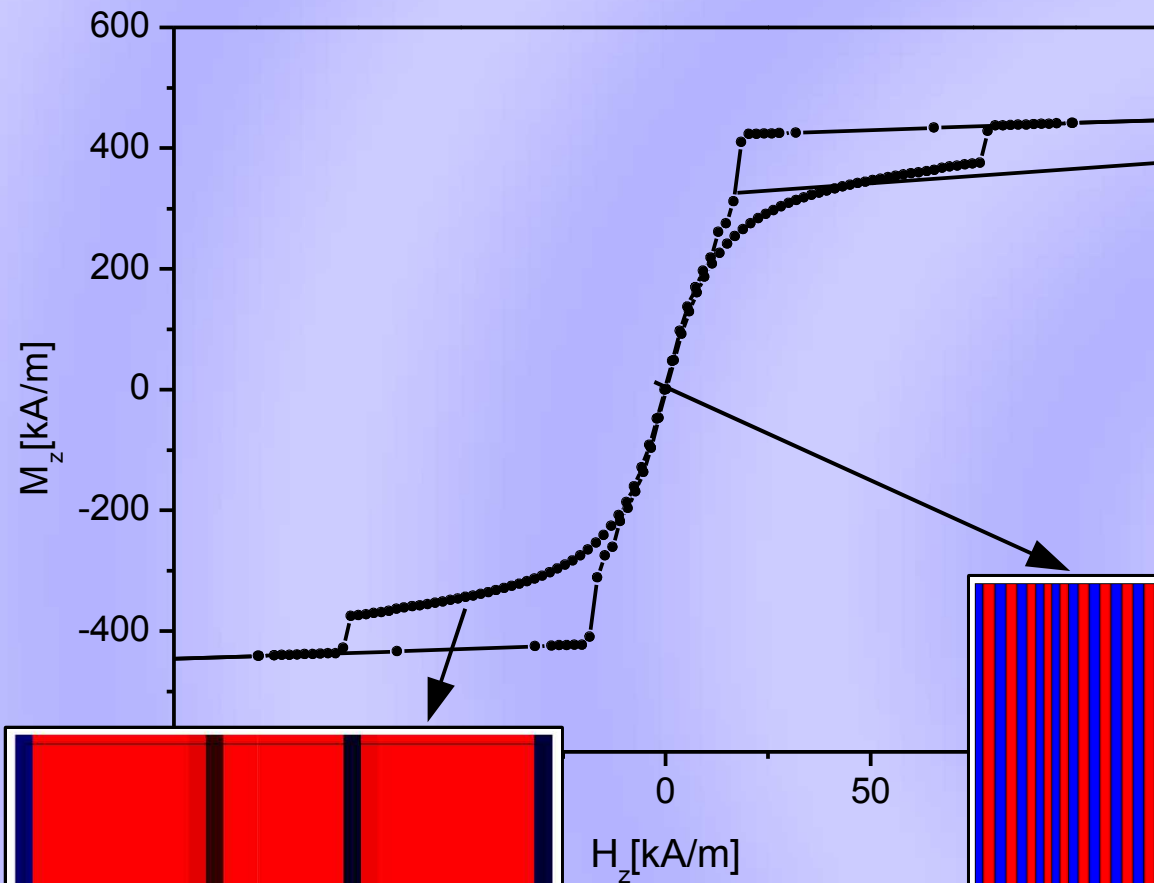


$[\text{Co}(1\text{nm})/\text{spacer}(1\text{nm})/\text{NiFe}(1\text{nm})/\text{spacer}(1\text{nm})]_4/\text{Co}(1\text{nm})$

No attempts were made to exactly mirror the $M(H)$ dependence ,i.e. , nucleation and annihilation fields.

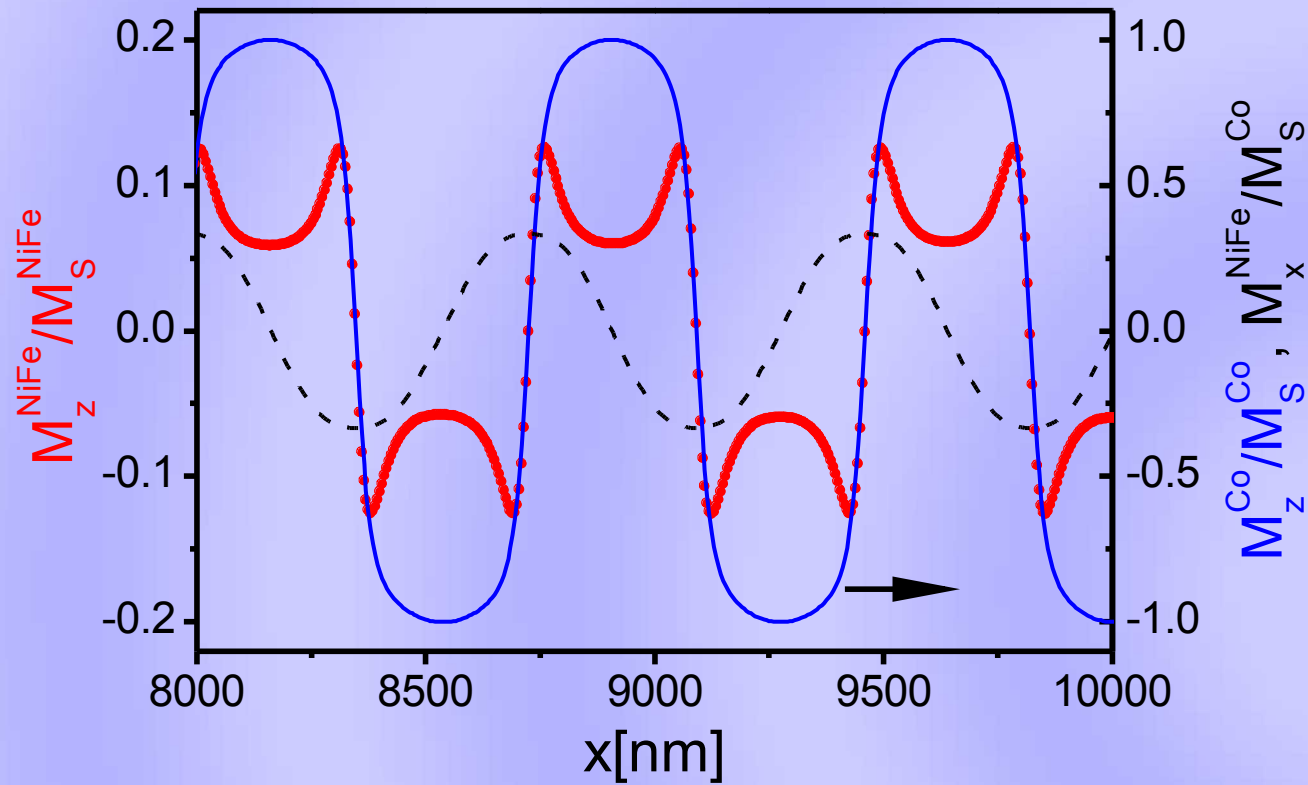
Magnetostatic coupling

Micromagnetic simulation

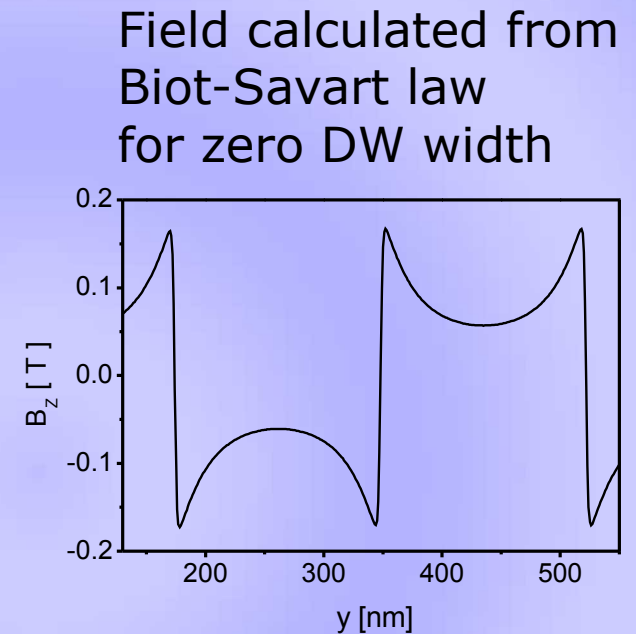


Stripe domains

Magnetostatic coupling

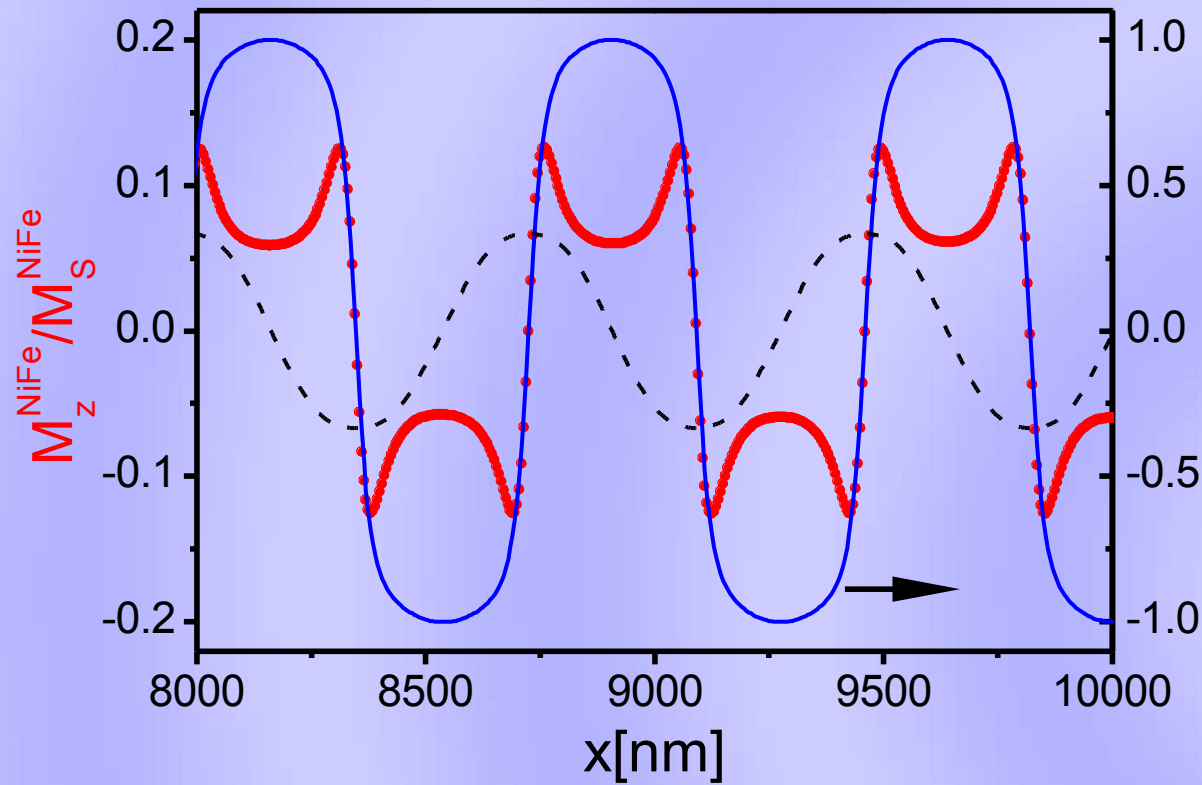


Micromagnetic simulation

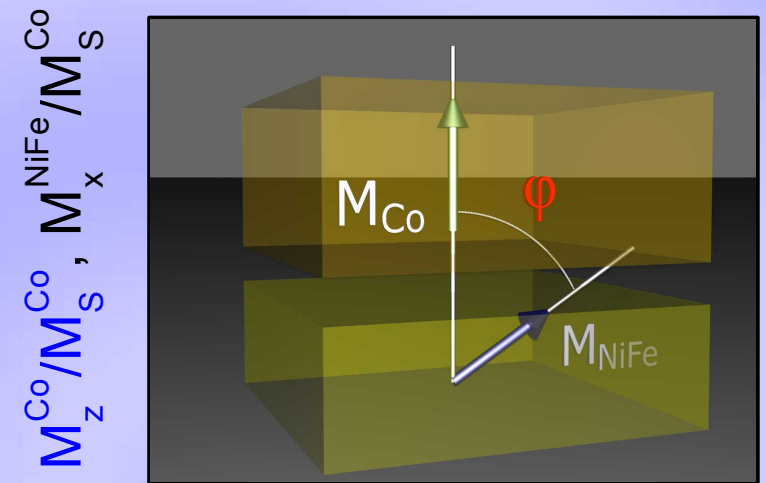


Magnetostatic interactions between Co and NiFe layers lead to the spatial replication of the z-component of magnetic moment of Co sublayers in NiFe sublayers \Rightarrow resistance decrease.

Magnetostatic coupling



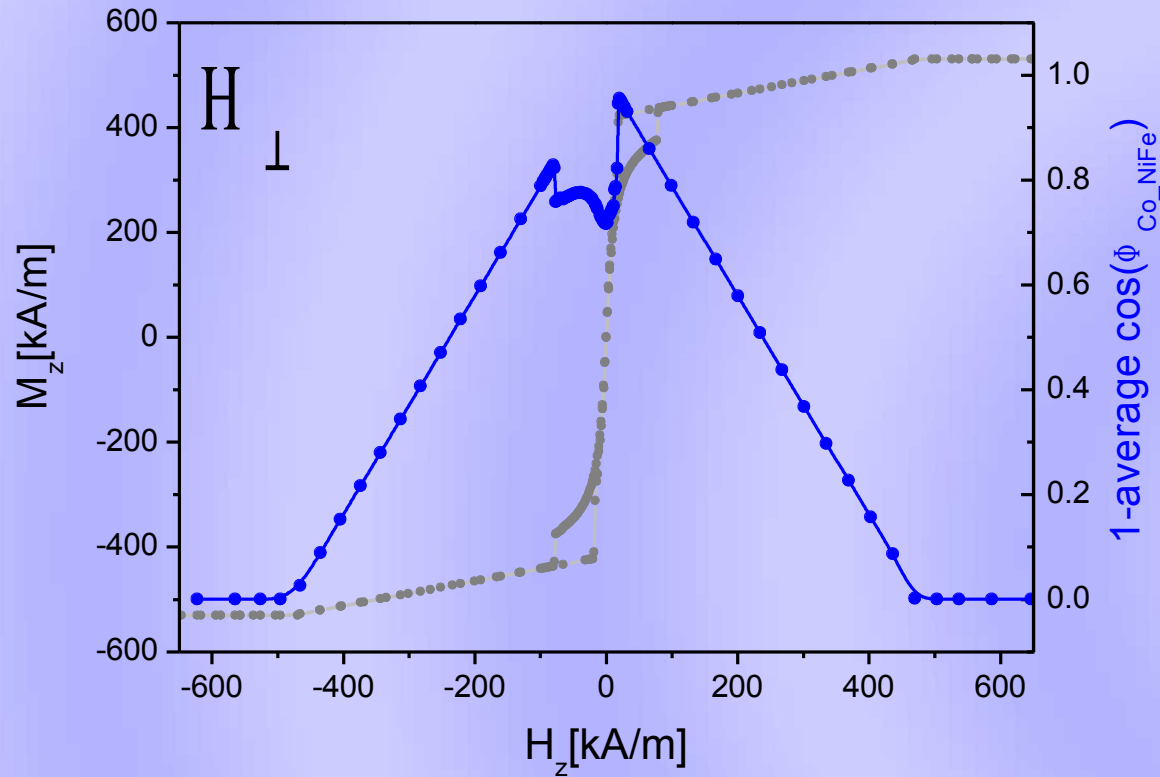
Micromagnetic simulation



$$\Delta R_{GMR} = \frac{1}{n} \sum_n \cos(\phi_{NiFe-Co})$$

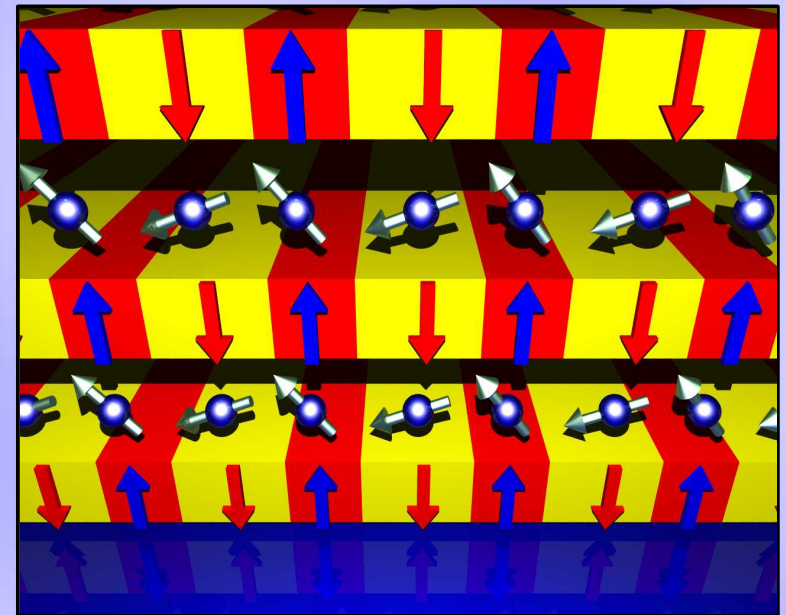


Magnetostatic coupling



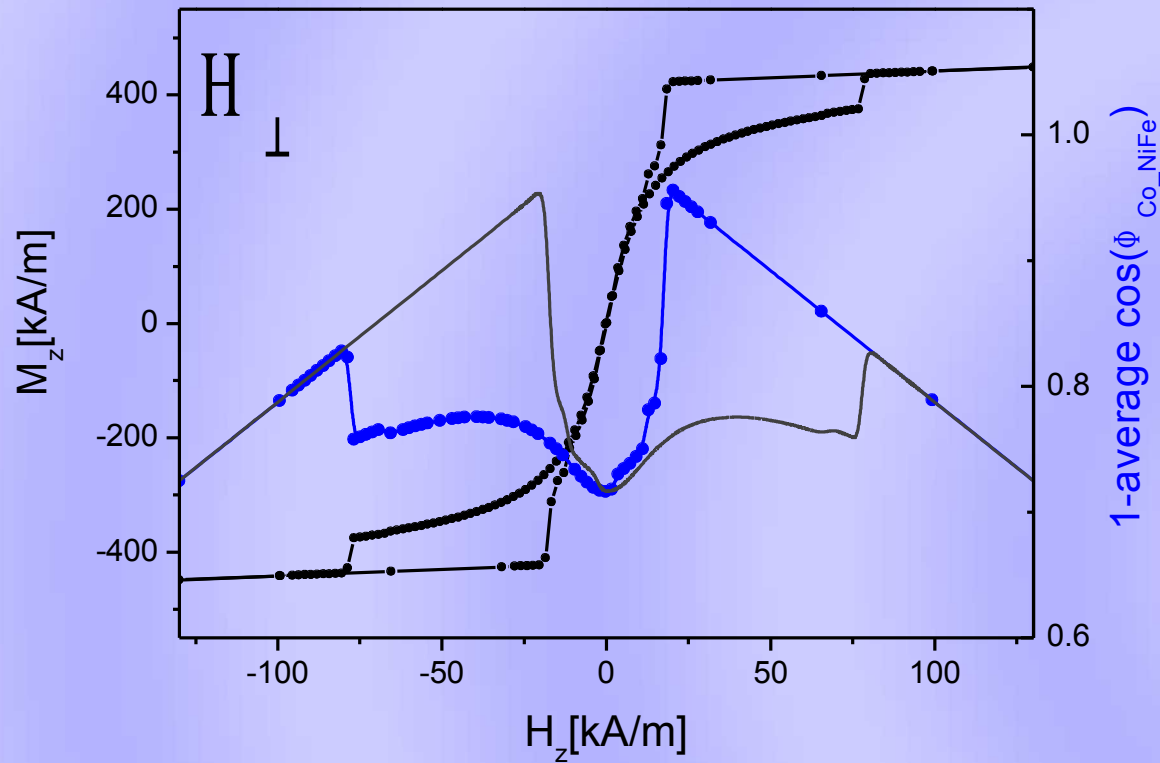
Micromagnetic simulation

The GMR(H) dependence was calculated as proportional to an average cosine of the angle between magnetic moment direction of juxtaposed NiFe and Co cells.



