

# Order in Quantum Compass and Orbital $e_g$ Models

Piotr Czarnik,<sup>1,2</sup> Jacek Dziarmaga,<sup>1</sup> and Andrzej M. Oleś<sup>1,3</sup>

<sup>1</sup>*Institute of Physics, Jagiellonian University, Kraków, Poland*

<sup>2</sup>*Institute of Physics, Universiteit van Amsterdam, The Netherlands*

<sup>3</sup>*Max Planck Institute for Solid State Research, Stuttgart, Germany*

Exchange interactions in orbital models are frustrated even on a square lattice, where two  $T = 1/2$  pseudospin components  $T_i^\gamma(\theta)$  parameterized by angle  $\theta \in (0, \pi/2]$  interact by terms  $JT_i^\gamma(\theta)T_j^\gamma(\theta)$ . Maximal frustration in the quantum compass model with  $T_i^\gamma(\pi/2) \equiv \frac{1}{2}\sigma_i^\gamma$ , where  $\sigma_i^\gamma$  is the Pauli matrix, is reduced to moderate frustration for the  $e_g$  orbital model at  $\theta = \pi/3$  [1]. We investigate thermodynamic phase transitions at temperature  $T_c$  on an infinite square lattice by variational tensor network renormalization (VTNR) in imaginary time. From the linear susceptibility (order parameter) in the symmetric (symmetry-broken) phase the onset of nematic order in the quantum compass model is estimated at  $T_c/J = 0.0606(4)$  [2], in good agreement with Quantum Monte Carlo (QMC). For the  $e_g$  orbital model one finds: (i) a very accurate VTNR estimate of  $T_c/J = 0.3566 \pm 0.0001$  while QMC fails due to the sign problem, and (ii) that the critical exponents are within the Ising universality class. Remarkably large difference in frustration and entanglement results in so distinct  $T_c$ .

## References:

- [1] L. Cincio, J. Dziarmaga, and A. M. Oleś, Phys. Rev. B **82**, 104416 (2010).
- [2] P. Czarnik, J. Dziarmaga, and A. M. Oleś, Phys. Rev. B **93**, 184410 (2016).