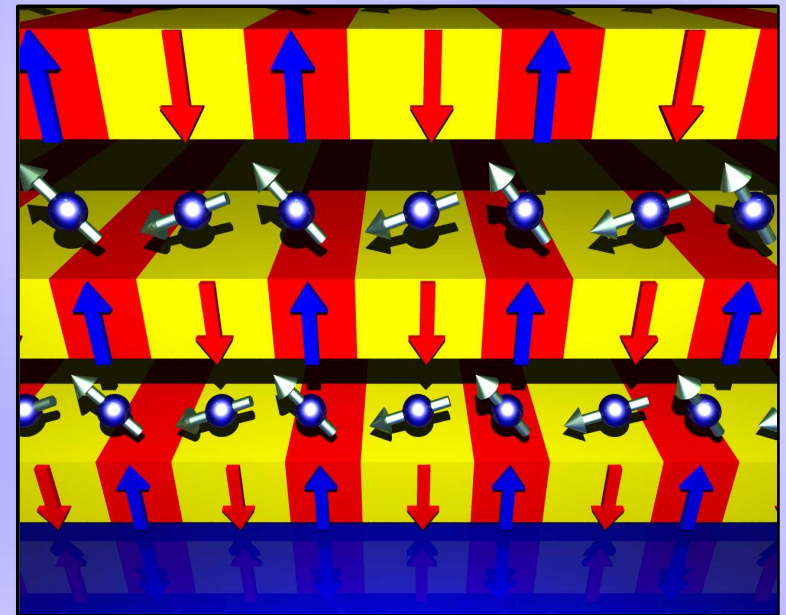
Giant magnetoresistance dependencies of NiFe/Au/Co/Au multilayers can be approximated from micromagnetic simulations.

Magnetostatic coupling



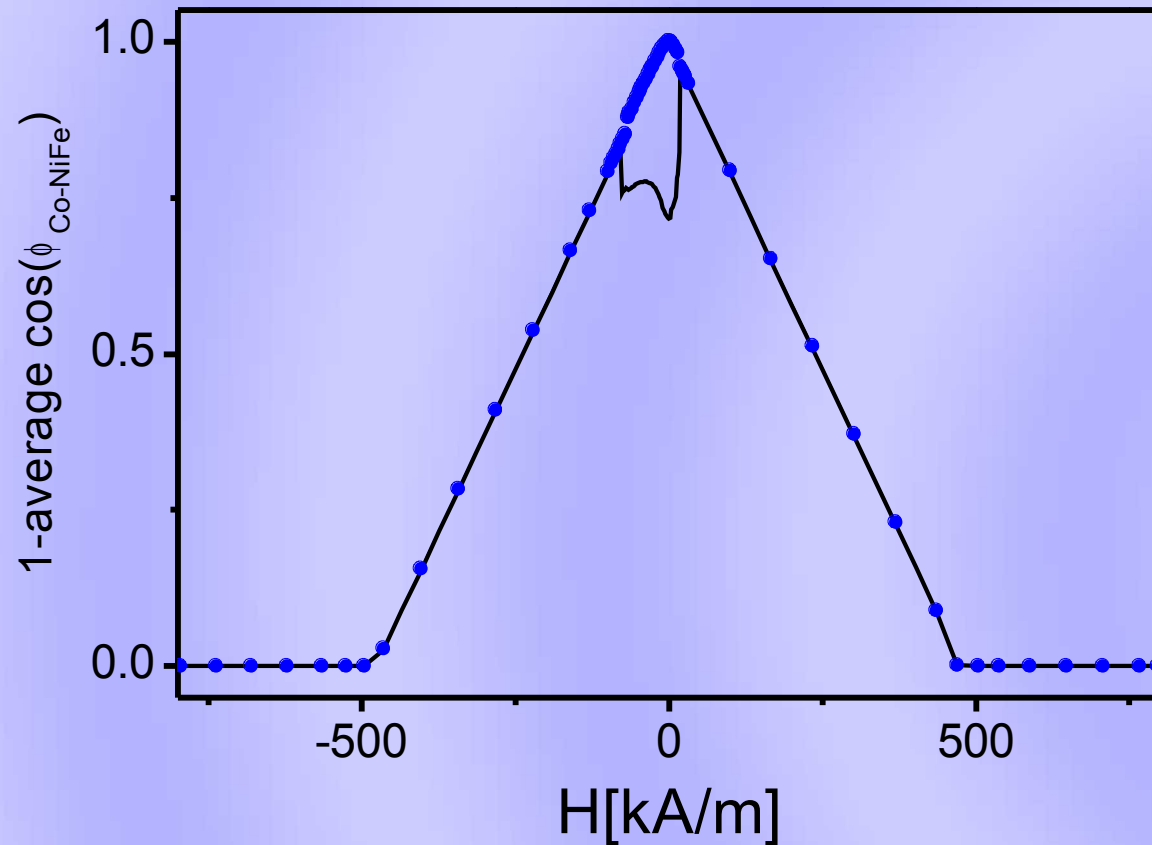
Micromagnetic simulation

The GMR(H) dependence was calculated as proportional to an average cosine of the angle between magnetic moment direction of juxtaposed NiFe and Co cells.



Giant magnetoresistance dependencies of NiFe/Au/Co/Au multilayers can be approximated from micromagnetic simulations.

Magnetostatic coupling



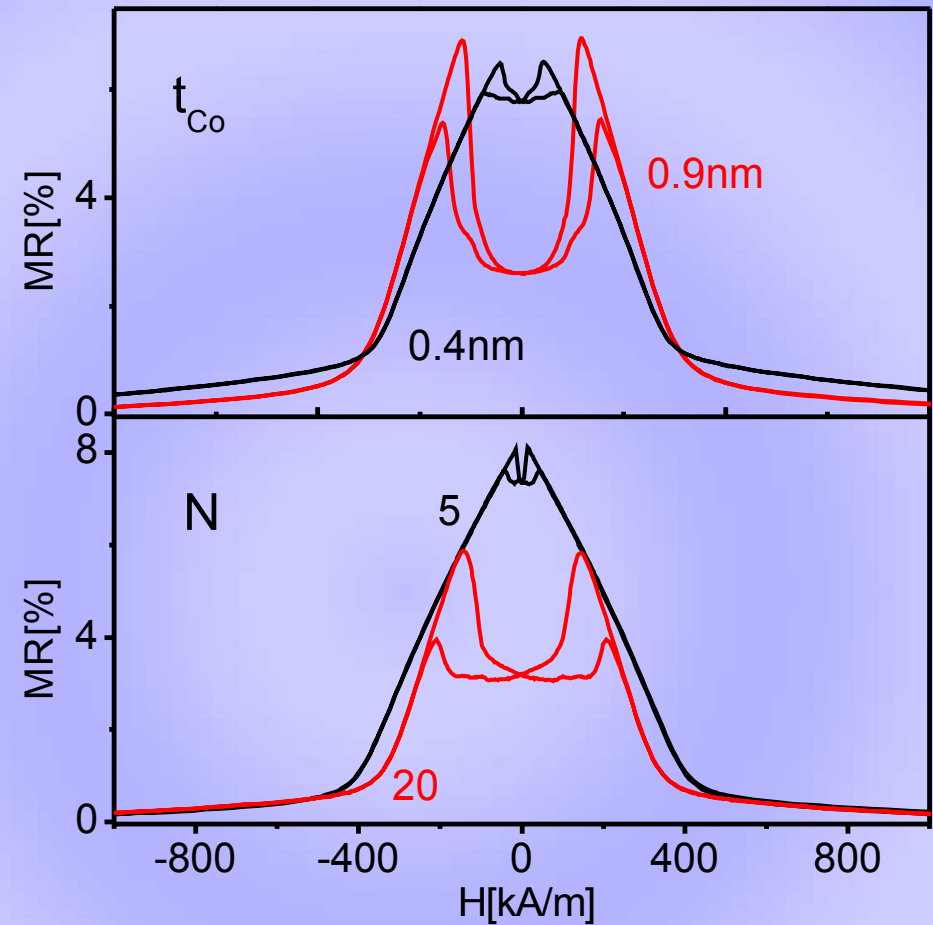
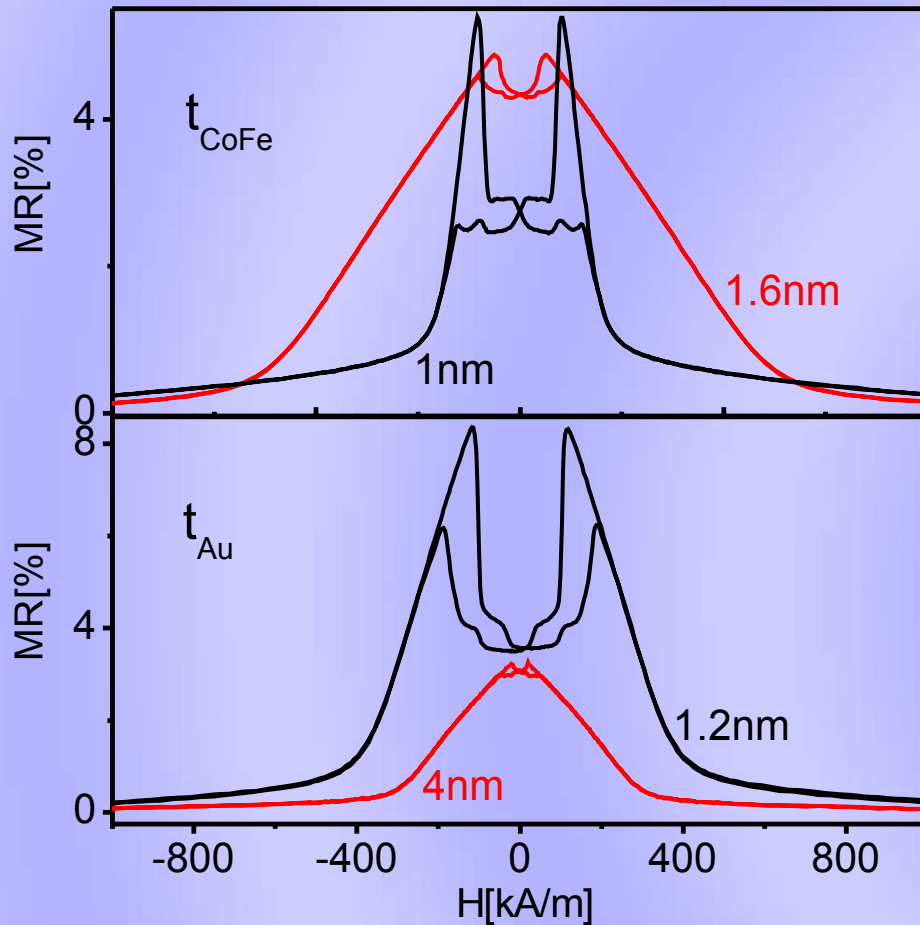
Micromagnetic simulation

The GMR(H) dependence was calculated as proportional to an average cosine of the angle between magnetic moment direction of juxtaposed NiFe and Co cells.

••••• – no coupling

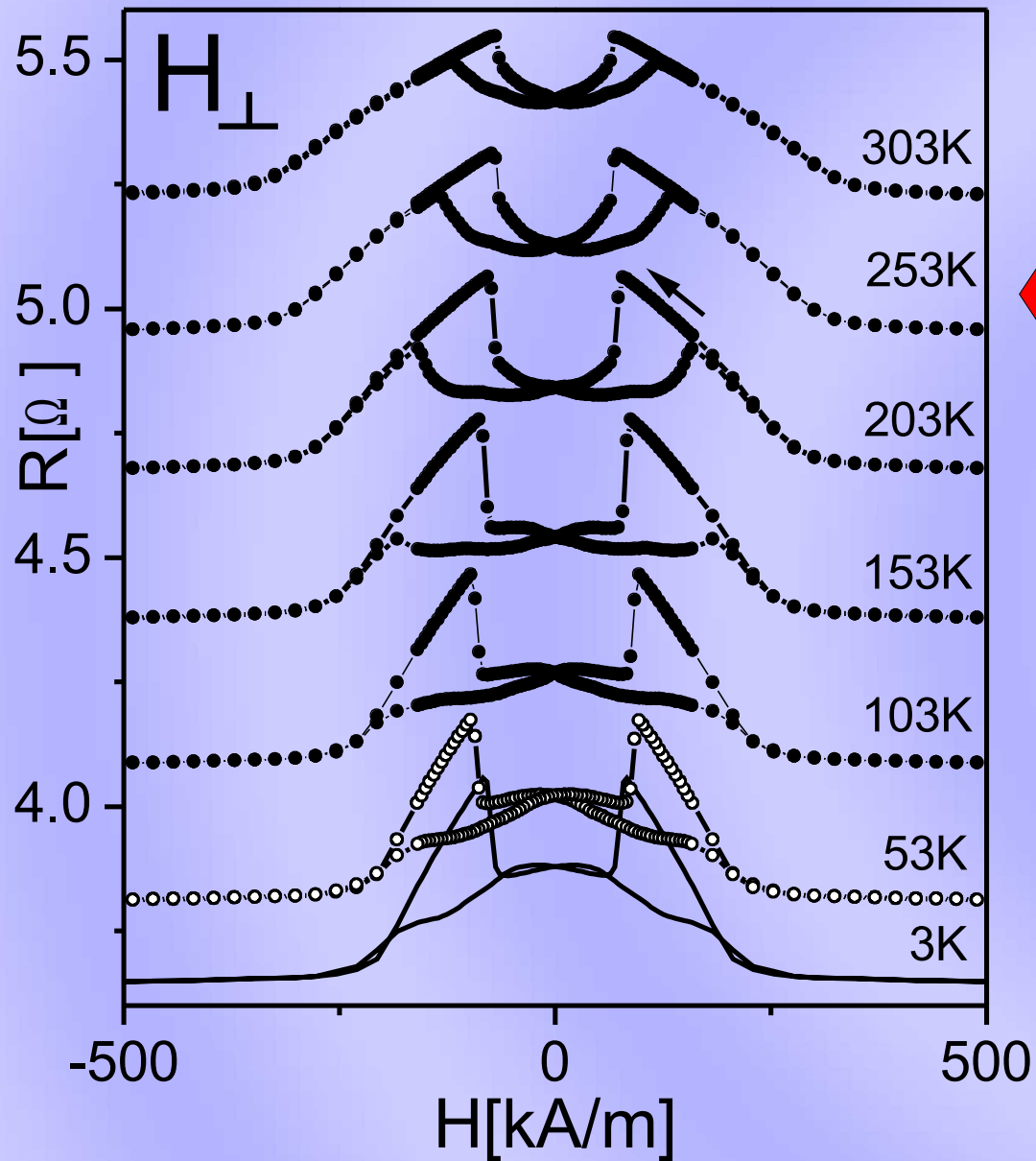
Without magnetostatic coupling between Co and NiFe layers there are no local minima of resistance.

Magnetostatic coupling

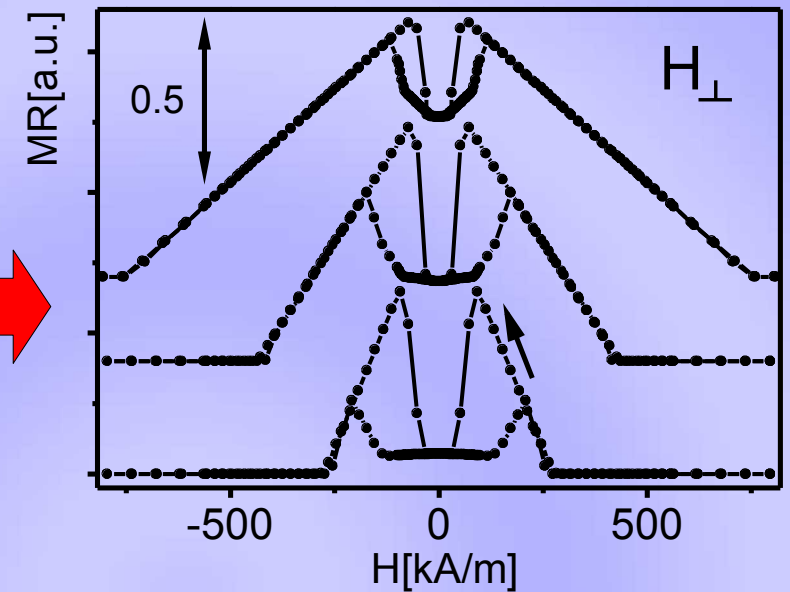


The strength of coupling between sublayers with perpendicular and in-plane anisotropy (depth of resistance minimum) depends on thicknesses of all types of sublayers and on the number of repetitions.

Magnetostatic coupling



Micromagnetic simulation



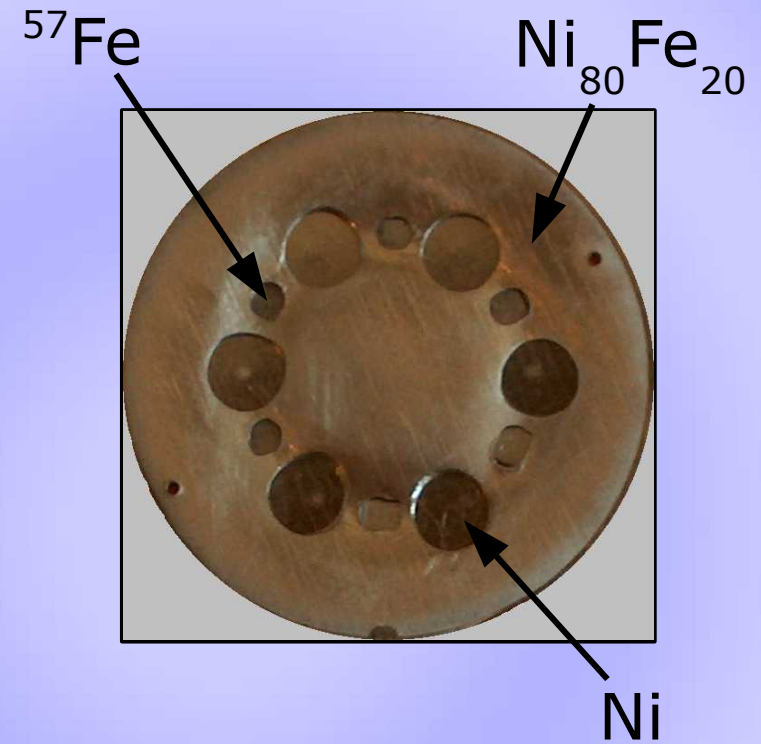
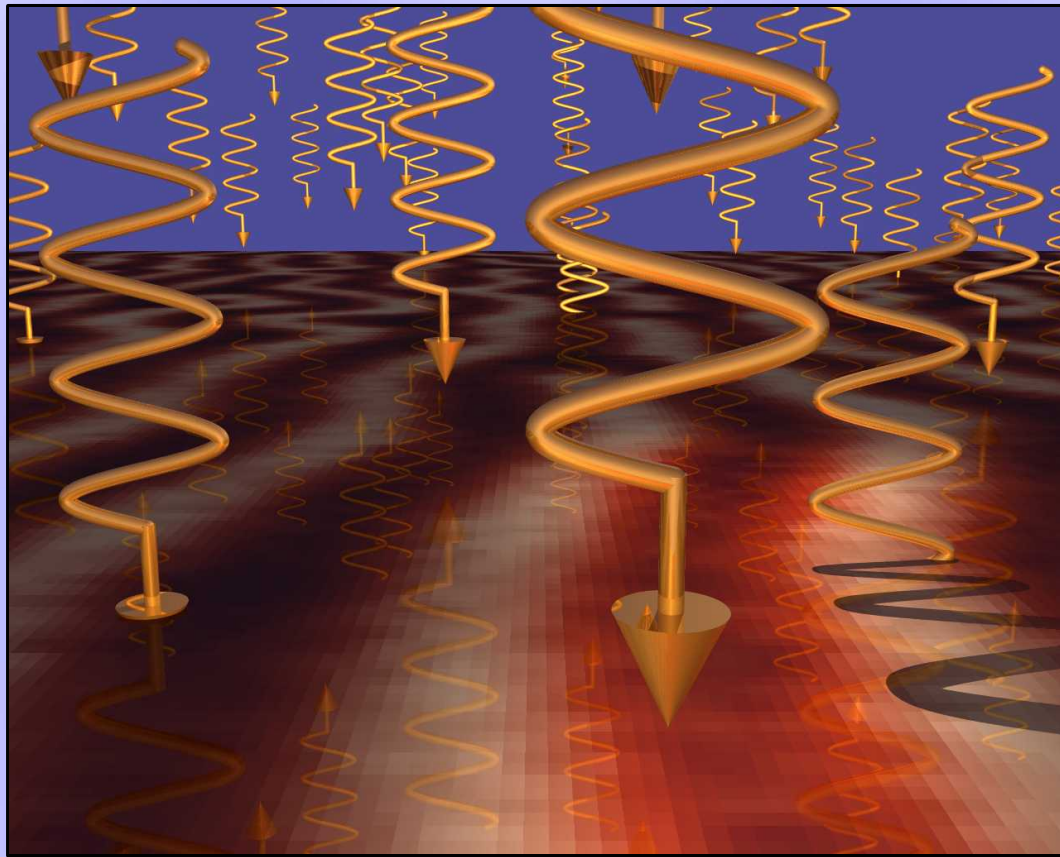
OOMMF simulation with changing perpendicular anisotropy of CoFe sublayers.

from top to bottom:
 $K_U = 0.7 \times 10^6 \text{ J/m}^3$, $K_U = 1 \times 10^6 \text{ J/m}^3$
and $K_U = 1.15 \times 10^6 \text{ J/m}^3$.

$[\text{Co}_{83}\text{Fe}_{17}(1.2 \text{ nm})/\text{Au}(2.2 \text{ nm})/\text{Co}(0.8 \text{ nm})/\text{Au}(2.2 \text{ nm})]_{10}$

Element selective measurements

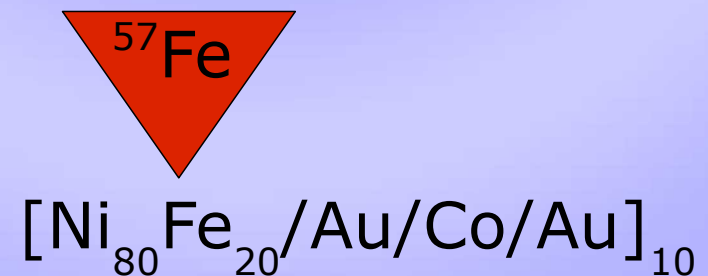
Mössbauer spectroscopy

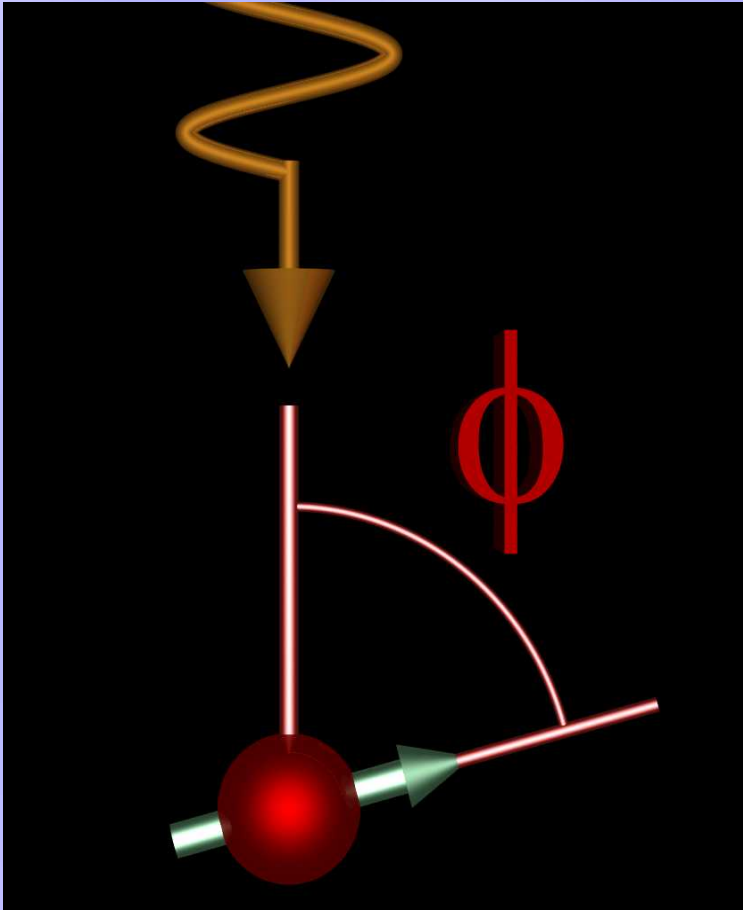


Conversion electron Mössbauer spectroscopy (CEMS)

^{57}Co source

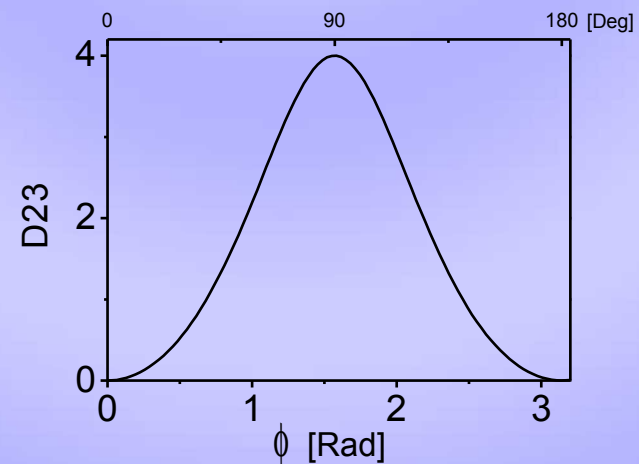
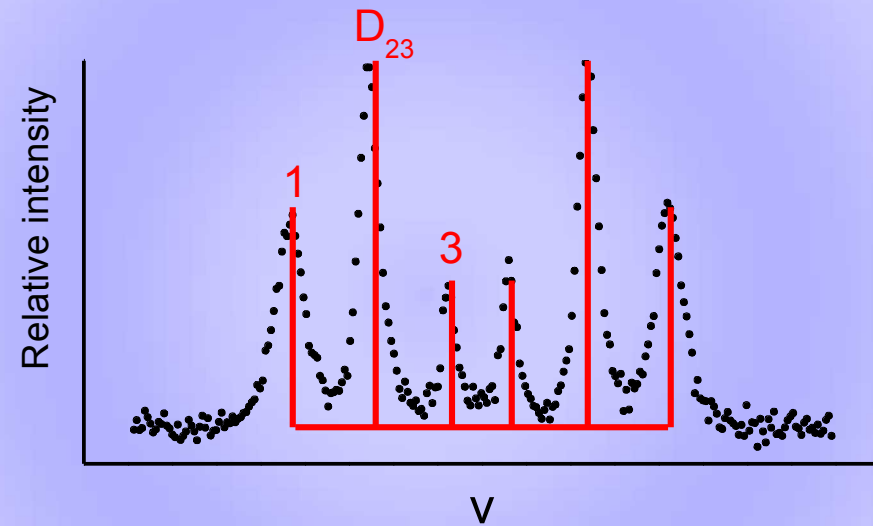
^{57}Fe 95.3 at. %

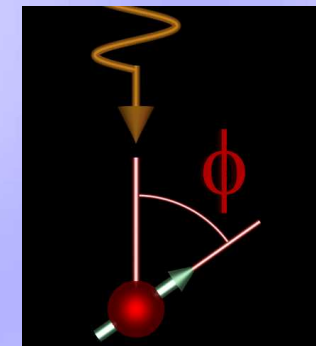
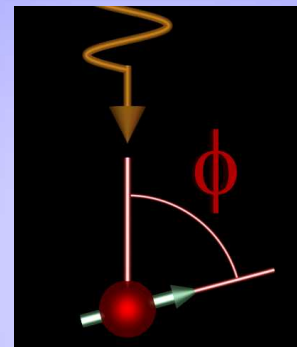
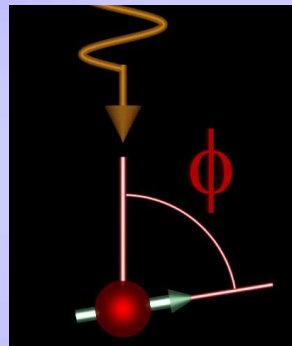
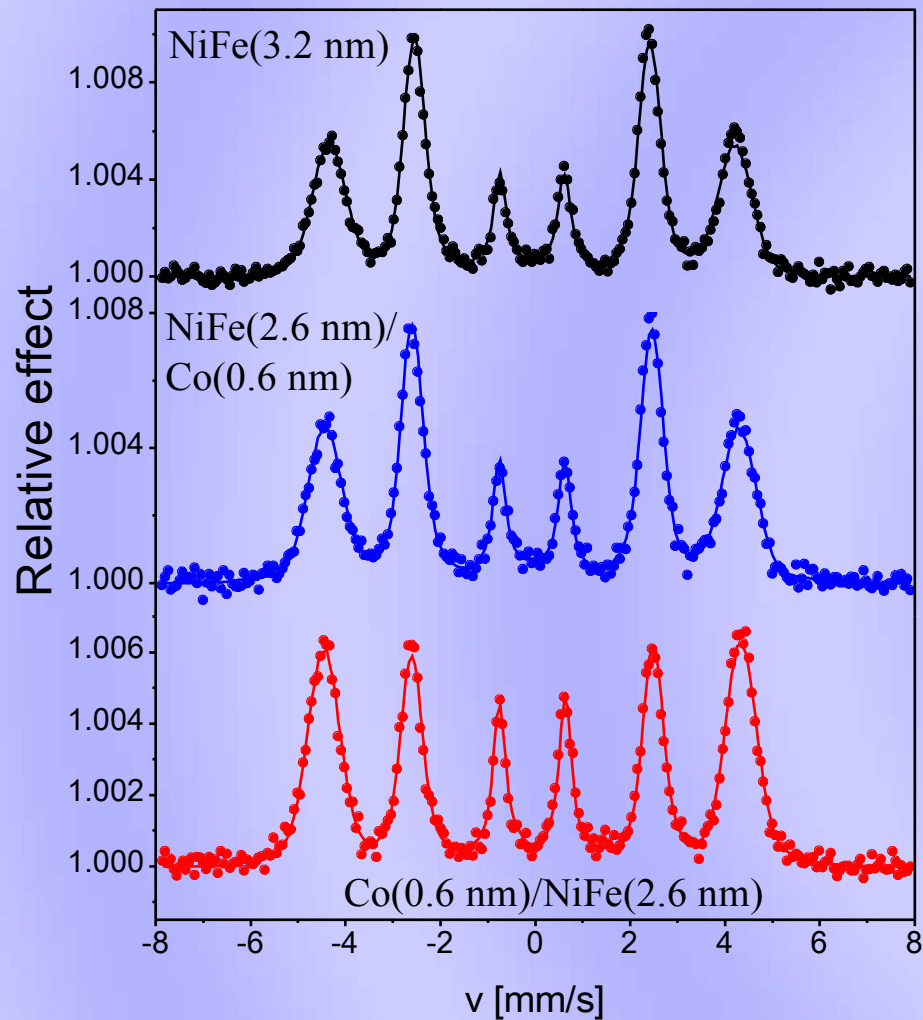




Relative intensities of the hyperfine lines vary with the angle ϕ between the incident γ -ray and the magnetic moment.

$$D_{23} = \frac{4 \sin^2(\phi)}{(1 + \cos^2(\phi))}$$



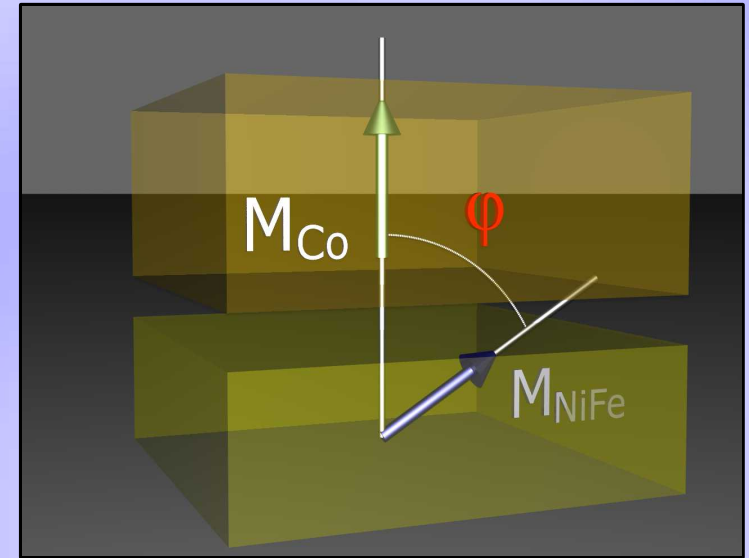
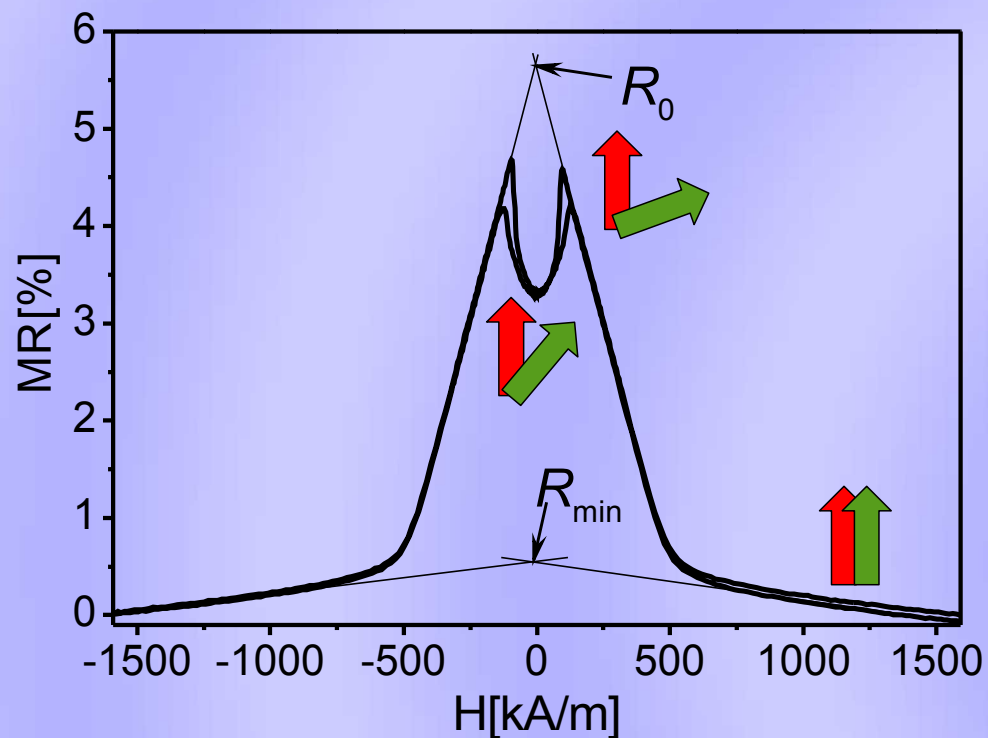


$[\text{Ni}_{80}\text{Fe}_{20}(3.2 \text{ nm})/\text{Au}(2.4 \text{ nm})/\text{Co}(0.8 \text{ nm})/\text{Au}(2.4 \text{ nm})]_{10}$

$[\text{Ni}_{80}\text{Fe}_{20}(2.6 \text{ nm})/\text{Co}(0.6 \text{ nm})/\text{Au}(2.4 \text{ nm})/\text{Co}(0.8 \text{ nm})/\text{Au}(2.4 \text{ nm})]_{10}$

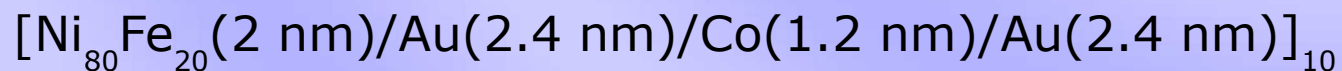
$[\text{Co}(0.6 \text{ nm})/\text{Ni}_{80}\text{Fe}_{20}(2.6 \text{ nm})/\text{Au}(2.4 \text{ nm})/\text{Co}(0.8 \text{ nm})/\text{Au}(2.4 \text{ nm})]_{10}$

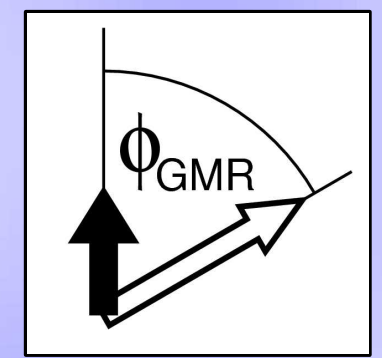
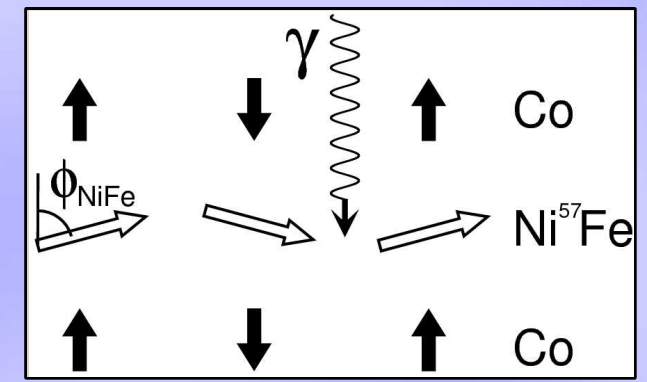
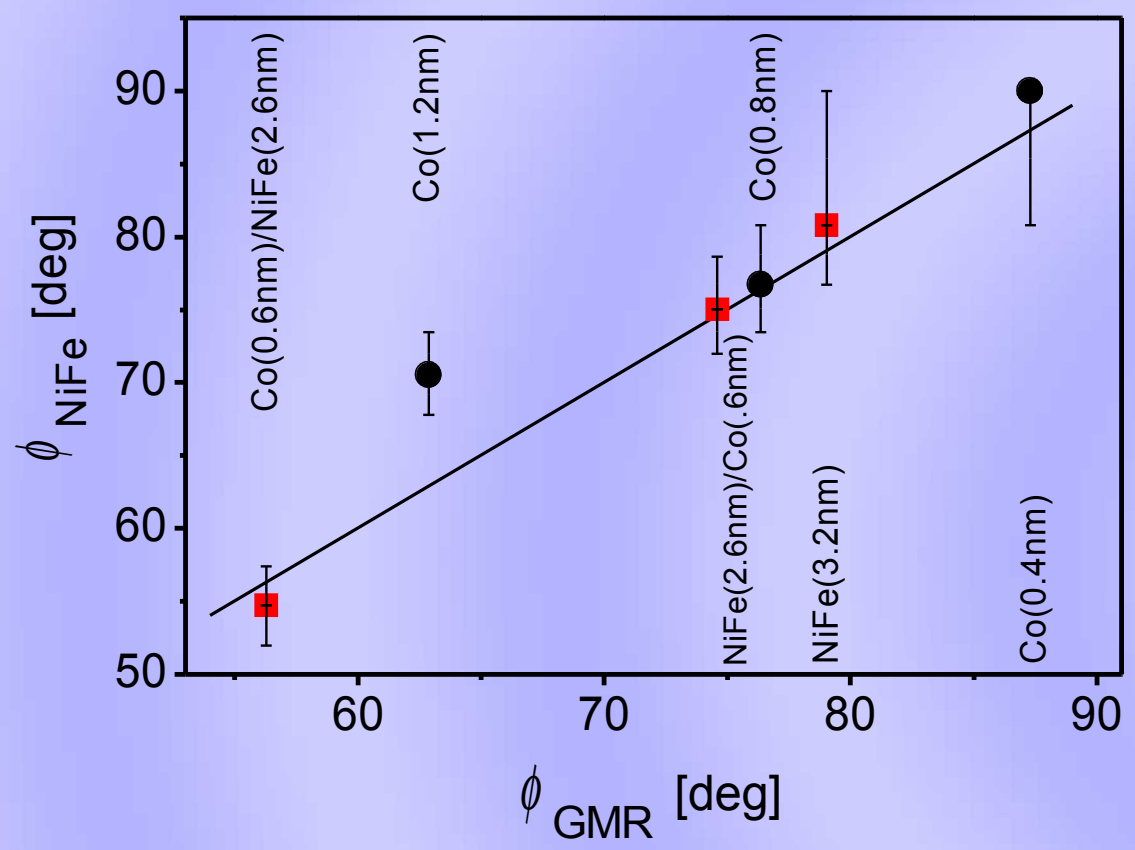
Element selective measurements



$$R = R_0 - (R_0 - R_{min}) \cos(\varphi_{Co-NiFe})$$

Resistance measurements allow the determination of the average cosine of the angle between magnetic moments of **Co** and **NiFe** layers.

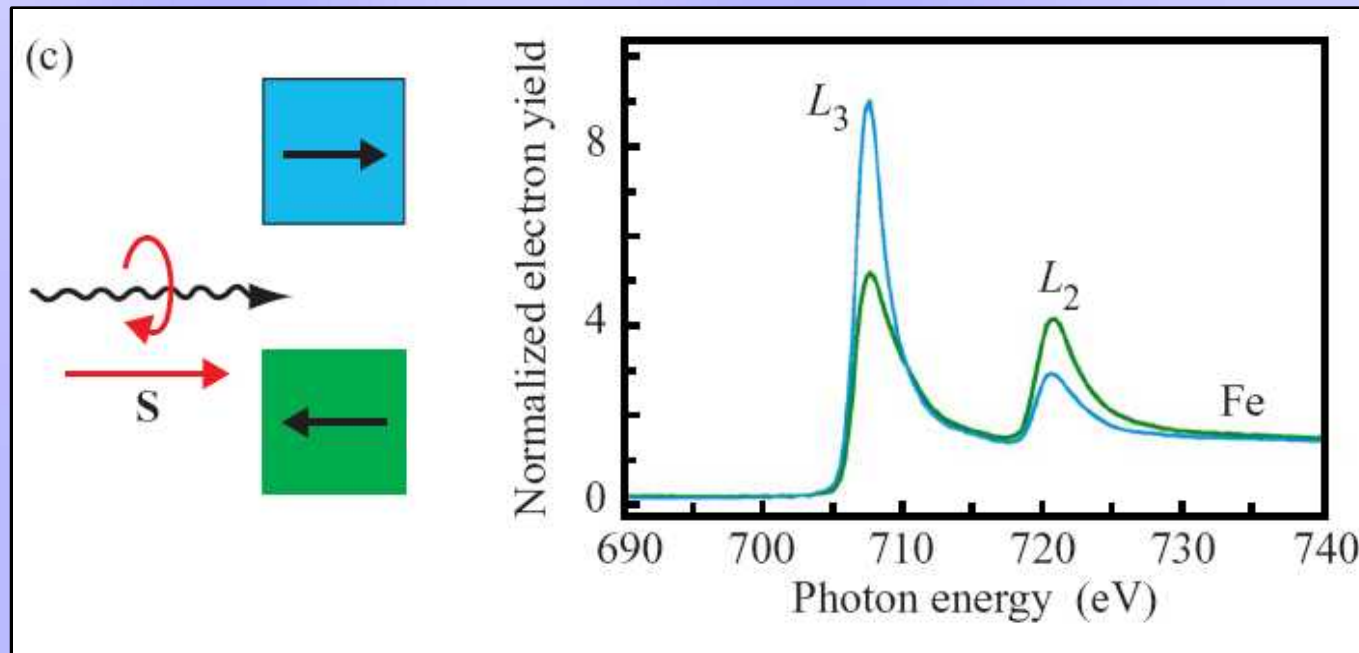




The magnetostatic fields of the Co domains cause the deflection of the magnetic moments of the NiFe layers. The deflection is stronger if the effective easy-plane anisotropy of NiFe layers is weaker.

- [**X**/Au(2.4 nm)/Co(0.8 nm)/Au(2.4 nm)]₁₀
- [Ni₈₀Fe₂₀(2 nm)/Au(2.4 nm)/**Co**/Au(2.4 nm)]₁₀

Soft x-ray resonant magnetic scattering (SXRMS)



Circularly polarized light

$\lambda \approx 1.4$ nm

interaction with core electrons

photon energy tuned to absorption edge



elemental selectivity

*

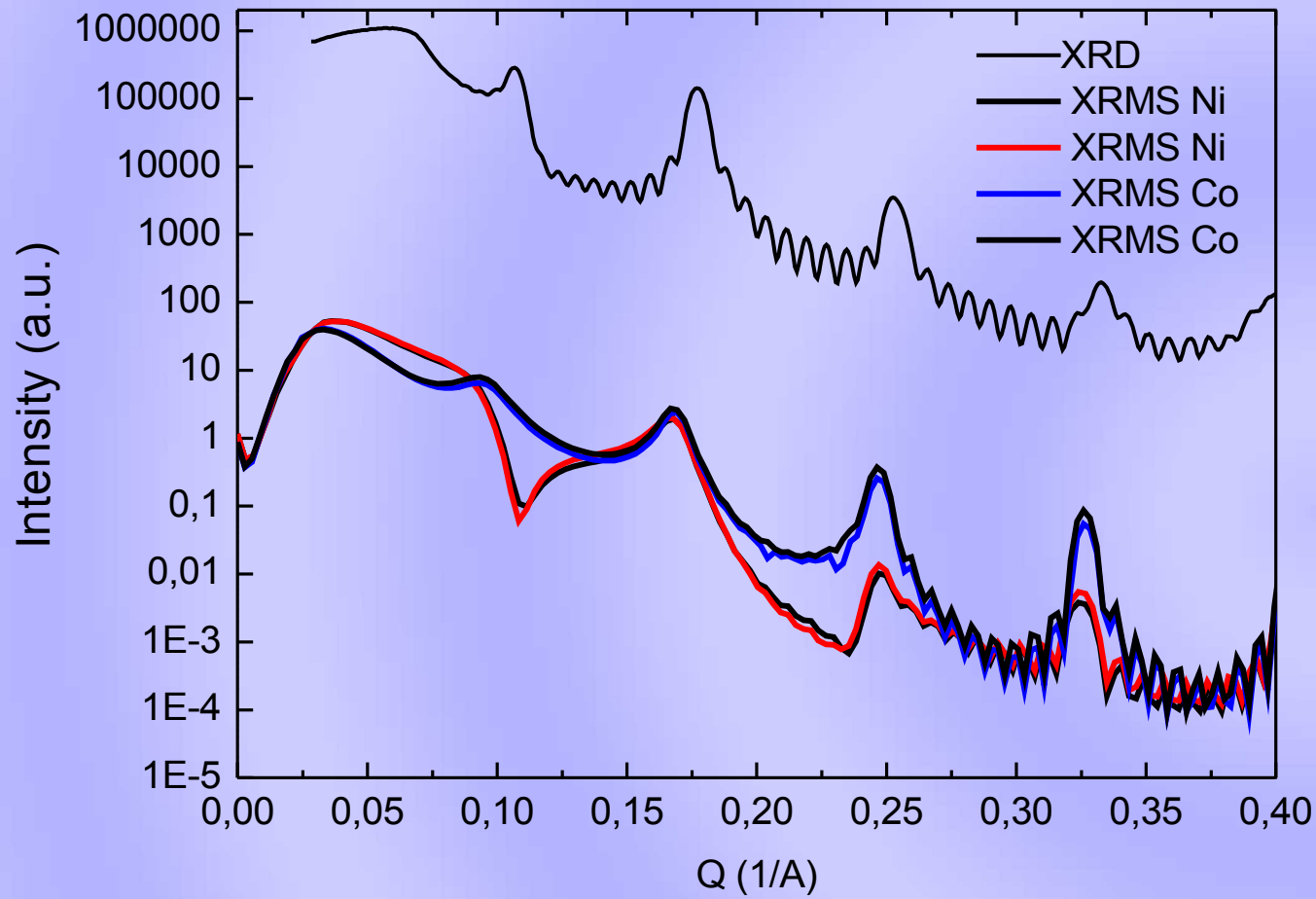
SXRMS at BESSY – measurement of the intensity of a reflected X-ray versus the external magnetic field (θ - 2θ geometry).

Sampling depth ~ 10 nm

ALICE diffractometer at the undulator beamline UE56/2-PGM2 at BESSY II (Berlin)

*graphics source: ssrl.slac.stanford.edu/stohr/xmcd.htm
see ssrl.slac.stanford.edu/stohr/X-Rays_and_Magnetism.ppt

Soft x-ray resonant magnetic scattering (SXRMS)



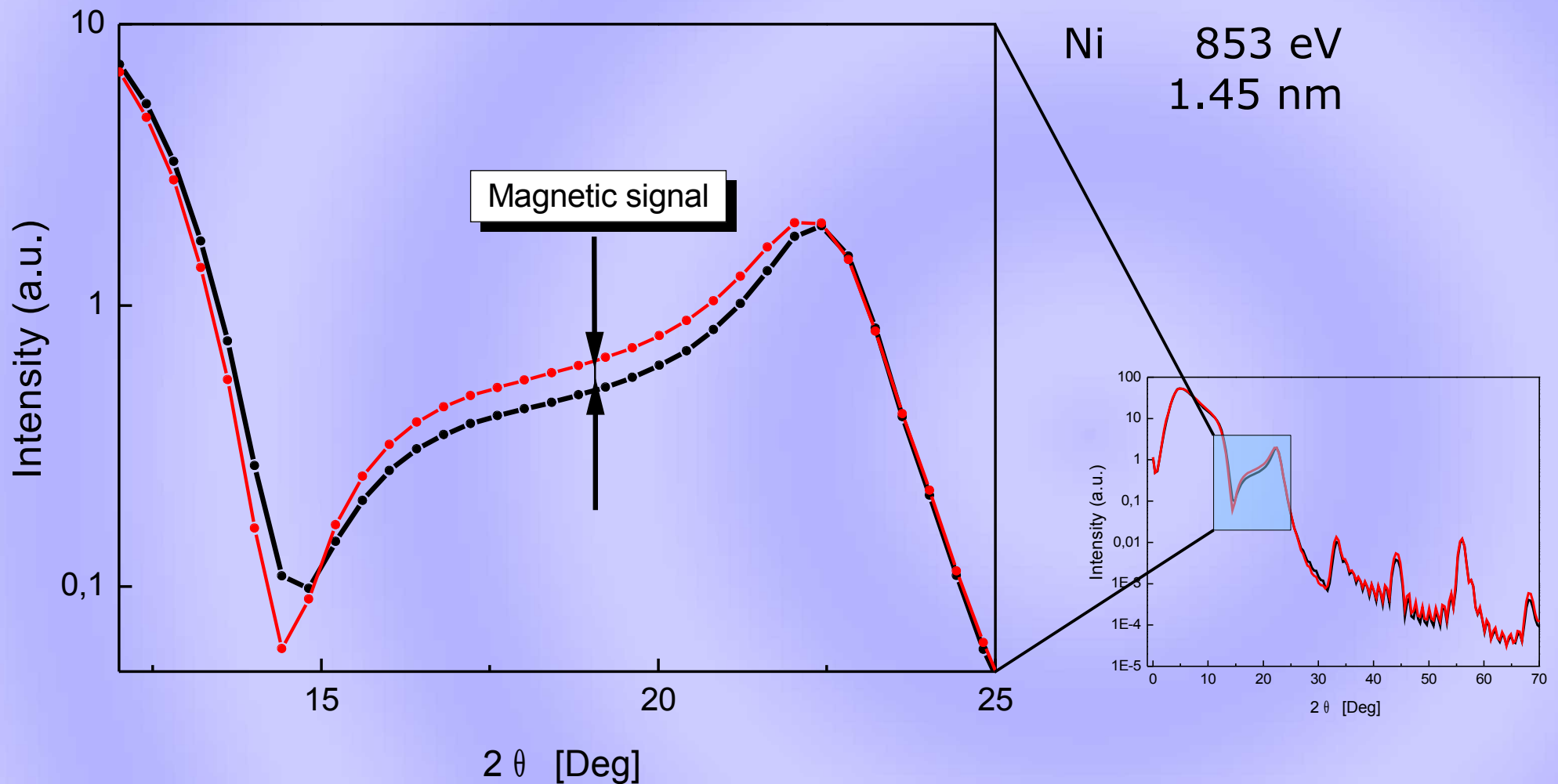
Cu K_α **8048** eV

Ni 853 eV

Co 778 eV

$[\text{Ni}_{80}\text{Fe}_{20}(2\text{ nm})/\text{Au}(2\text{ nm})/\text{Co}(0.8\text{ nm})/\text{Au}(2\text{ nm})]_{10}$

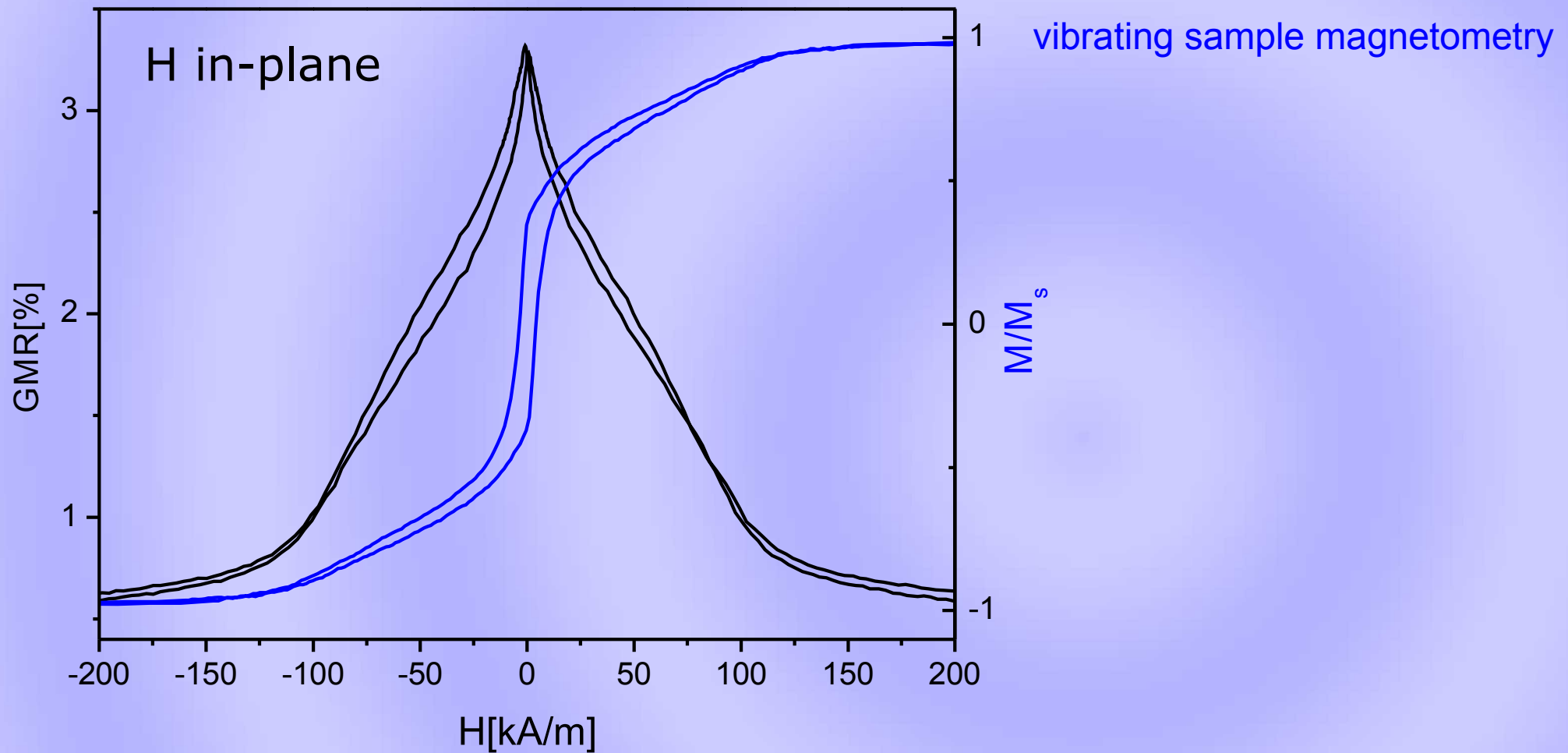
Soft x-ray resonant magnetic scattering (SXRMS)



$[\text{Ni}_{80}\text{Fe}_{20}(2\text{ nm})/\text{Au}(2\text{ nm})/\text{Co}(0.8\text{ nm})/\text{Au}(2\text{ nm})]_{10}$

ALICE diffractometer at the undulator beamline UE56/2-PGM2
at BESSY II (Berlin)

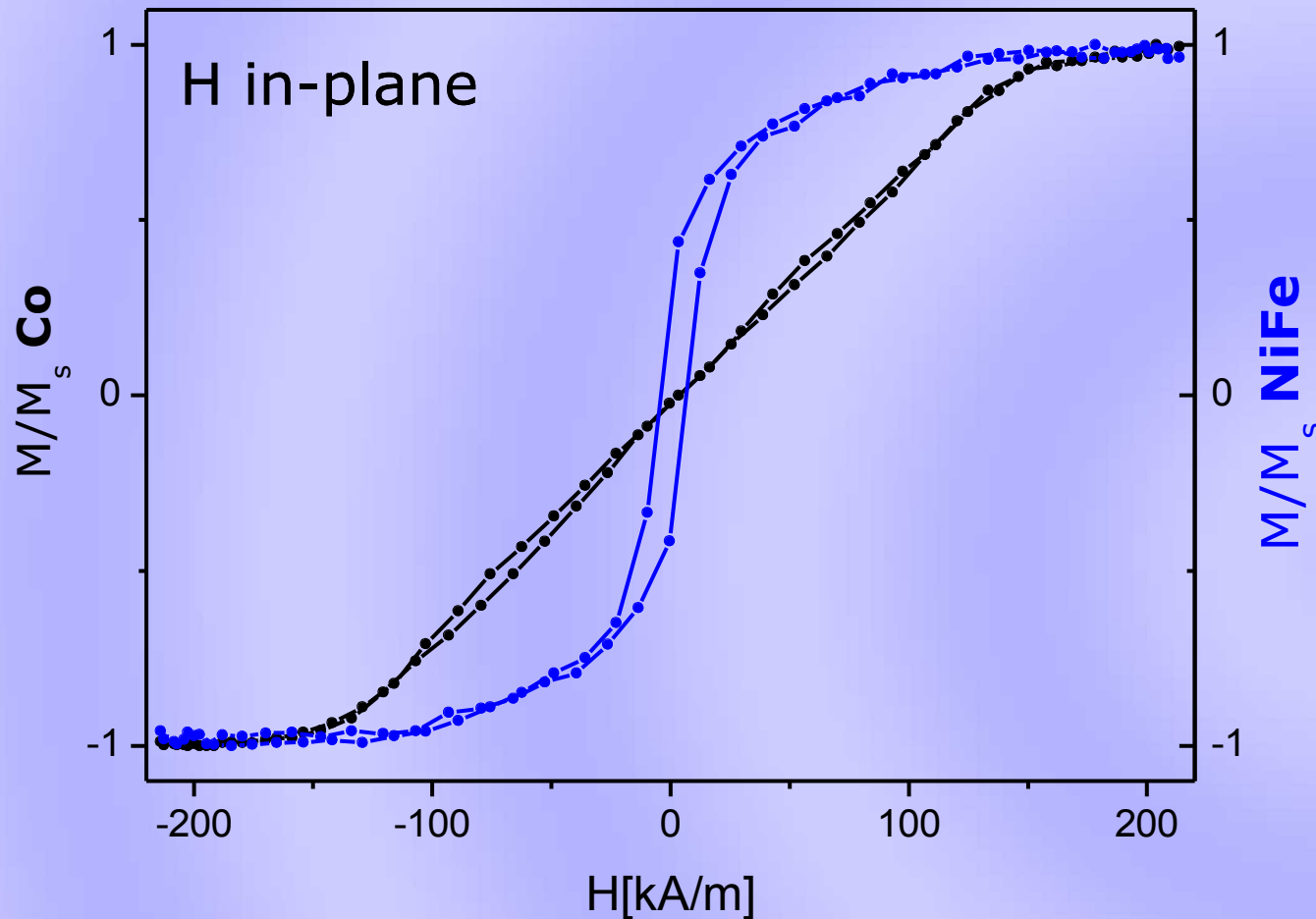
Soft x-ray resonant magnetic scattering (SXRMS)



$R(H) \leftrightarrow M(H)$

$[\text{Ni}_{80}\text{Fe}_{20}(2 \text{ nm})/\text{Au}(2 \text{ nm})/\text{Co}(1.1 \text{ nm})/\text{Au}(2 \text{ nm})]_{10}$

Soft x-ray resonant magnetic scattering (SXRMS)



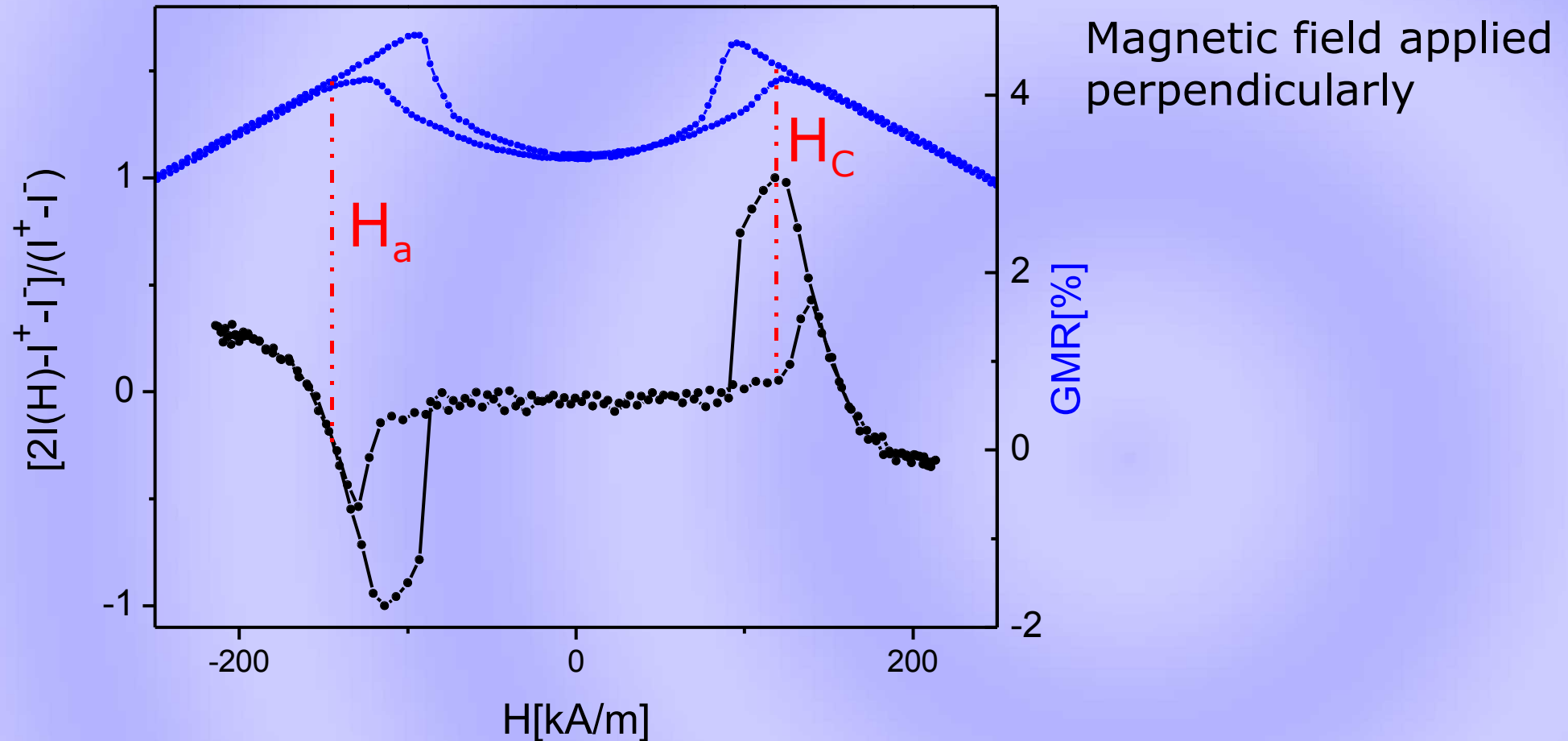
XRMS allows independent measurement of $M(H)$ dependence of **Co** and **NiFe** layers.

$\Theta = 8.5 \text{ Deg}$
 $E = 853 \text{ eV (Ni } L_3)$
 $E = 778 \text{ eV (Co } L_3)$

$$M/M_s \propto [2I(H) - I^+ - I^-] / (I^+ - I^-)$$

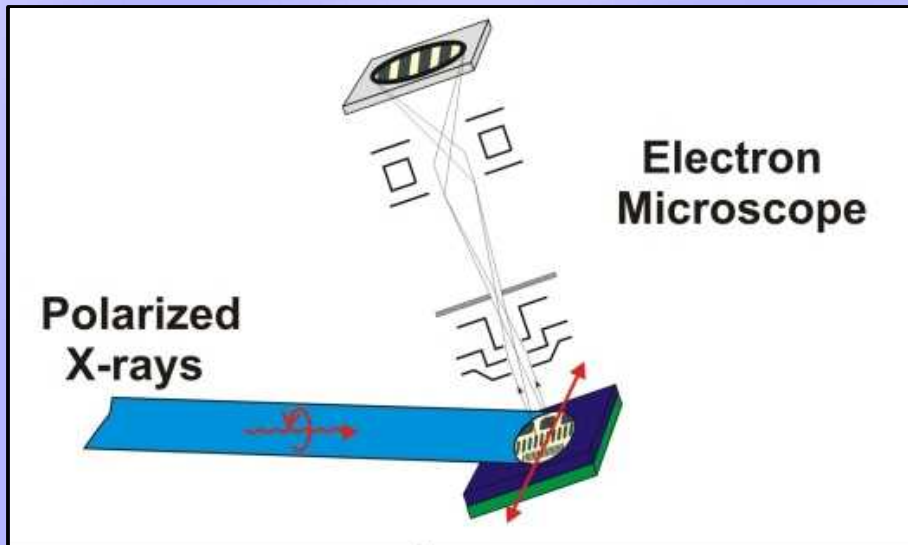
$[\text{Ni}_{80}\text{Fe}_{20} (2 \text{ nm}) / \text{Au} (2 \text{ nm}) / \text{Co} (1.1 \text{ nm}) / \text{Au} (2 \text{ nm})]_{10}$

Soft x-ray resonant magnetic scattering (SXRMS)

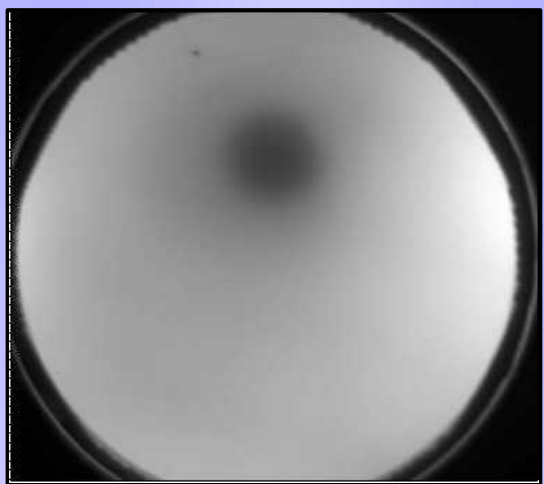


SXRMS signal from **NiFe layers** shows fields characteristic for Co layers reversal:
-creation of the stripe domain structure (H_C)
-annihilation field of domain structure (H_a)

$[\text{Ni}_{80}\text{Fe}_{20}(\mathbf{2\text{ nm}})/\text{Au}(2\text{ nm})/\text{Co}(1.1\text{ nm})/\text{Au}(2\text{ nm})]_{10}$

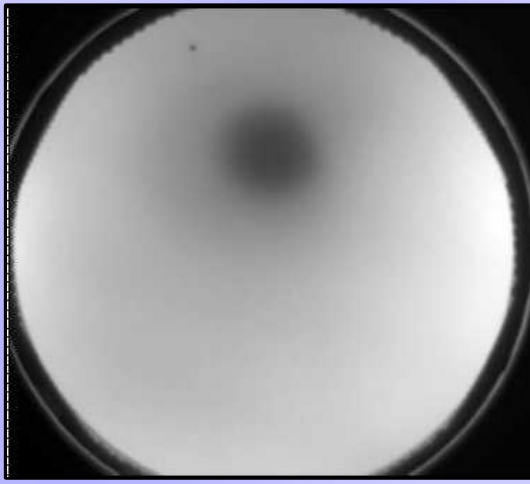


*graphics from: ssrl.slac.stanford.edu/stohr/xmcd.htm



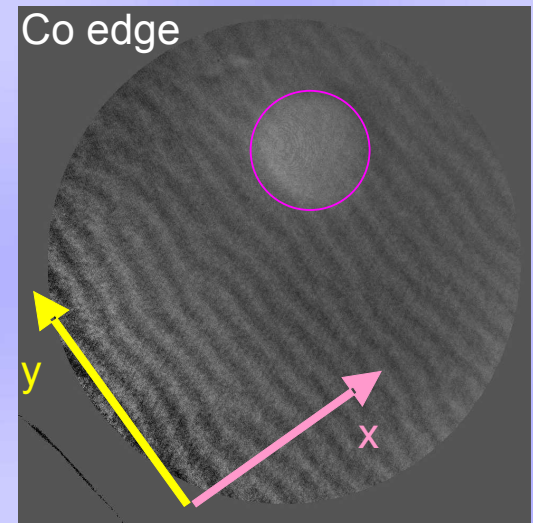
σ^L

-

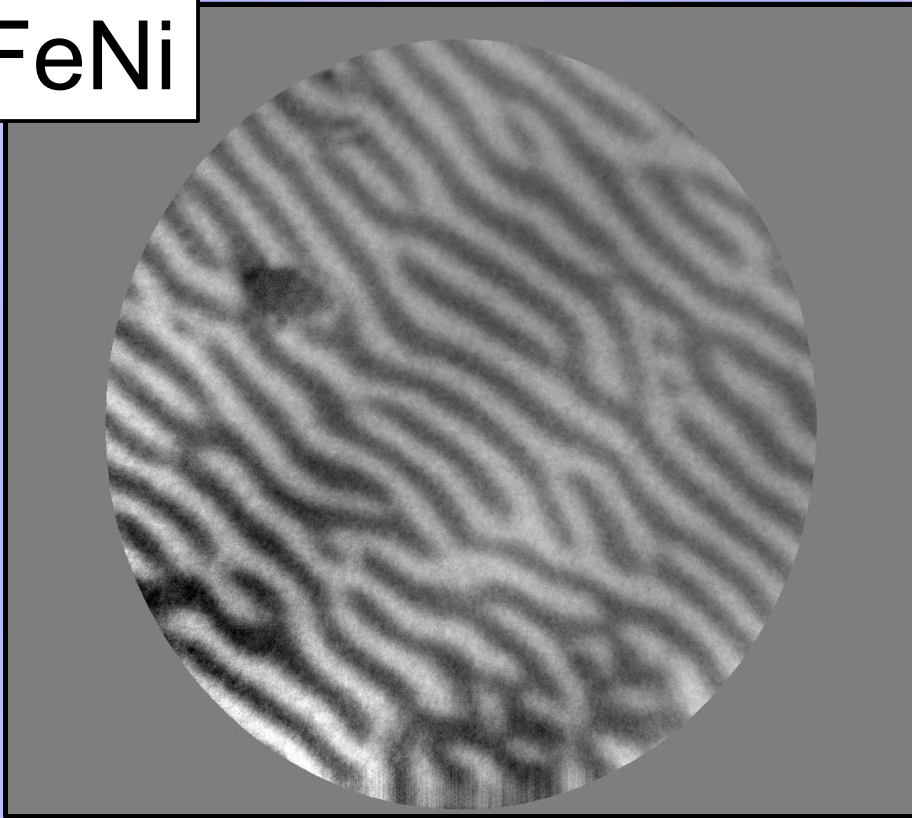


σ^R

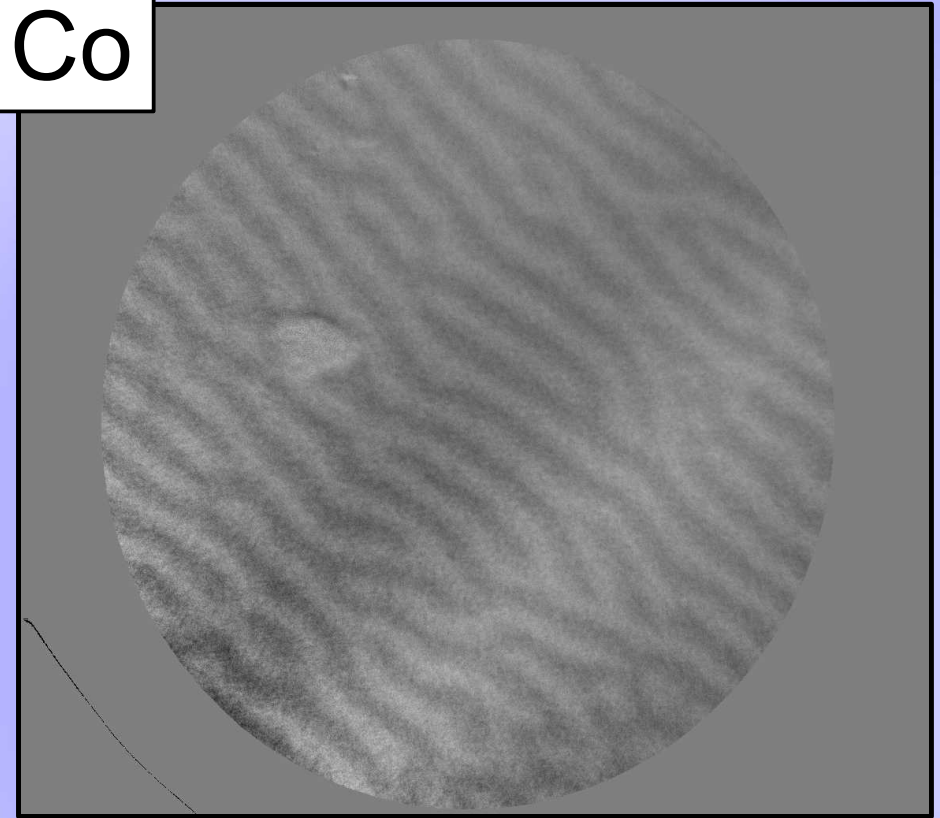
=



FeNi



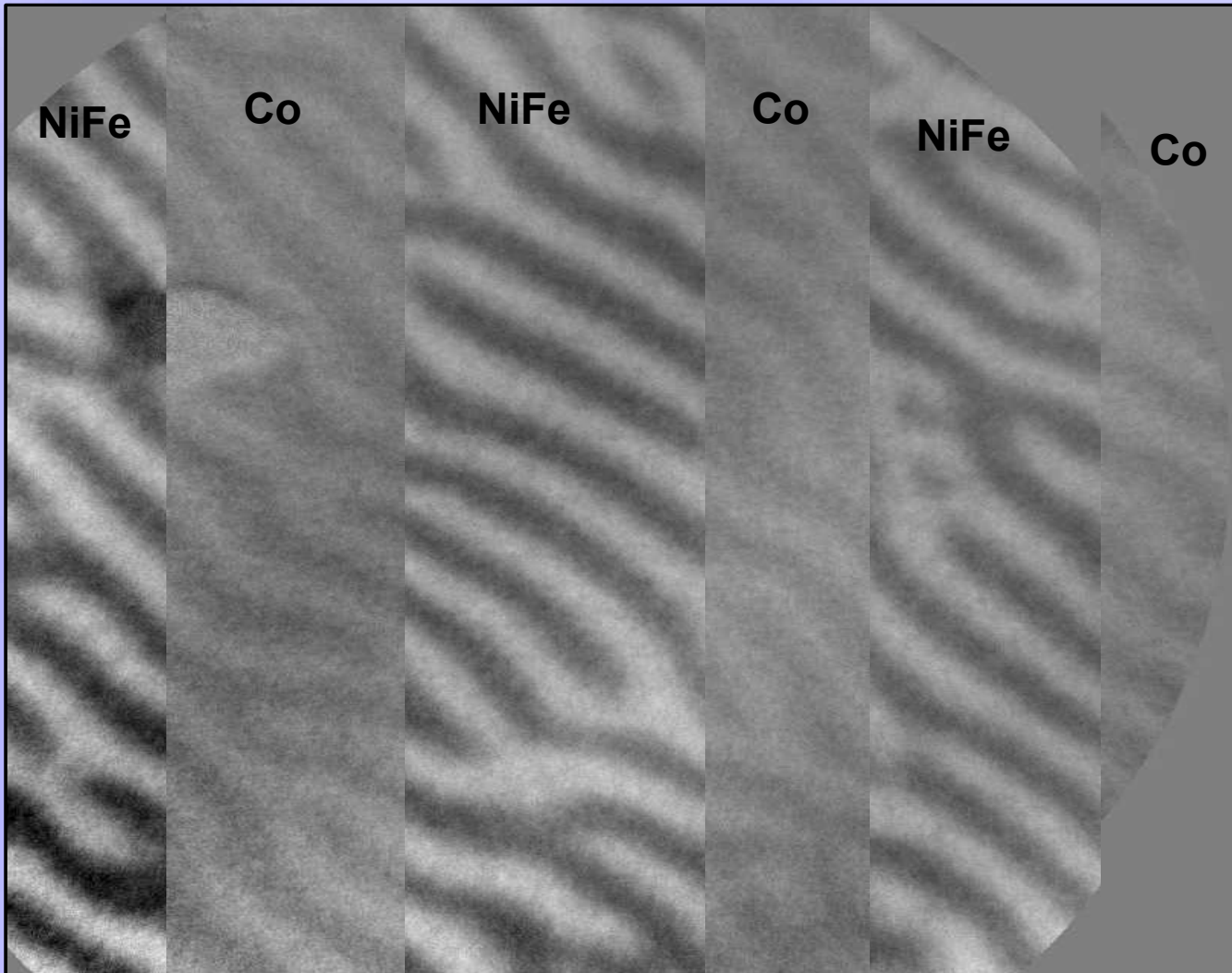
Co



Experimental confirmation of the replication of the Co stripe domains in the perpendicular component of NiFe sublayers magnetization.

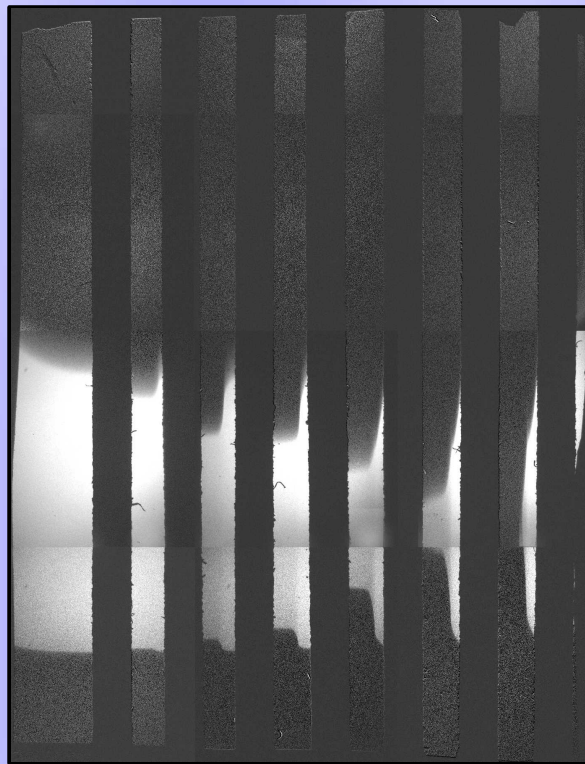
$[\text{Ni}_{80}\text{Fe}_{20}(2\text{ nm})/\text{Au}(2\text{ nm})/\text{Co}(0.8\text{ nm})/\text{Au}(2\text{ nm})]_{10}/\text{Ni}_{80}\text{Fe}_{20}(2\text{ nm})$

after in-plane ex-situ magnetizing in 0.7T

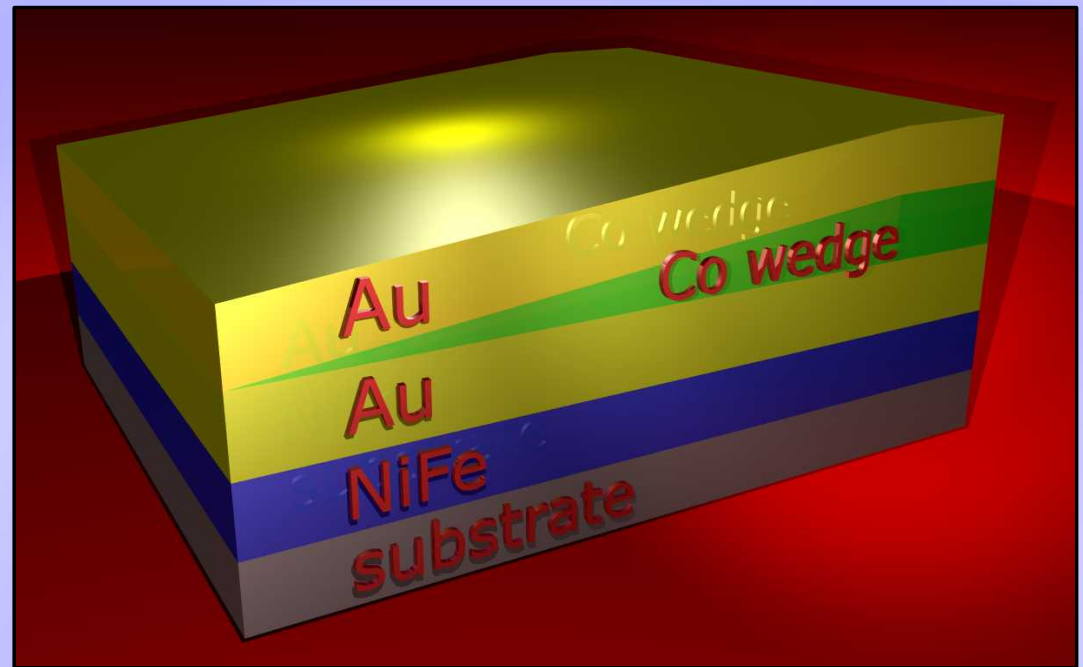


Experimental confirmation of the replication of the Co stripe domains in the perpendicular component of NiFe sublayers magnetization.

He Ion bombardment*



Co
4nm



Magneto-optical Kerr effect observation of magnetic structuration caused by ion bombardment:

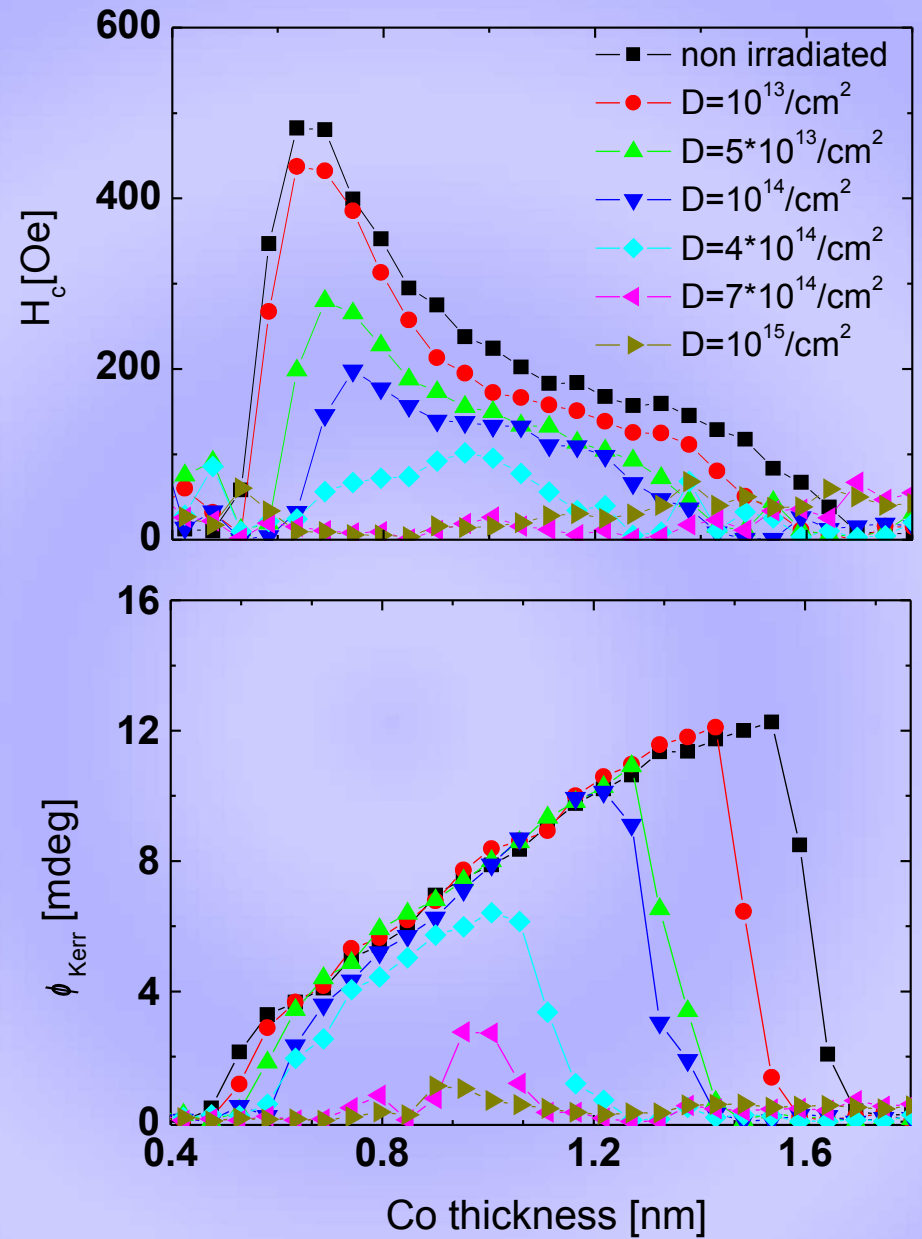
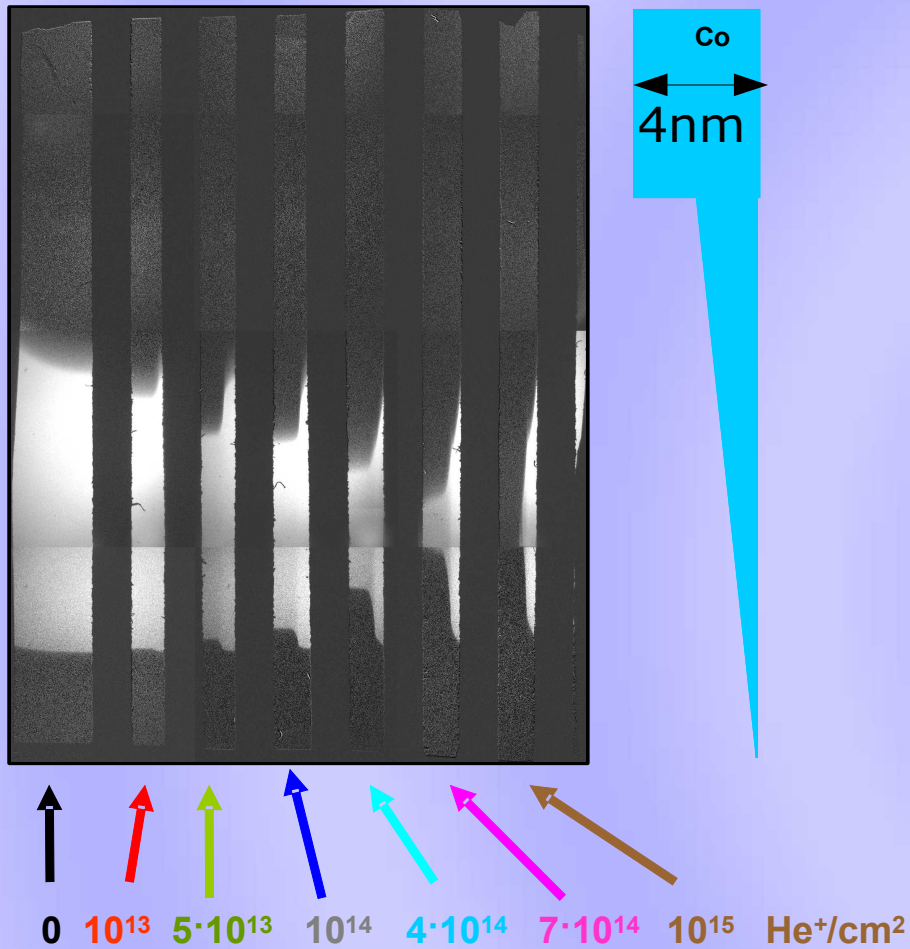
10keV He⁺

non-topological patterning

Si(100)/buffer/Ni₈₀Fe₂₀-2nm/Au-3nm/Co wedge/Au-3nm

*P. Kuświk *et al.*, *ACTA PHYSICA POLONICA A* 113, 651 (2008)

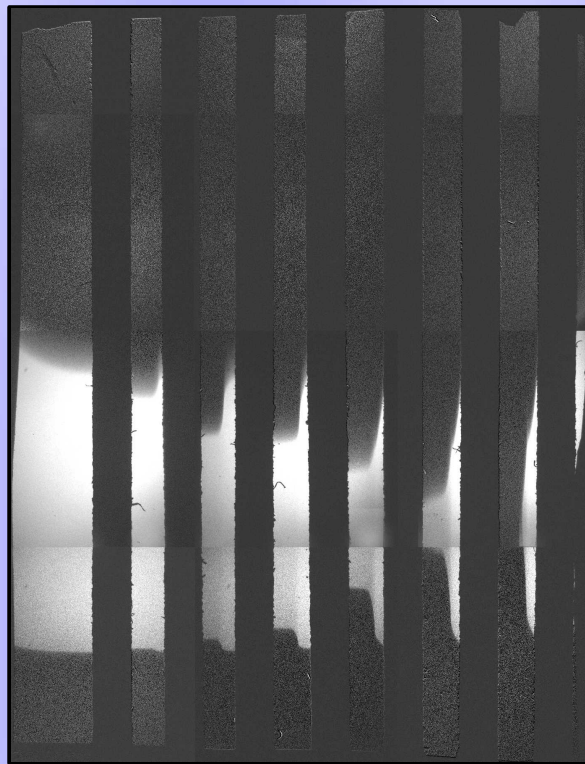
He Ion bombardment*



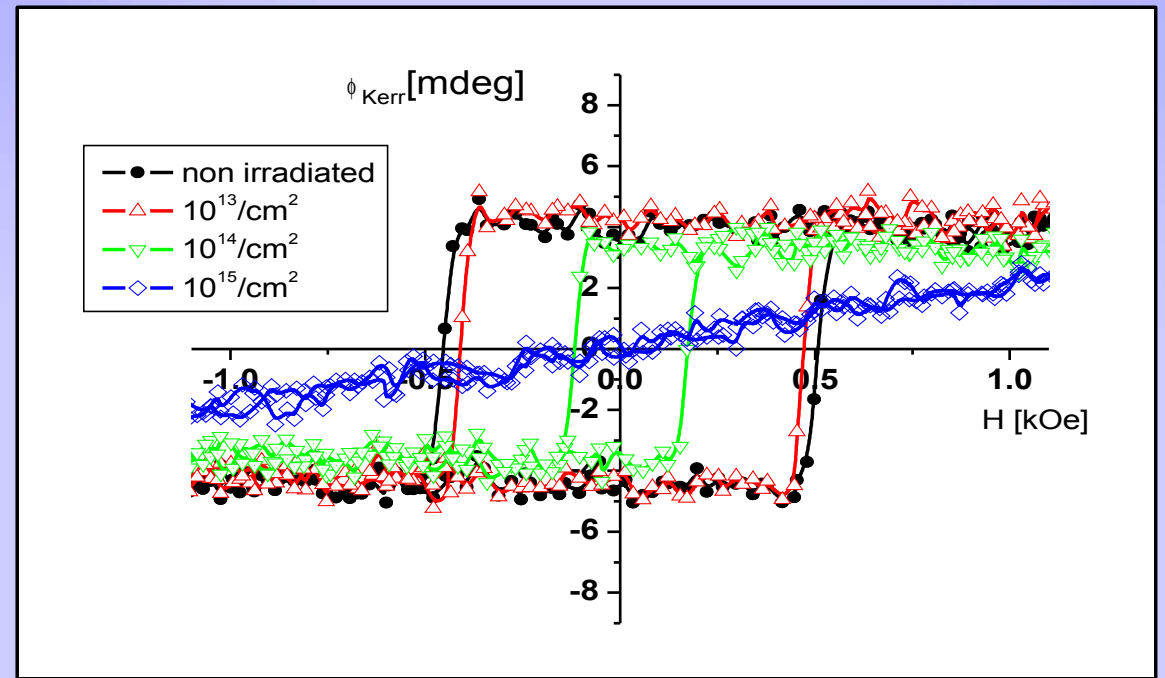
Si(100)/buffer/Ni₈₀Fe₂₀-2nm/Au-3nm/Co wedge/Au-3nm

*P. Kuświk *et al.*, *ACTA PHYSICA POLONICA A* 113, 651 (2008)

He Ion bombardment*



Co
4nm



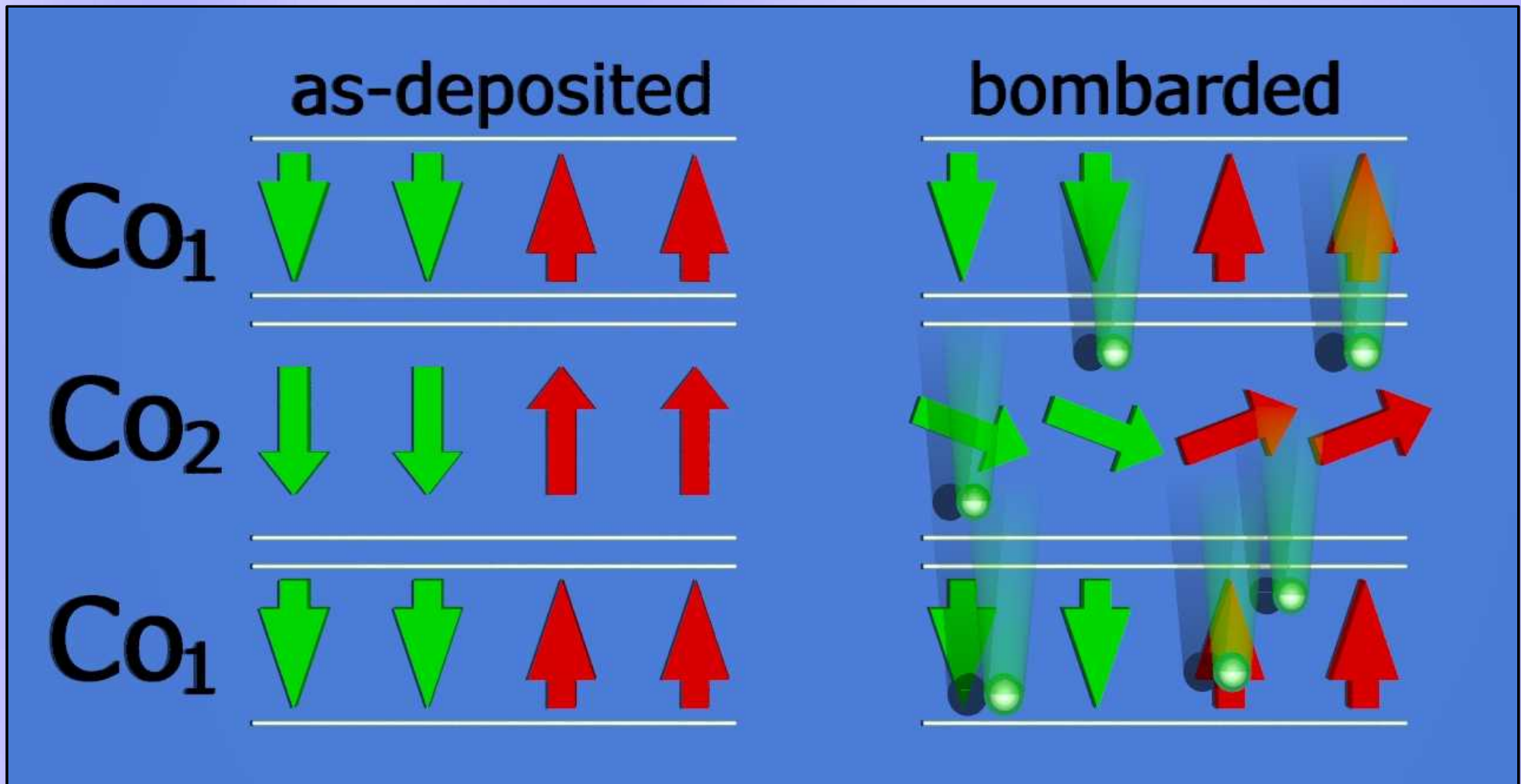
$t_{\text{Co}} = 0.6 \text{ nm}$, $M(H)$ for Co sublayers only

↑
0 10¹³ 5·10¹³ 10¹⁴ 4·10¹⁴ 7·10¹⁴ 10¹⁵ He⁺/cm²

Negligible changes of magnetic properties for doses $\leq 10^{13}$ ion/cm²

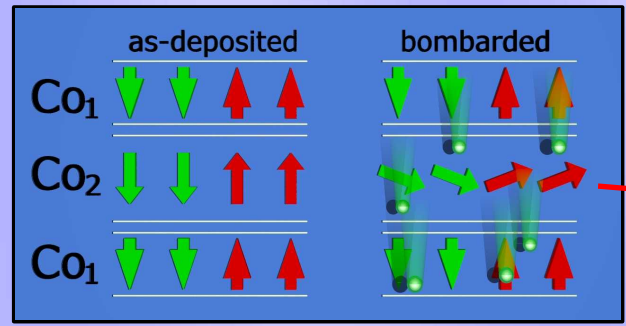
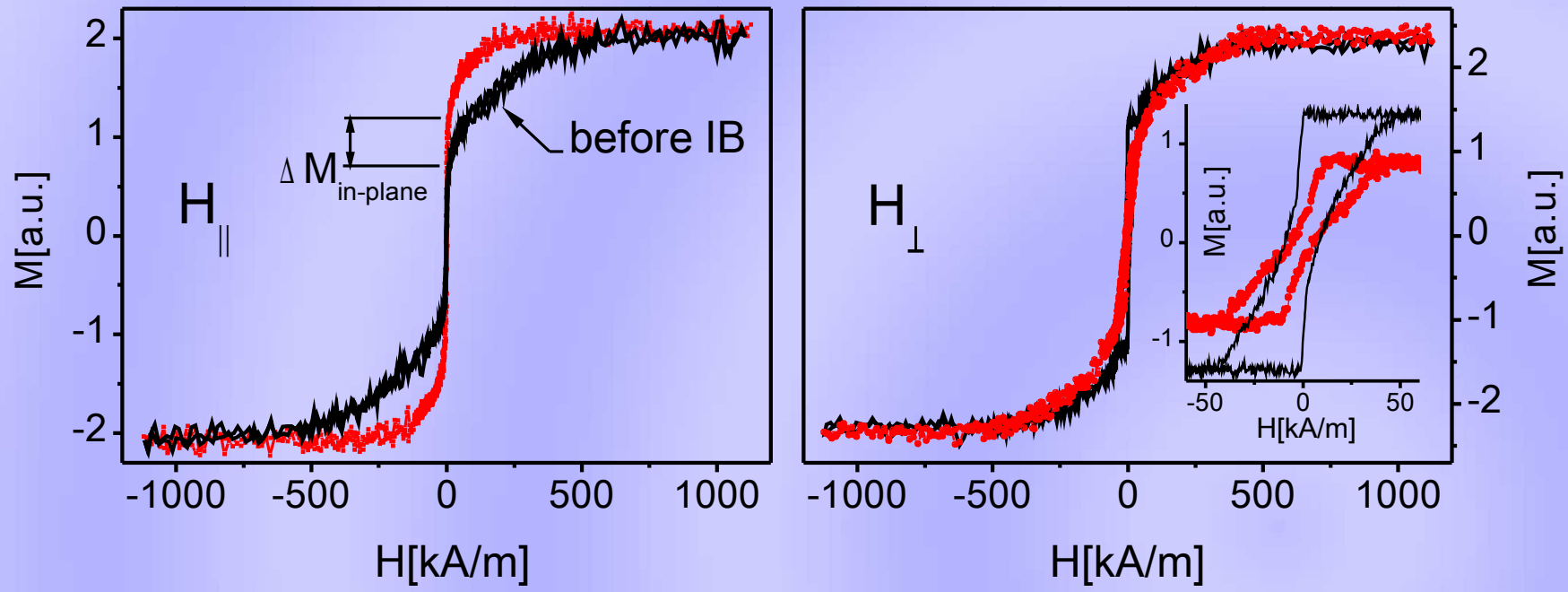
Si(100)/buffer/Ni₈₀Fe₂₀-2nm/Au-3nm/Co wedge/Au-3nm

*P. Kuświk *et al.*, *ACTA PHYSICA POLONICA A* 113, 651 (2008)



[Co₁(0.6 nm)/Au(4 nm)/Co₂(1 nm)/Au(4 nm)]₄

He Ion bombardment*

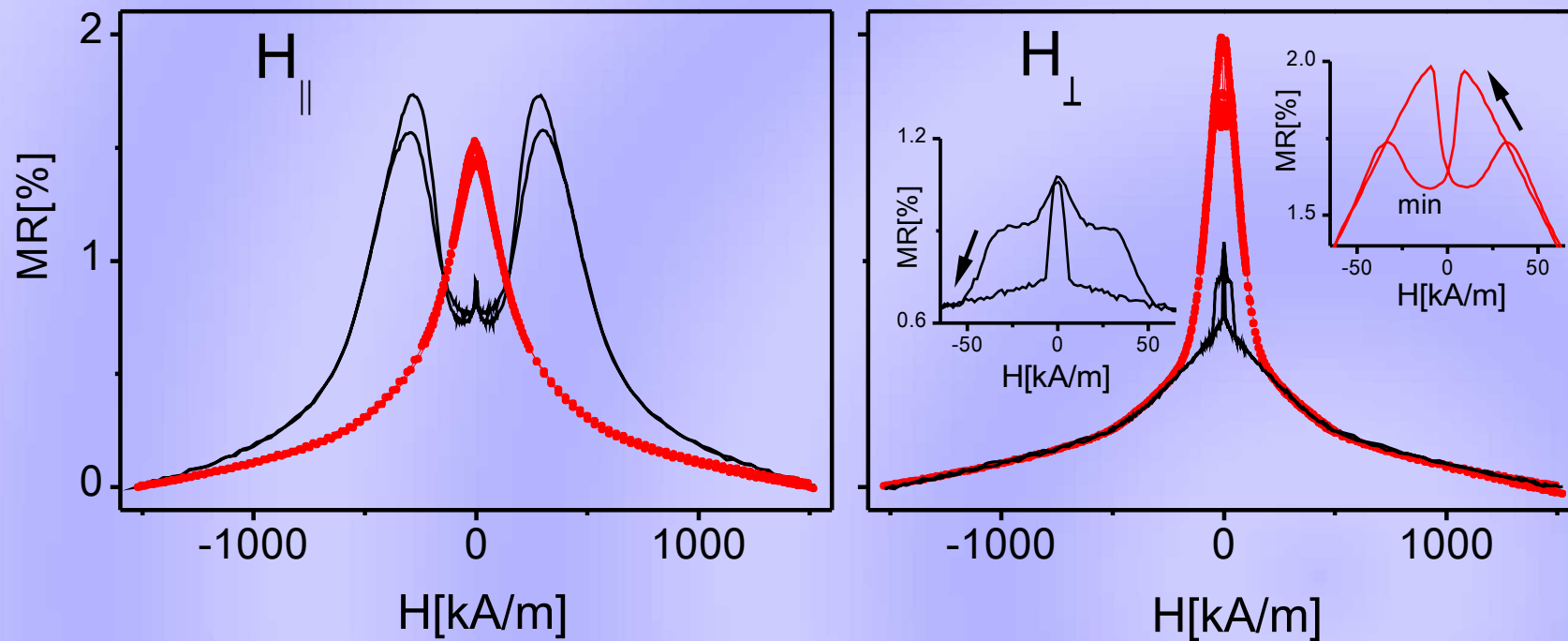


$$K_{eff} = \frac{2 K_{1s}}{t_{Co}} + K_{1v} - \frac{1}{2} \mu_0 (M_S^{Co})^2$$

[Co1(0.6 nm)/Au(4 nm)/Co2(1 nm)/Au(4 nm)]₄

*together with Ehresmann AG, Kassel (www.physik.uni-kassel.de/ehresmann)

He Ion bombardment



The resistance measurements confirm the observation inferred from the $M(H)$ measurements: the IB led to the switching of the EA direction in the 1 nm thick Co layers while the 0.6 nm thick layers preserved the perpendicular effective anisotropy.

$[\text{Co1}(0.6 \text{ nm})/\text{Au}(4 \text{ nm})/\text{Co2}(1 \text{ nm})/\text{Au}(4 \text{ nm})]_4$

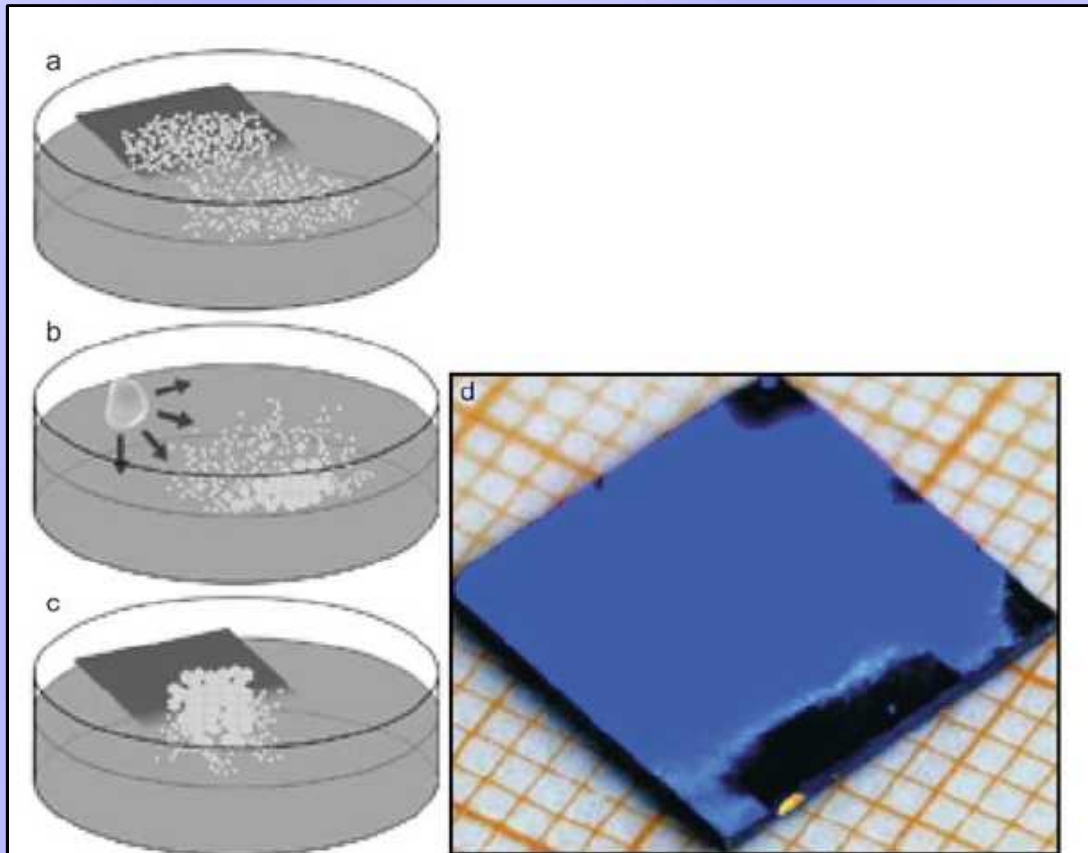
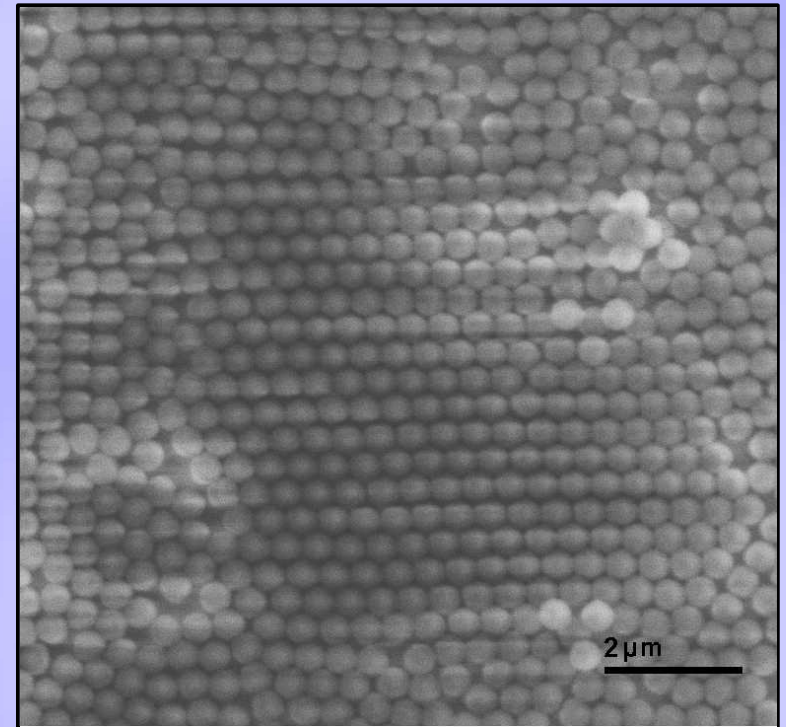
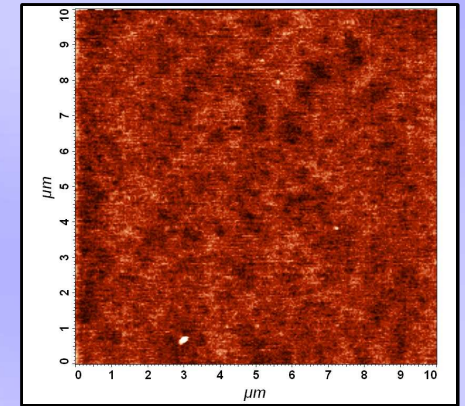
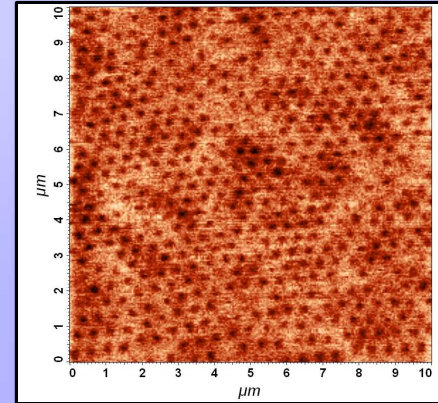
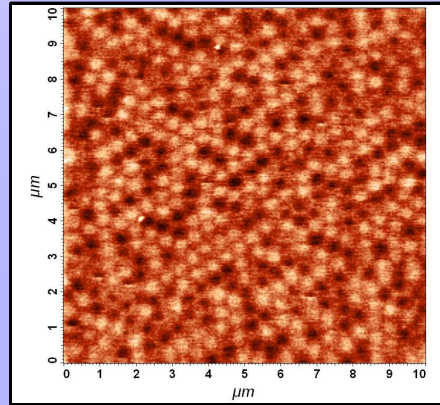
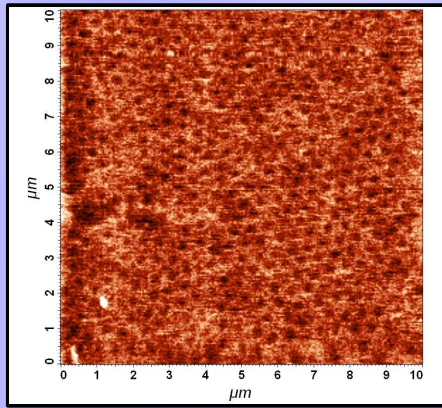


Fig. 1. The preparation process of monolayer masks: application of latex beads onto water surface (a); consolidation of particles (b); and liftoff of ordered monolayer (c). The $1 \times 1 \text{ cm}^2$ silicon wafer covered with monolayer built from 496 nm PS-latex beads: most of the surface does not contain any grain boundaries, which is represented as a monochrome light interference color (d).



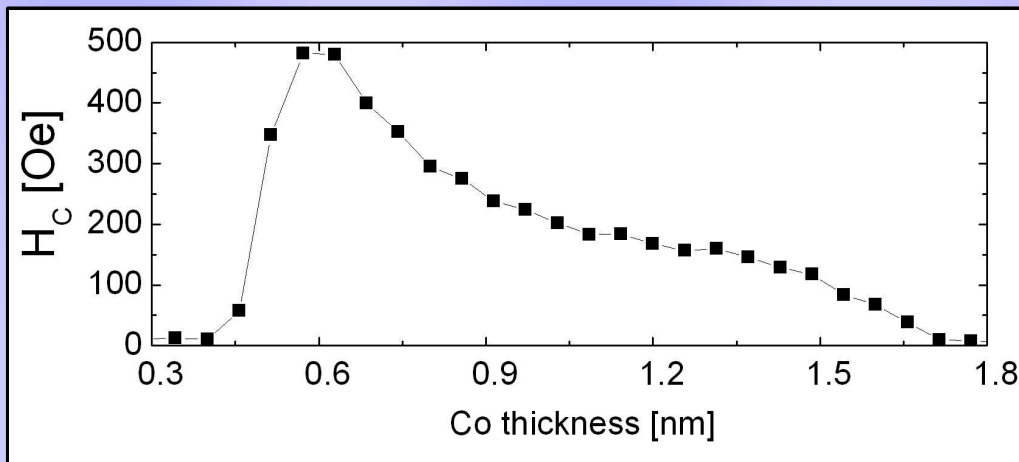
*W. Glapka, P. Kuświk *et al.*,
ACTA PHYSICA POLONICA A 115, 348 (2009)

polystyrene nanospheres
(**diameter 470 nm**) were
deposited on the multilayer
surfaces via a self-assembly
process realized by **a dip
coating**.



Co wedge

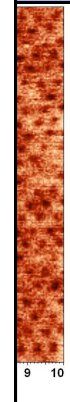
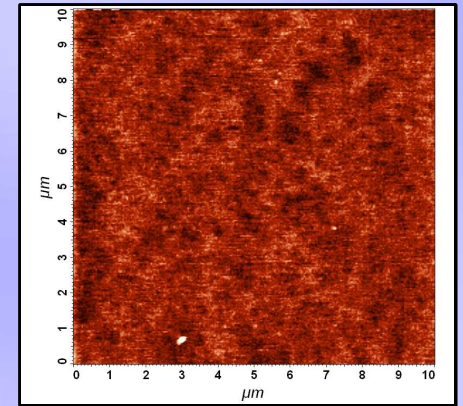
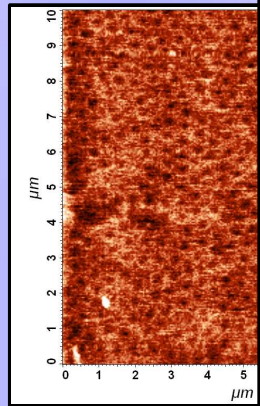
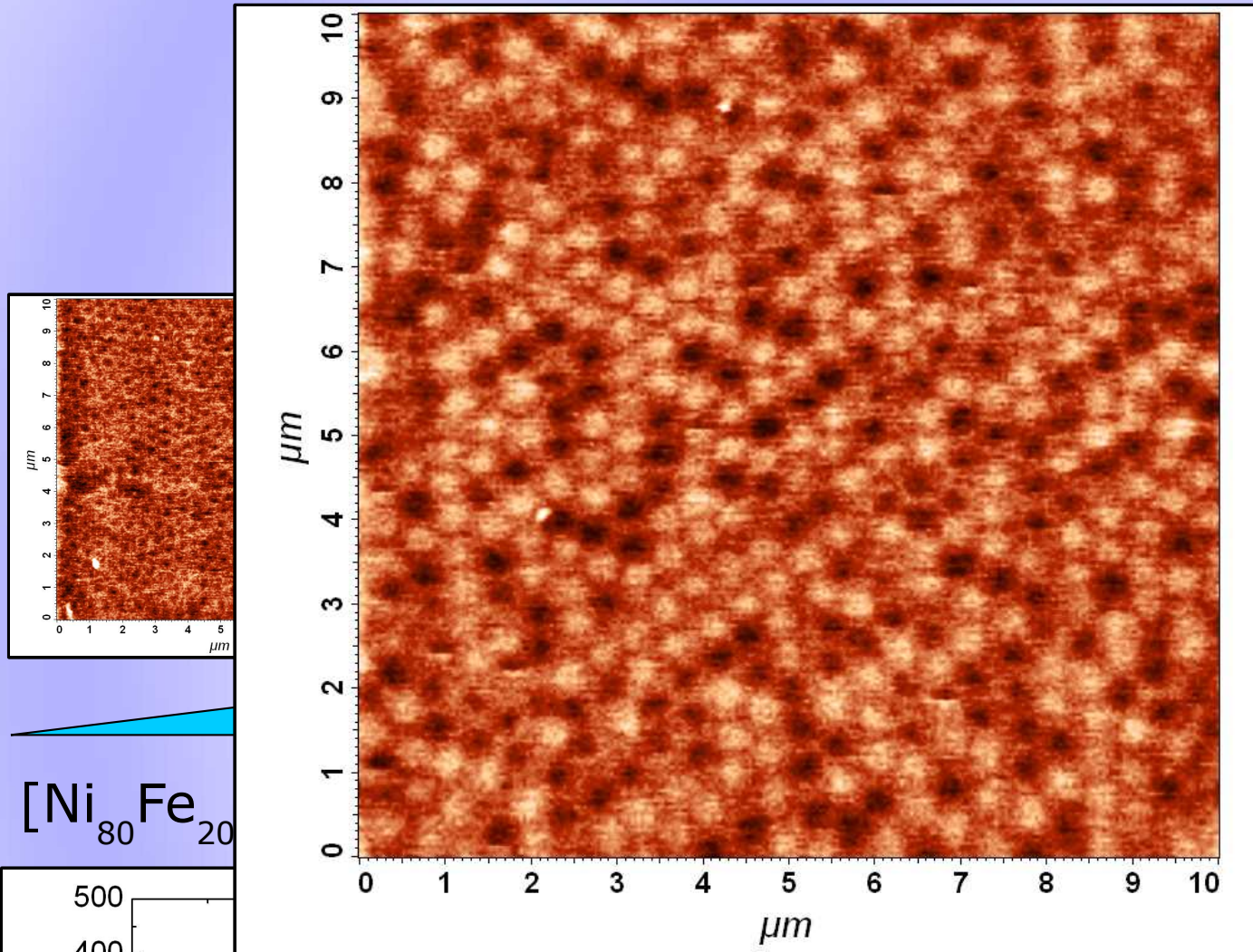
$[\text{Ni}_{80}\text{Fe}_{20} (2 \text{ nm})/\text{Au}(3 \text{ nm})/\text{Co}(\text{wedge})/\text{Au}(3 \text{ nm})]_{10}$



*W. Glapka, P. Kuświk *et al.*,
ACTA PHYSICA POLONICA A 115, 348 (2009)

He Ion bombardment

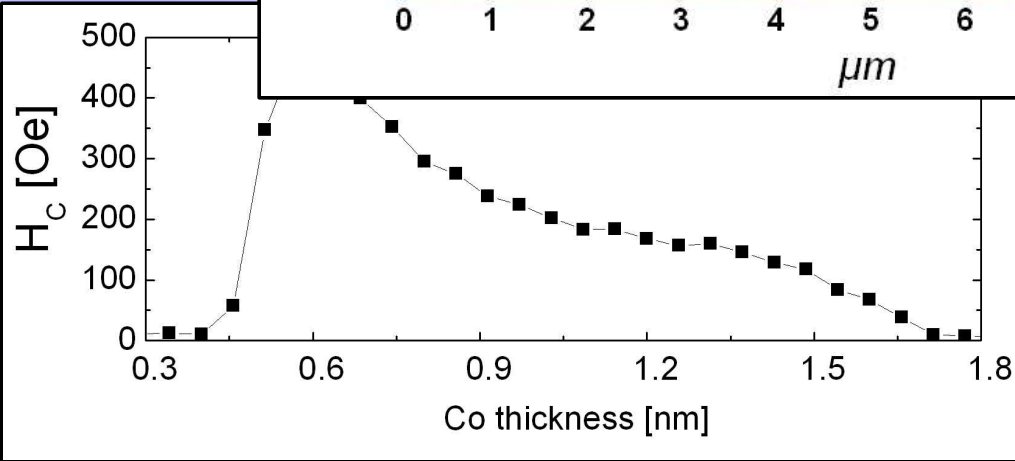
Magnetic patterning



[Ni₈₀Fe₂₀]

ge

n)]₁₀



*W. Glapka, P. Kuświk *et al.*,
ACTA PHYSICA POLONICA A 115, 348 (2009)

Conclusions

- $F_{\parallel}/\text{Au}/F_{\perp}/\text{Au}$ MIs represent new type of spin-valves
- Magnetostatic coupling influences magnetic reversal of $F_{\parallel}/\text{Au}/F_{\perp}$ MIs
- $F_{\parallel}/\text{Au}/F_{\perp}$ MIs are suitable for a non-topological magnetic patterning

Zakopane
2009

**Thank you for
your
attention**